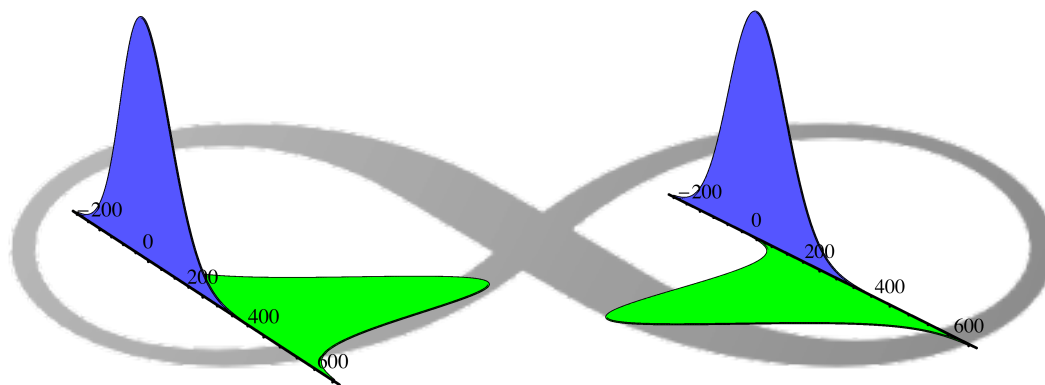


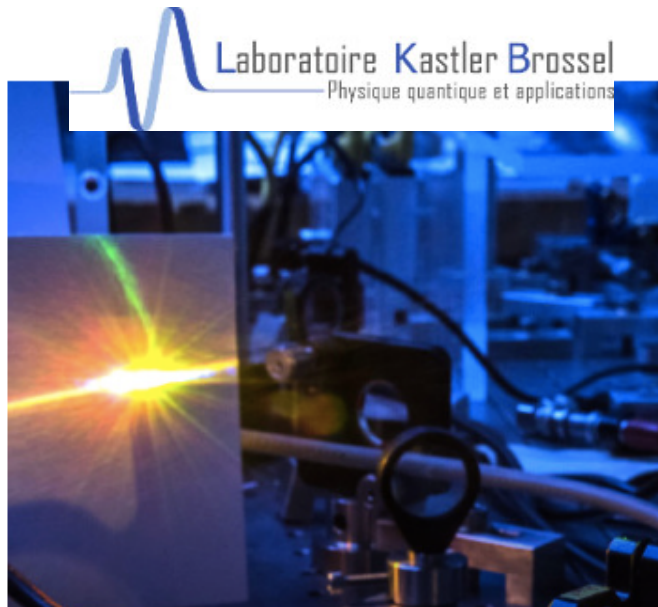
Quantum multimode resources based on optical frequency combs and simulation of complex quantum network

Valentina Parigi

Multimode quantum optics group



QuProCSII 07 April 2017



MULTIMODE QUANTUM OPTICS

NICOLAS TREPS, CLAUDE FABRE,
VALENTINA PARIGI



*C. Fabre, N. Trep*s

Femto Lab

Optical Fiber Networks

- Quantum Frequency Combs
- Quantum Information Processing
- Multimode analysis of laser dynamics
- Quantum metrology
- Complex quantum networks

CAILabs
Shaping the light


QCUMBER
Quantum Controlled Ultrafast Multimode
Entanglement and Measurement

ANR

SORBONNE
UNIVERSITÉS

Outline

Multimode quantum optics in Quantum Information technologies

Introduction

Multicolor entanglement

Towards measurement based quantum computing

Simulation of complex quantum networks

Probing a structured environment

Energy transport: some ideas

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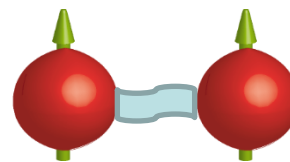
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Probing a structured environment

Energy transport: some ideas



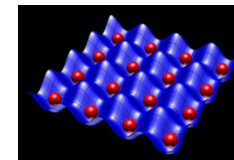
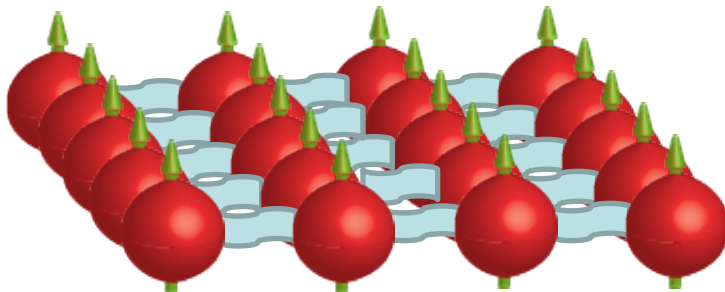
+



Quantum systems (q-bits)

Quantum Correlations (entanglement)

Quantum systems on large scale



Example
atoms

Technology

Real advantage over classical

Quantum communications

Safer (cryptography)

Quantum computation

Faster (prime factorization)

Quantum simulation

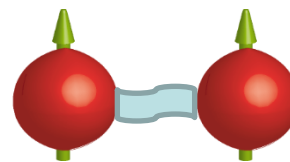
Solution to hard problem

Quantum metrology

More sensitive



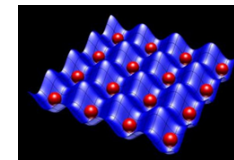
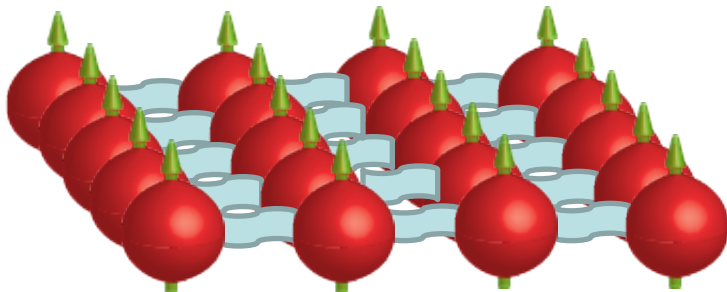
+



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Quantum systems on **large scale**



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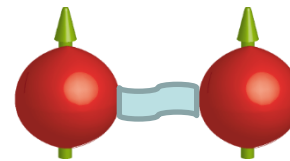
Quantum metrology

More sensitive

...with multimode quantum optics



+

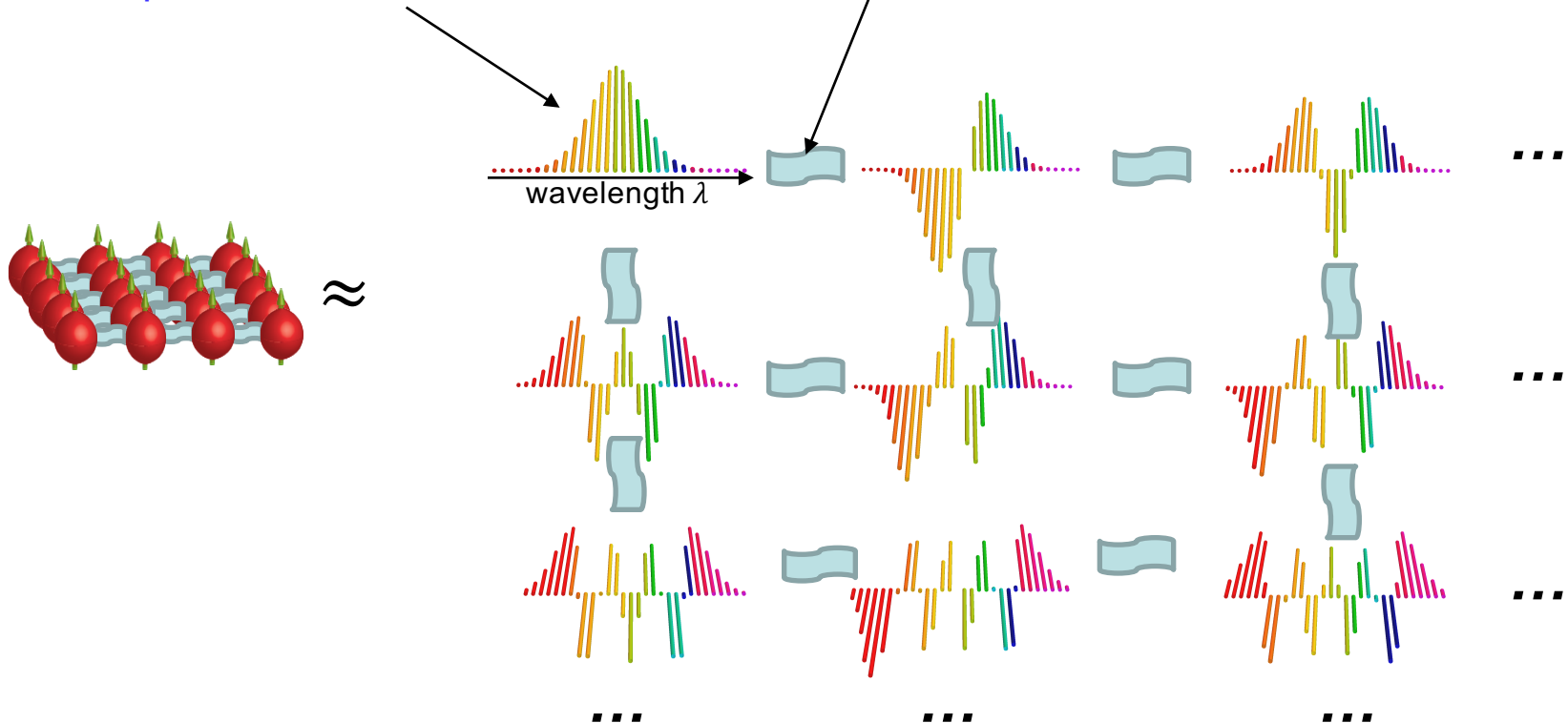


Quantum systems

spatial-temporal modes
of the electromagnetic field
= independent harmonic oscillators

Quantum Correlations (entanglement)

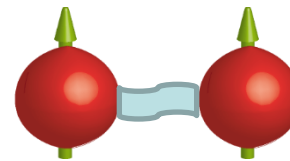
Entanglement between quadratures
(i.e. amplitude and phase)



...with multimode quantum optics



+

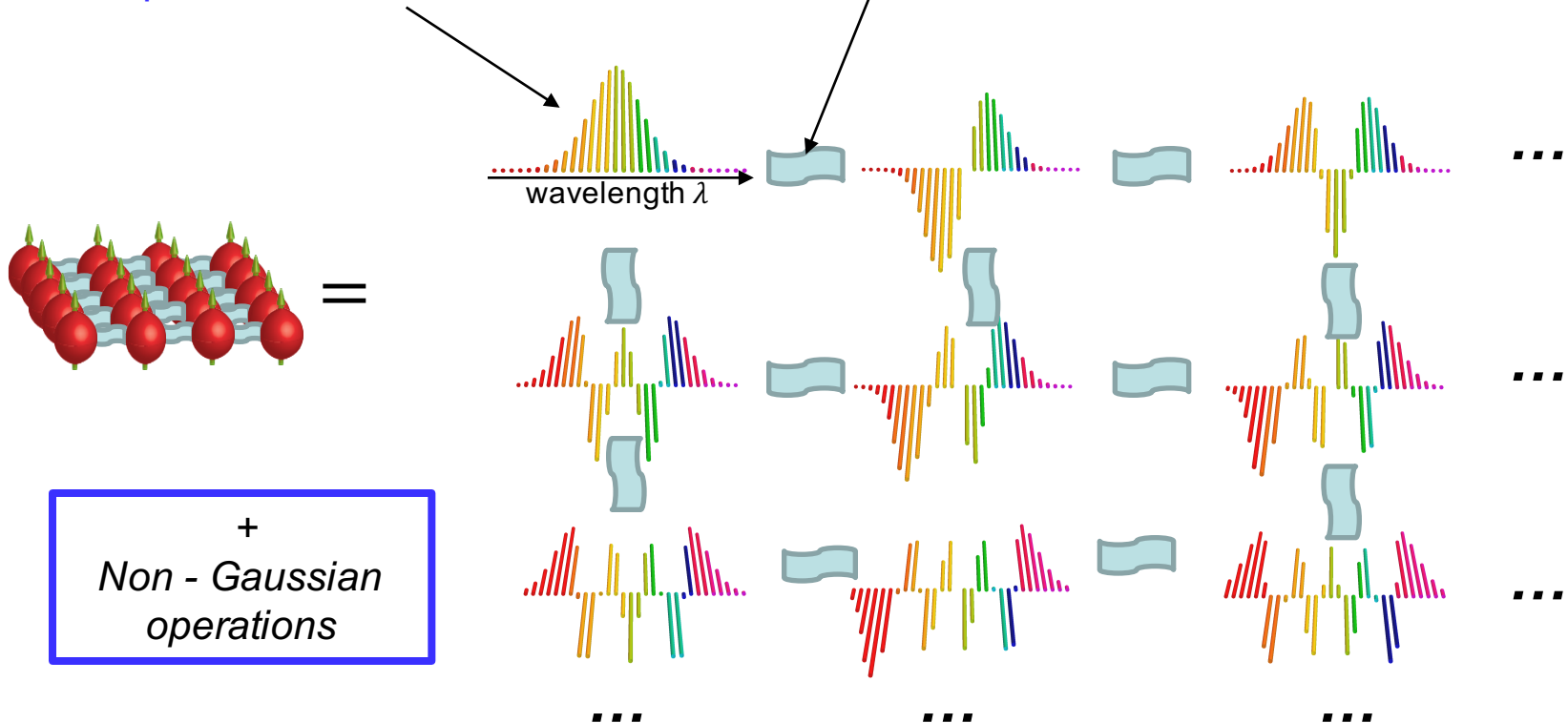


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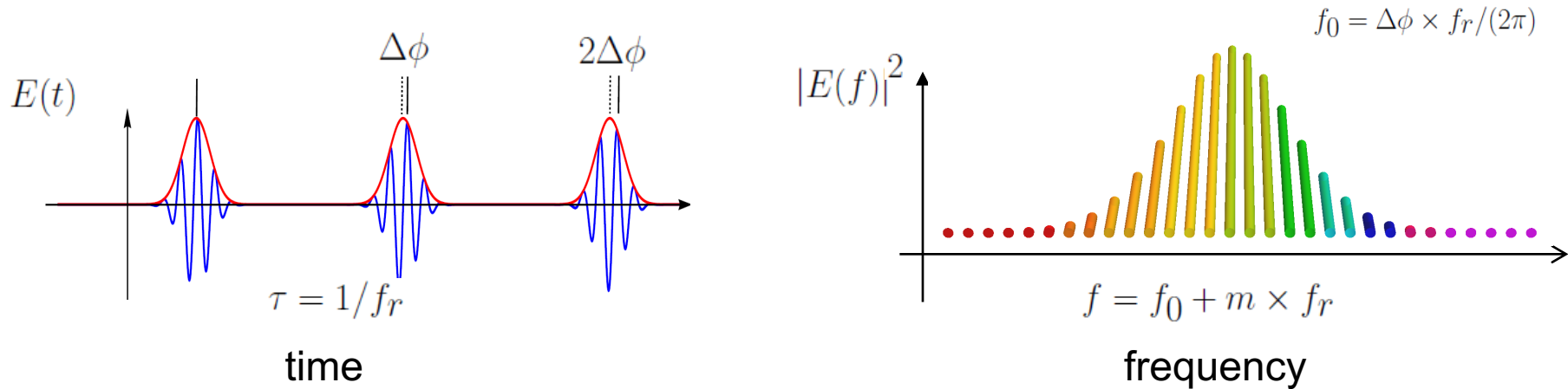
Simulation of complex quantum networks

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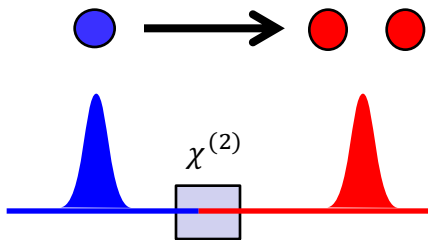
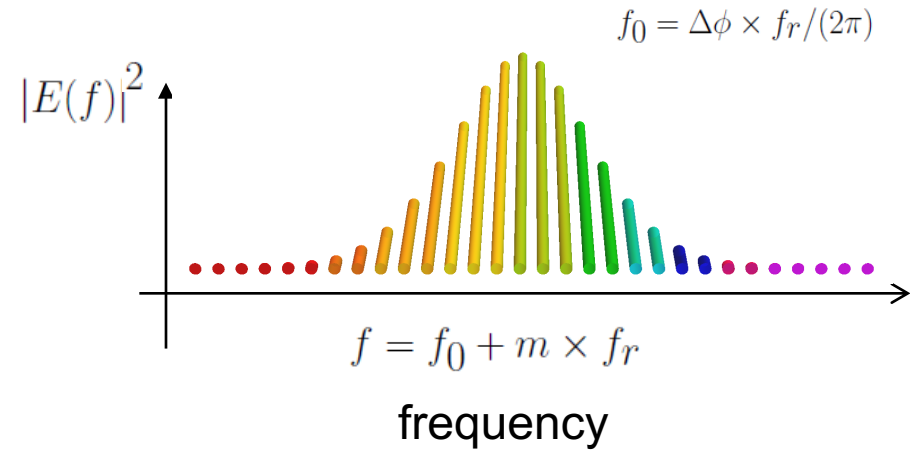
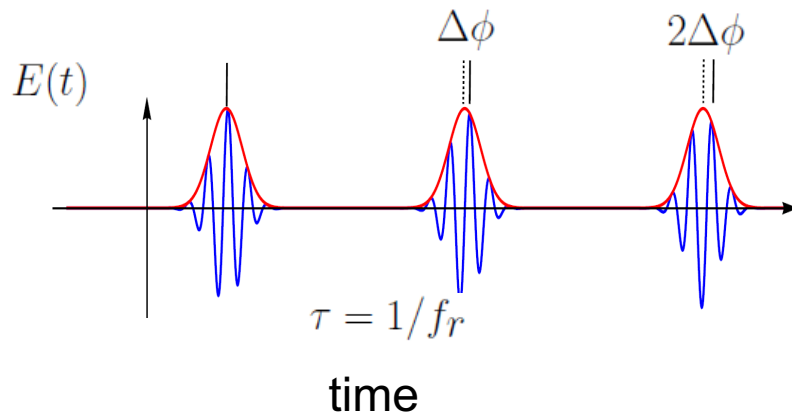
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Optical frequency comb + quantum optics



Femtosecond mode-locked laser \rightarrow Huge number of modes ($\sim 10^6$ frequencies)

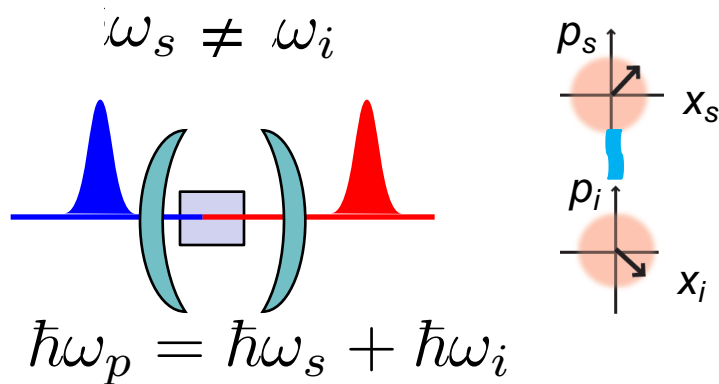
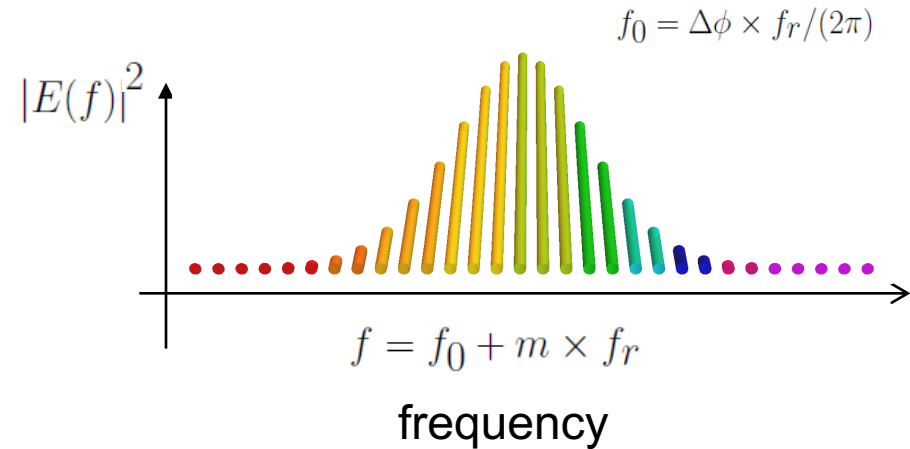
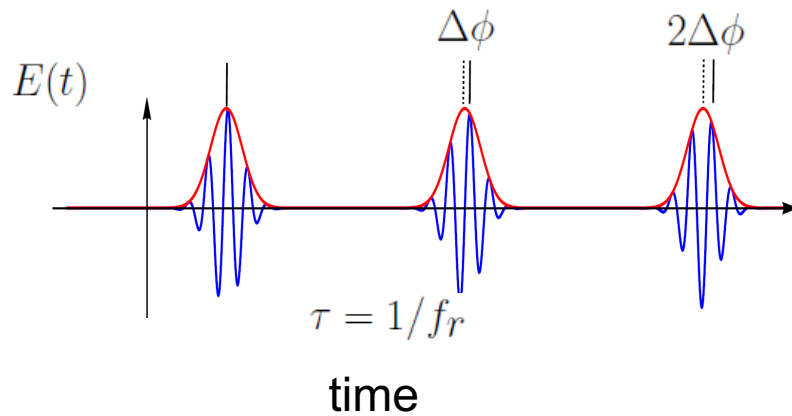
Optical frequency comb + quantum optics



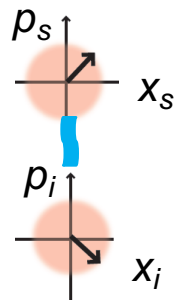
$$\hbar\omega_p = \hbar\omega_s + \hbar\omega_i$$

Parametric process

Optical frequency comb + quantum optics



Parametric process

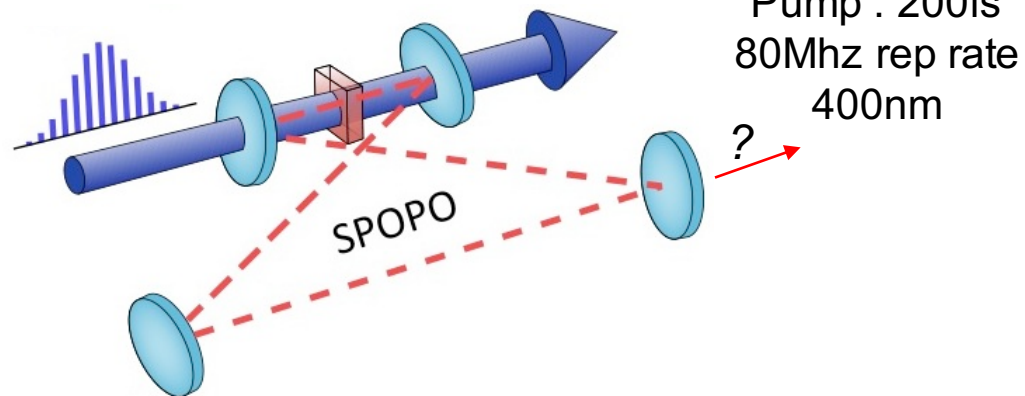


Quadrature entangled states

x_s et x_i correlated
 p_s et p_i anti-correlated

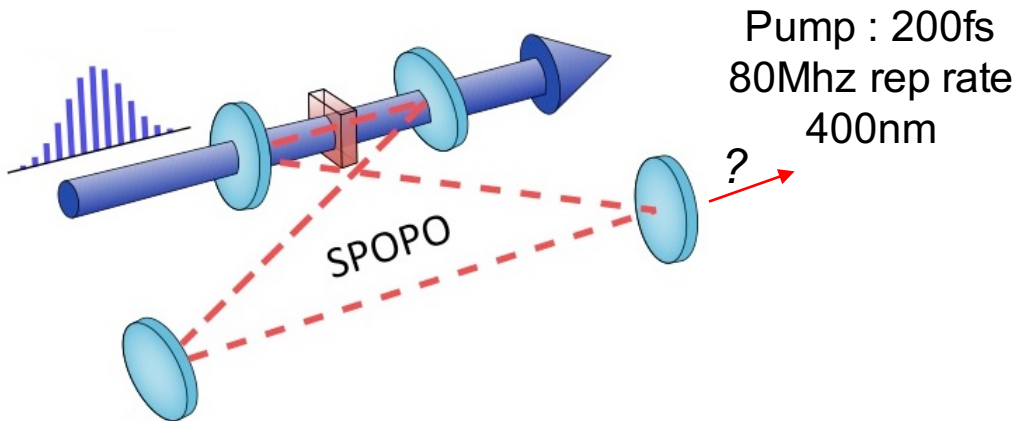
Optical frequency comb + quantum optics

Synchronously pumped OPO



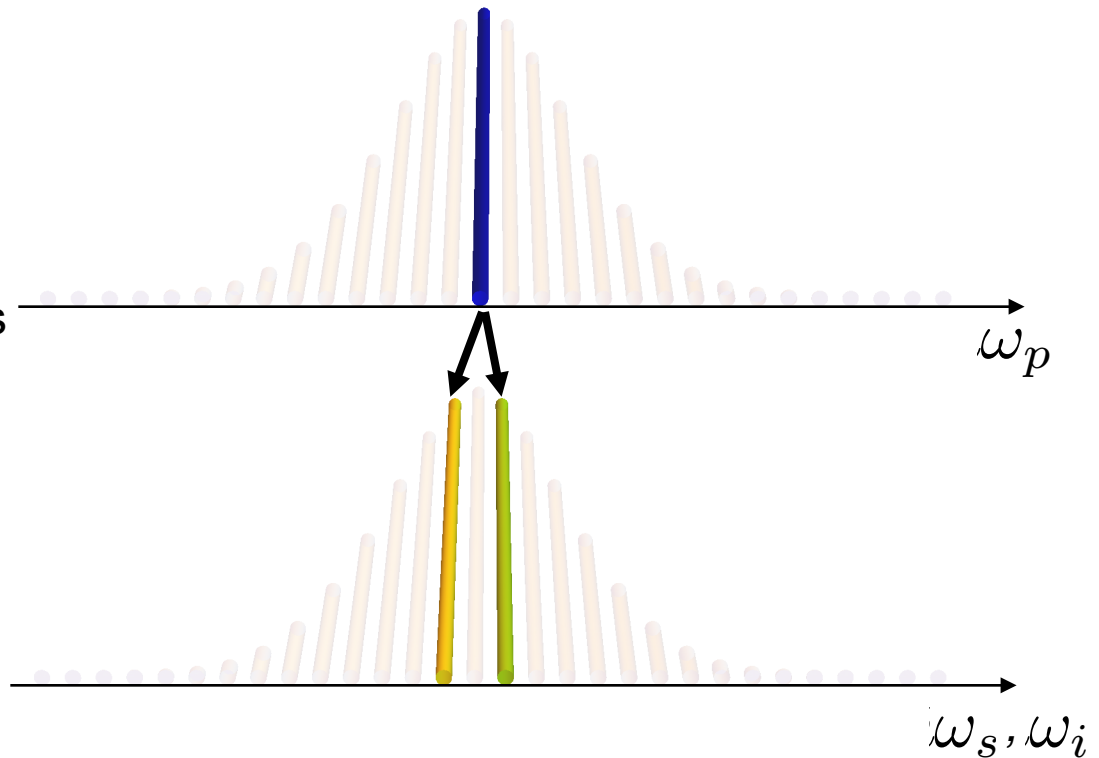
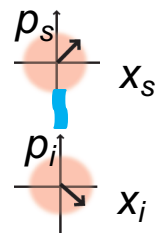
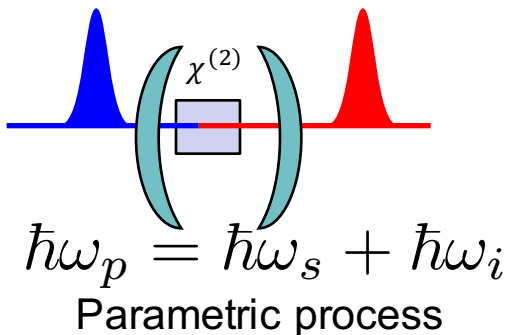
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Synchronously pumped OPO



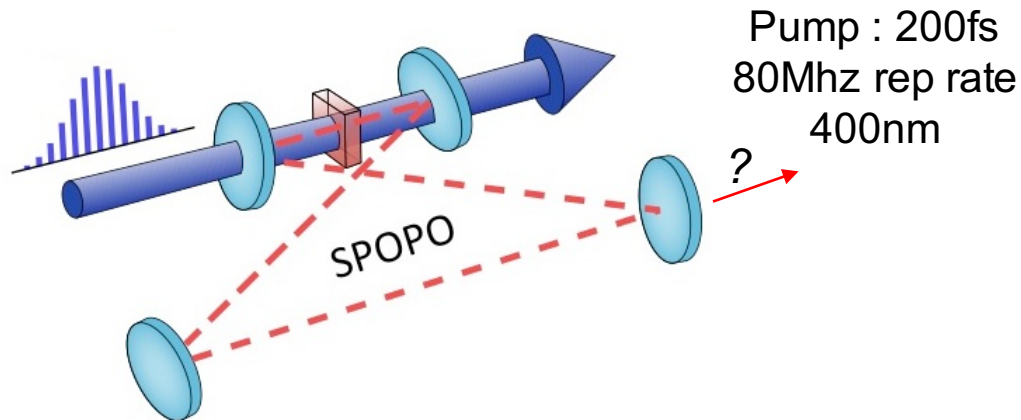
$$H = i\hbar \sum_{m,n} L_{-m,n} \hat{a}_{-m}^\dagger \hat{a}_n^\dagger$$

Symmetric Frequency Correlations (quadratures entanglement)



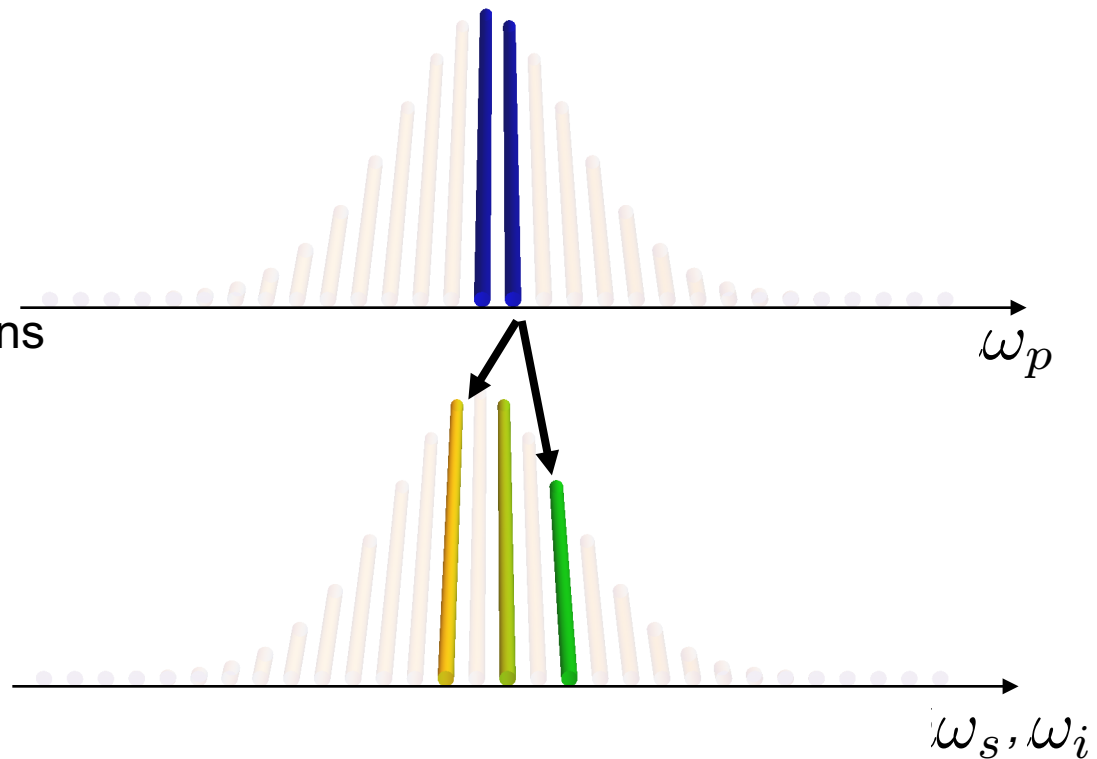
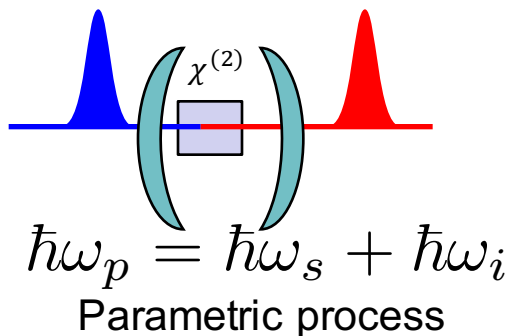
Optical frequency comb + quantum optics

Synchronously pumped OPO



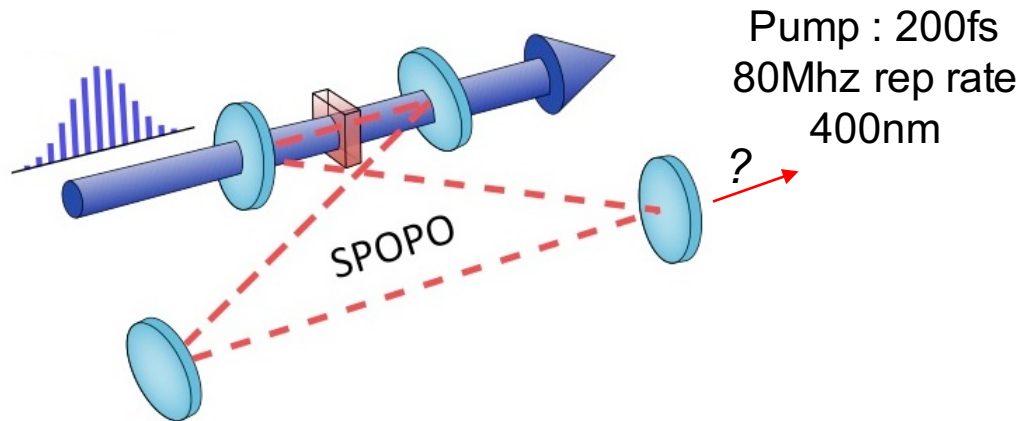
$$H = i\hbar \sum_{m,n} L_{-m,n} \hat{a}_{-m}^{\dagger} \hat{a}_n^{\dagger}$$

Asymmetric Frequency Correlations (quadratures entanglement)



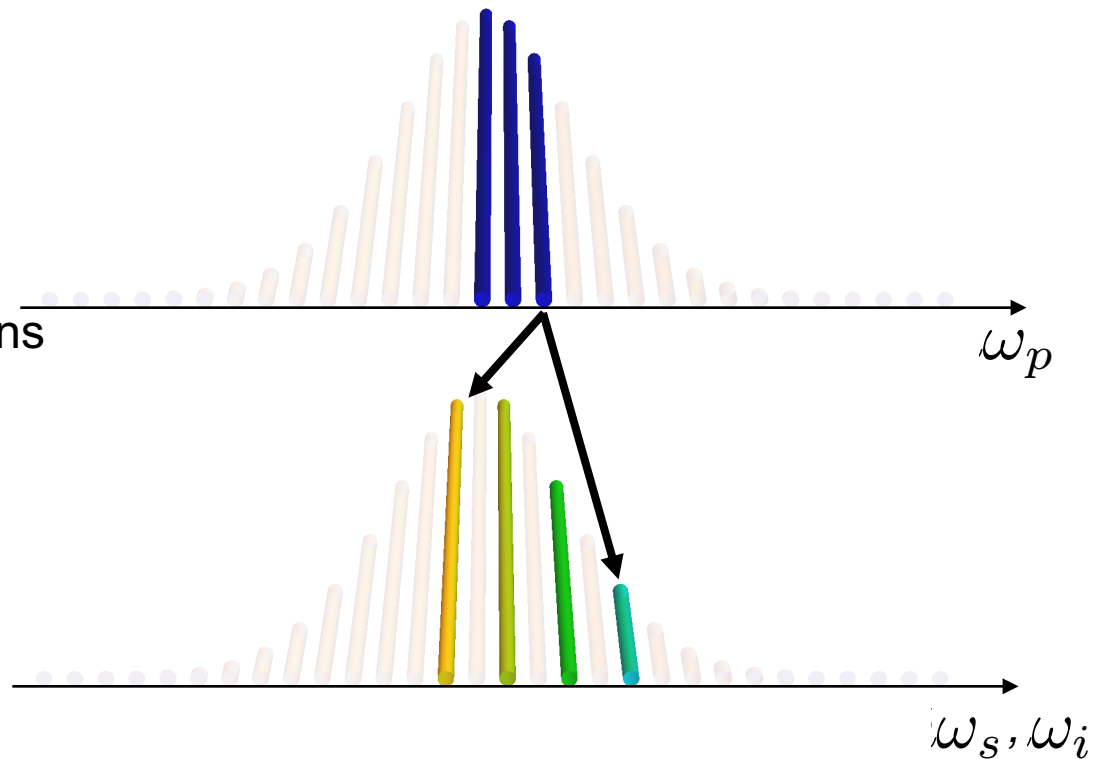
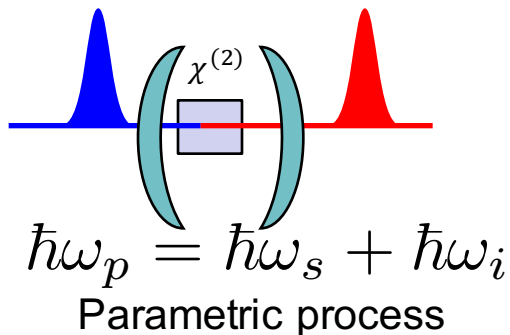
Optical frequency comb + quantum optics

Synchronously pumped OPO



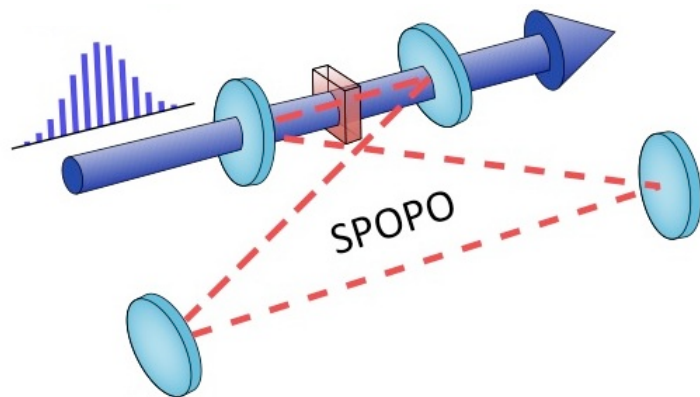
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Asymmetric Frequency Correlations (quadratures entanglement)



Optical frequency comb + quantum optics

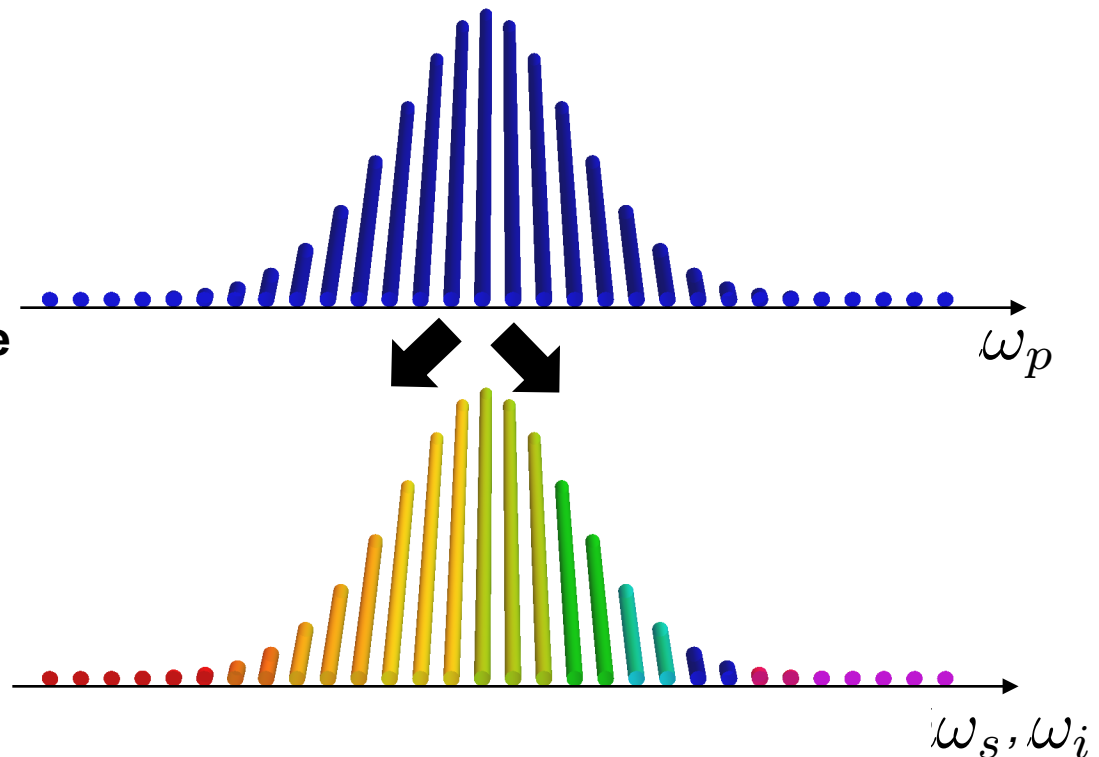
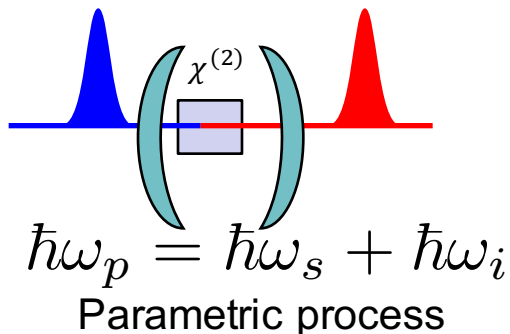
Synchronously pumped OPO



Pump : 200fs
80MHz rep rate
400nm

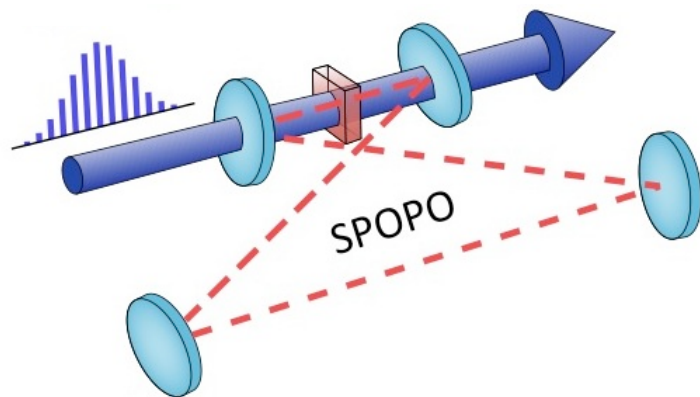
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Complex entanglement structure (quadratures entanglement)



Optical frequency comb + quantum optics

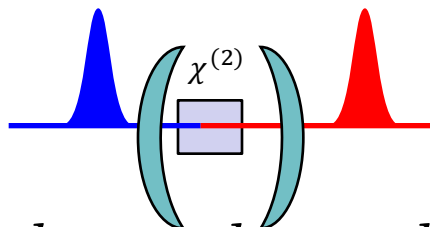
Synchronously pumped OPO



Pump : 200fs
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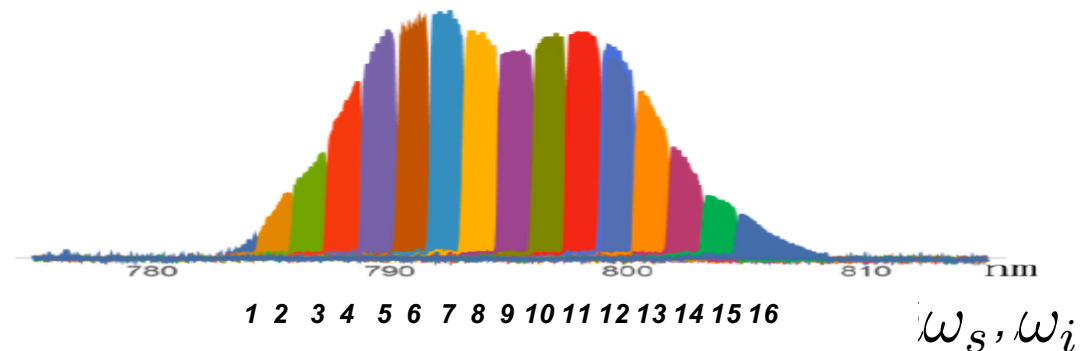
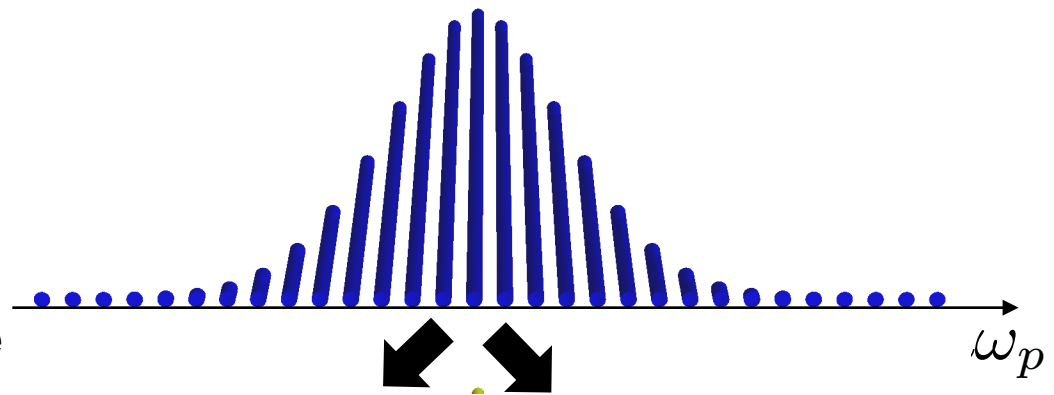
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Complex entanglement structure (quadratures entanglement)



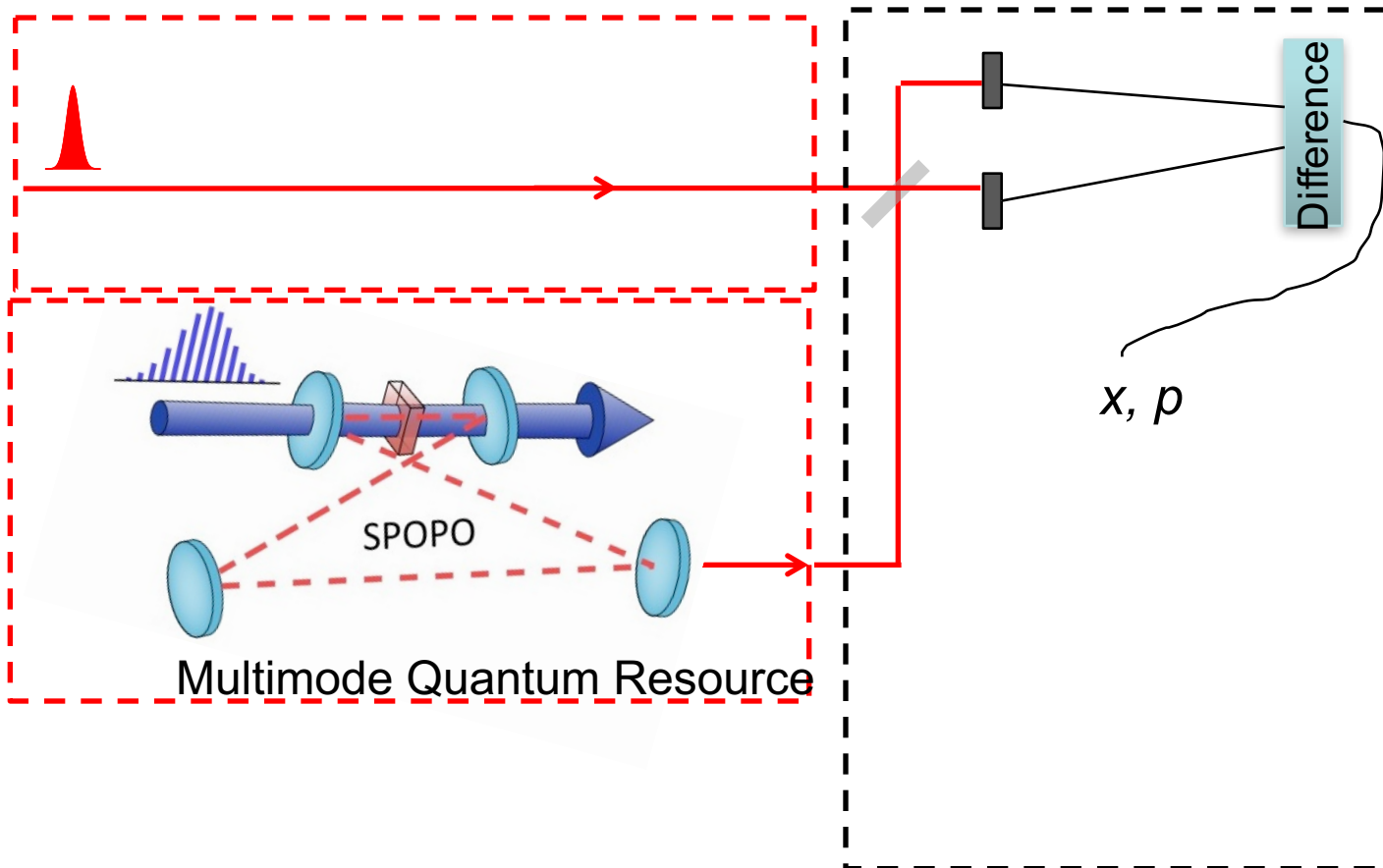
$$\hbar\omega_p = \hbar\omega_s + \hbar\omega_i$$

Parametric process

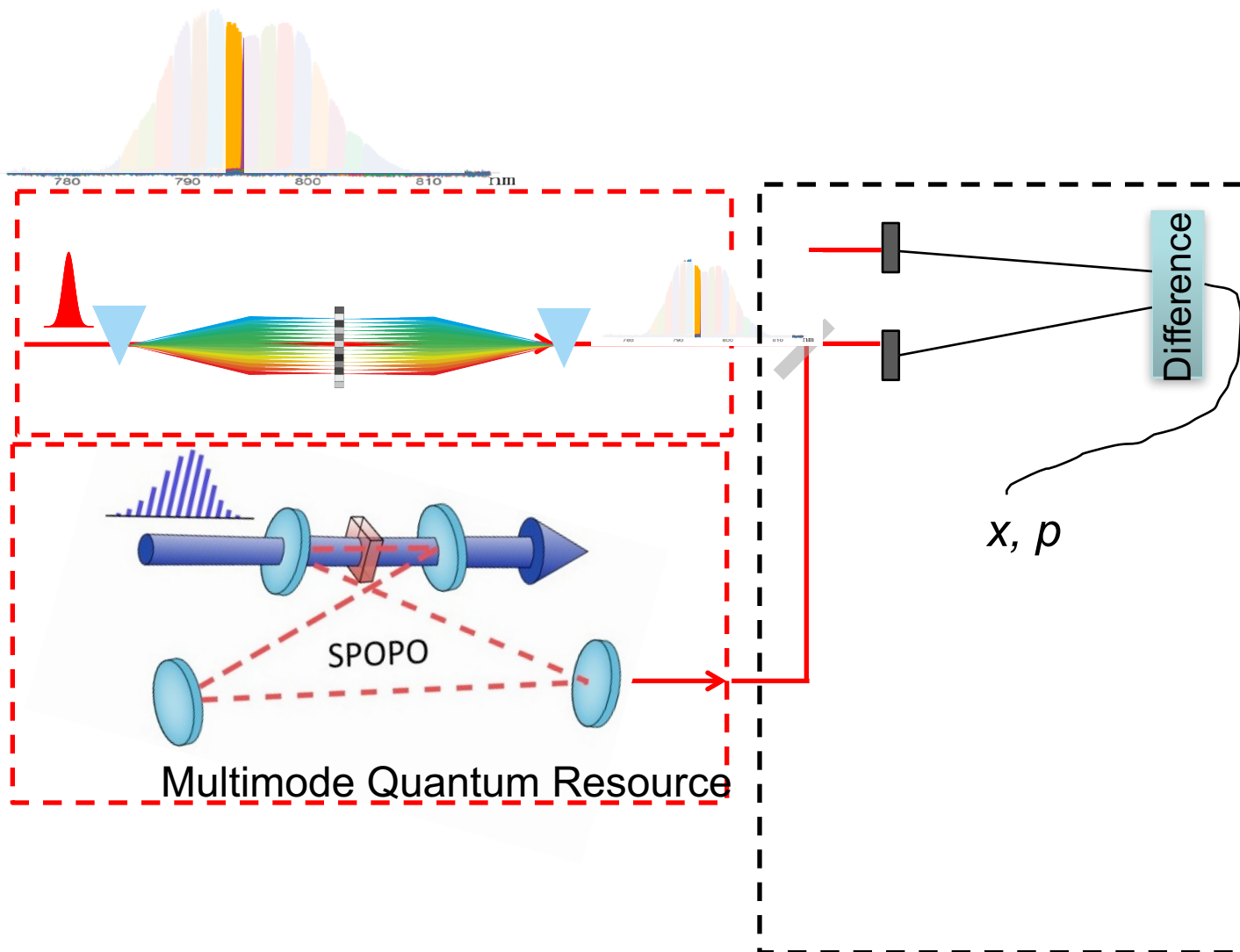


Characterization :

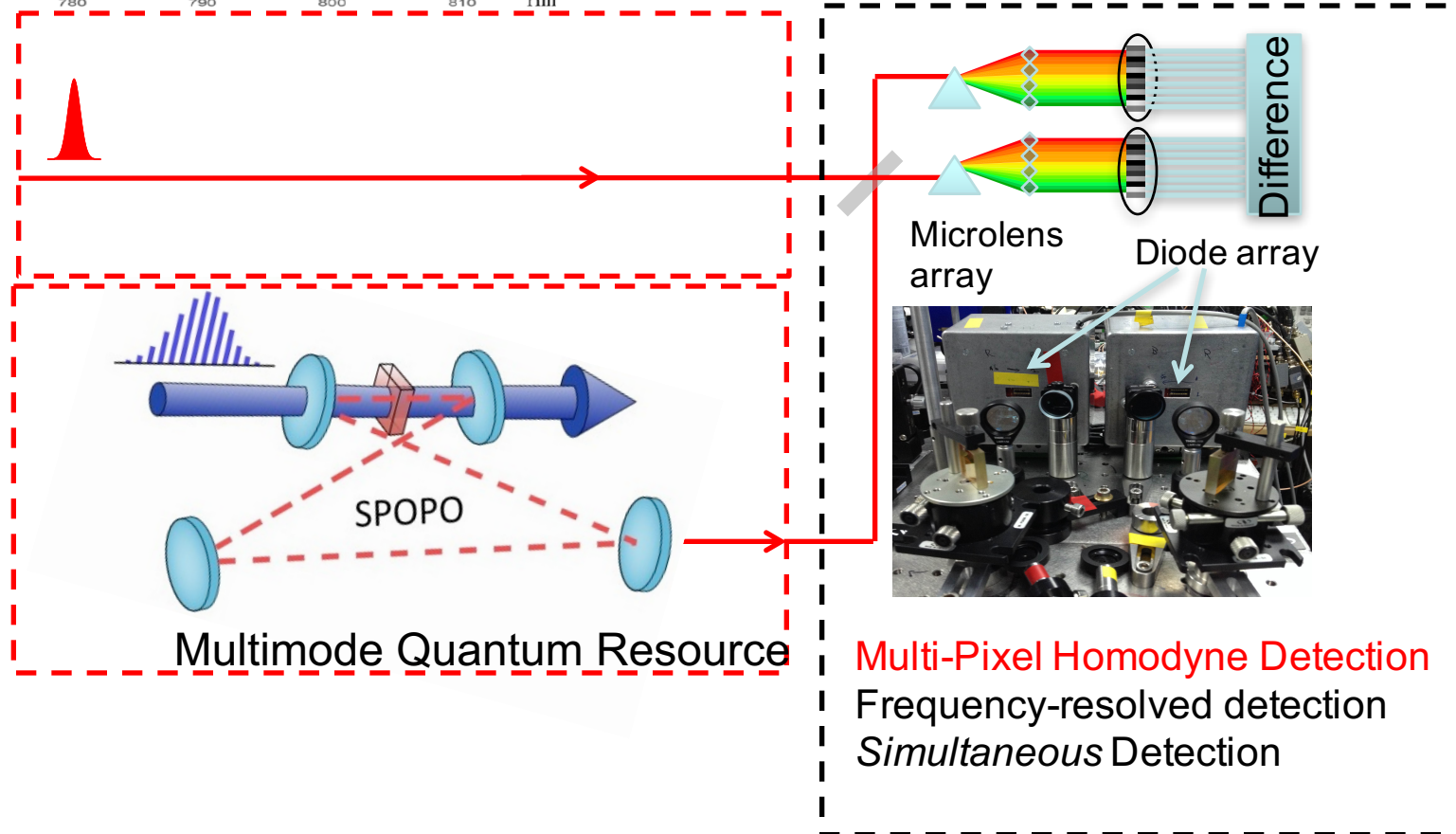
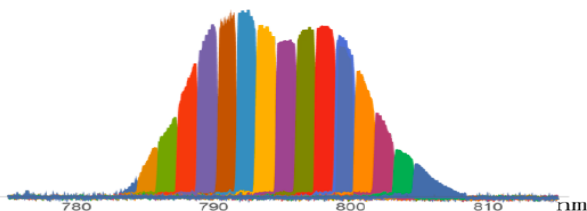
homodyne detection

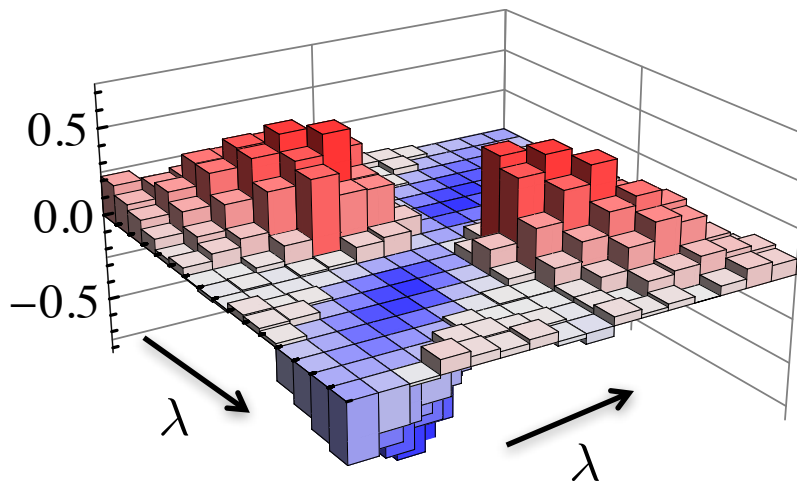
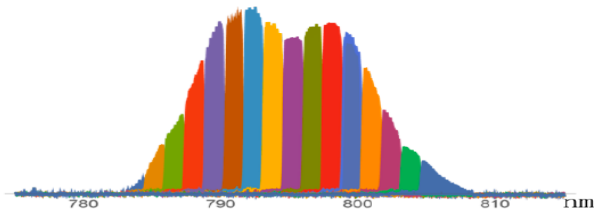


Characterization : mode-selective homodyne detection

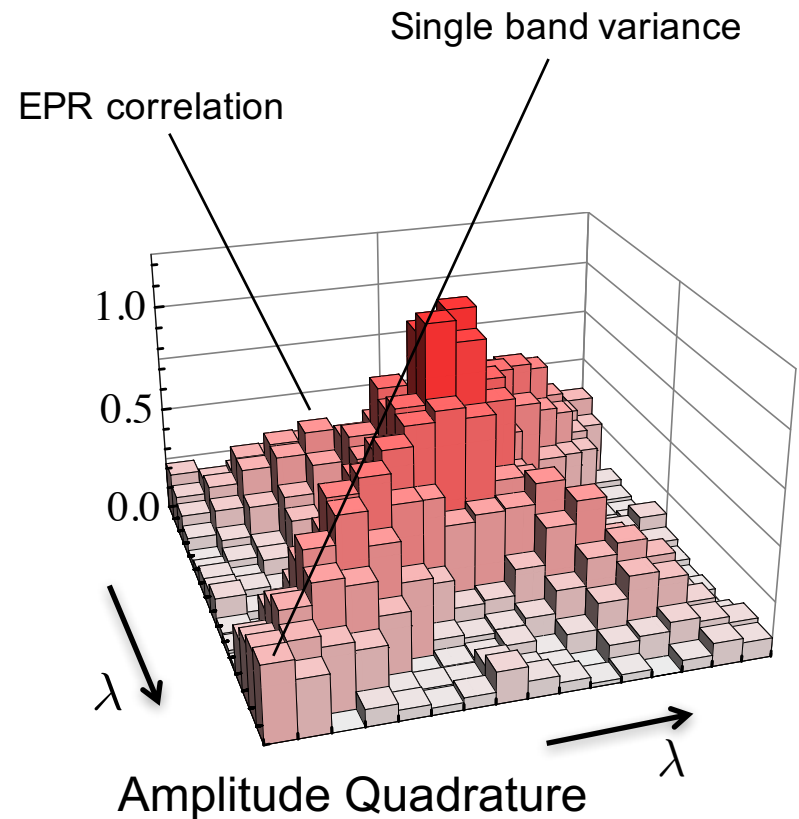


Characterization : multi-mode homodyne detection





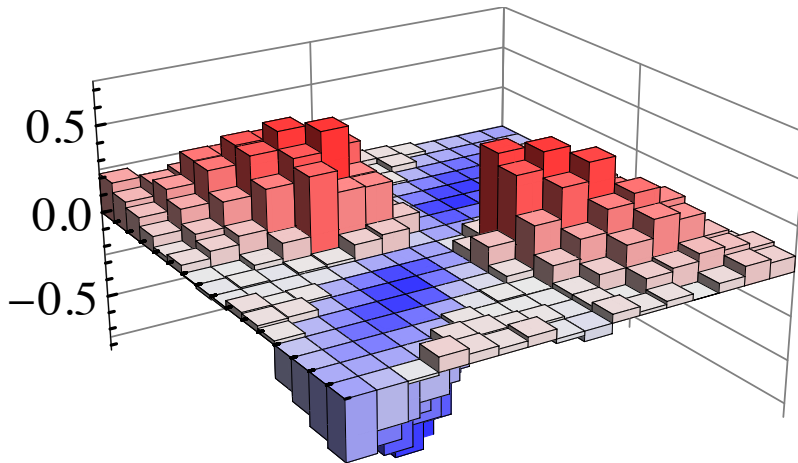
16-mode Covariance matrix
of Phase Quadrature



Diagonalization

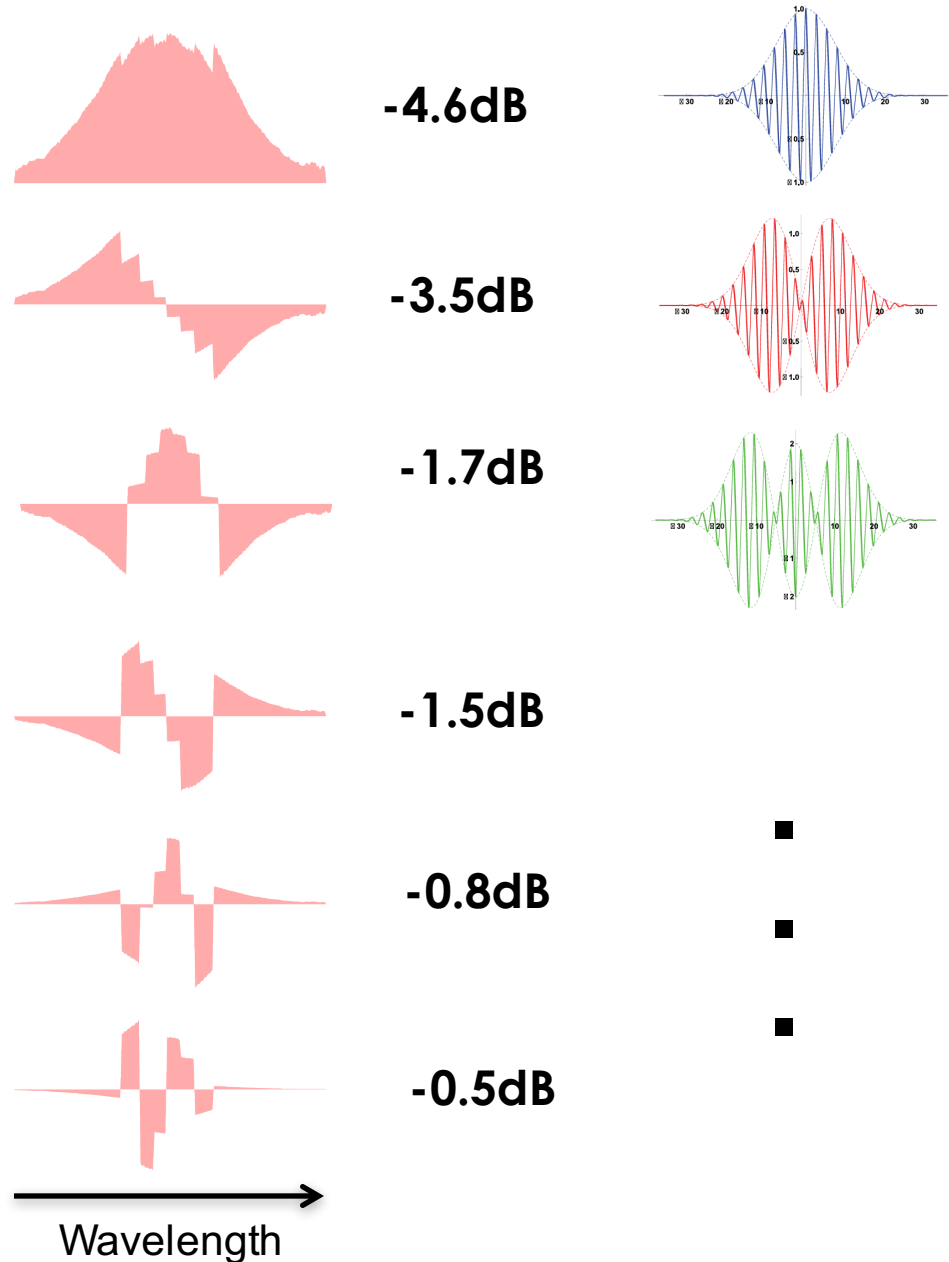


- Independent Squeezers
- Unique temporal / spectral pulse shapes



J. Roslund, R. Medeiros, S. Jiang,
C. Fabre and N. Treps,
Nature Photonics 8, 109–112 (2014)

Spectrum · Noise level · Pulse shape



Diagonalization



- Independent Squeezers
- Unique temporal / spectral pulse shapes

It corresponds to

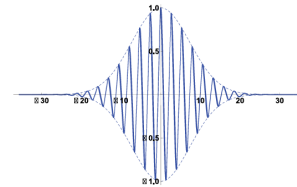
$$H = i\hbar \sum_{m,n} L_{-m,n} \hat{a}_{-m}^{\dagger} \hat{a}_n^{\dagger}$$



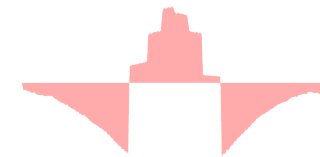
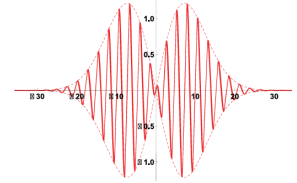
$$H = i\hbar \sum_k \Lambda_k \hat{S}_k^{\dagger 2}$$



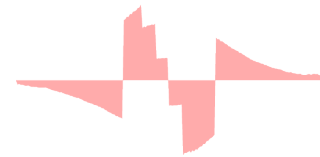
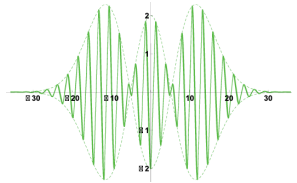
-4.6dB



-3.5dB



-1.7dB



-1.5dB



-0.8dB



-0.5dB



Wavelength

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QuProCSII 07 April 2017

Bloch-Messiah decomposition

PHYSICAL REVIEW A **71**, 055801 (2005)

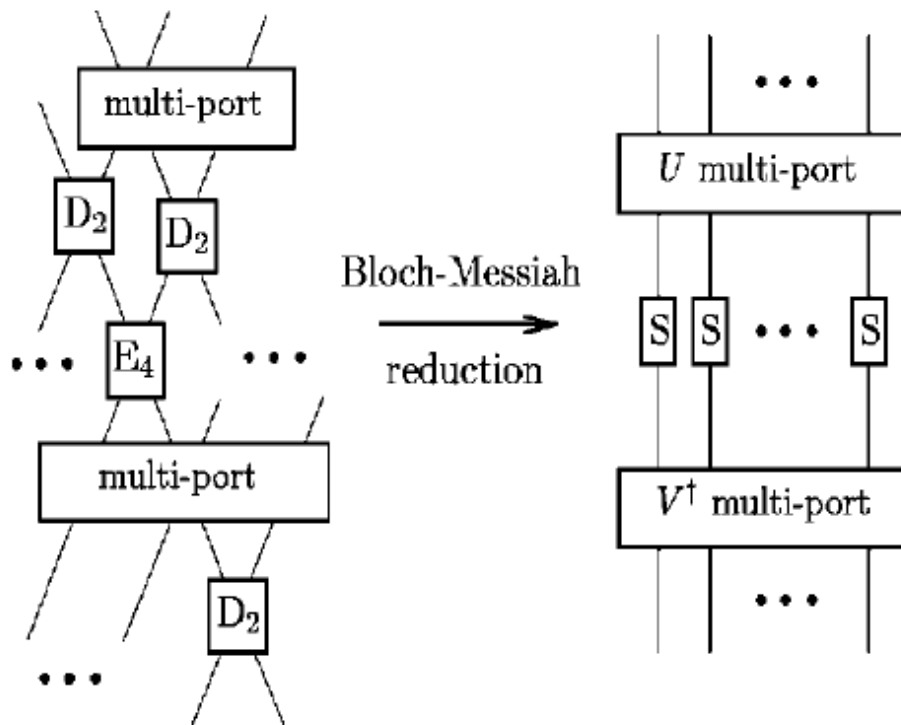
Squeezing as an irreducible resource

Samuel L. Braunstein

Computer Science, University of York, York YO10 5DD, United Kingdom

(Received 6 March 2005; published 31 May 2005)

multimode quantum
Gaussian resource =
Squeezed modes +
basis change



$$\vec{a} = (a_1, a_2, \dots, a_N).$$

collection of input modes

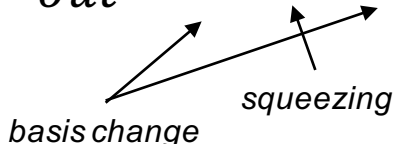
$$\vec{a}_{out} = V^\dagger S U \vec{a}$$

basis change

squeezing

Bloch-Messiah decomposition

$$\vec{a}_{out} = V^\dagger S U \vec{a}$$



basis change *squeezing*

if \vec{a} collection of vacuum states

$$\vec{a}_{out} = V^\dagger S \vec{a}$$

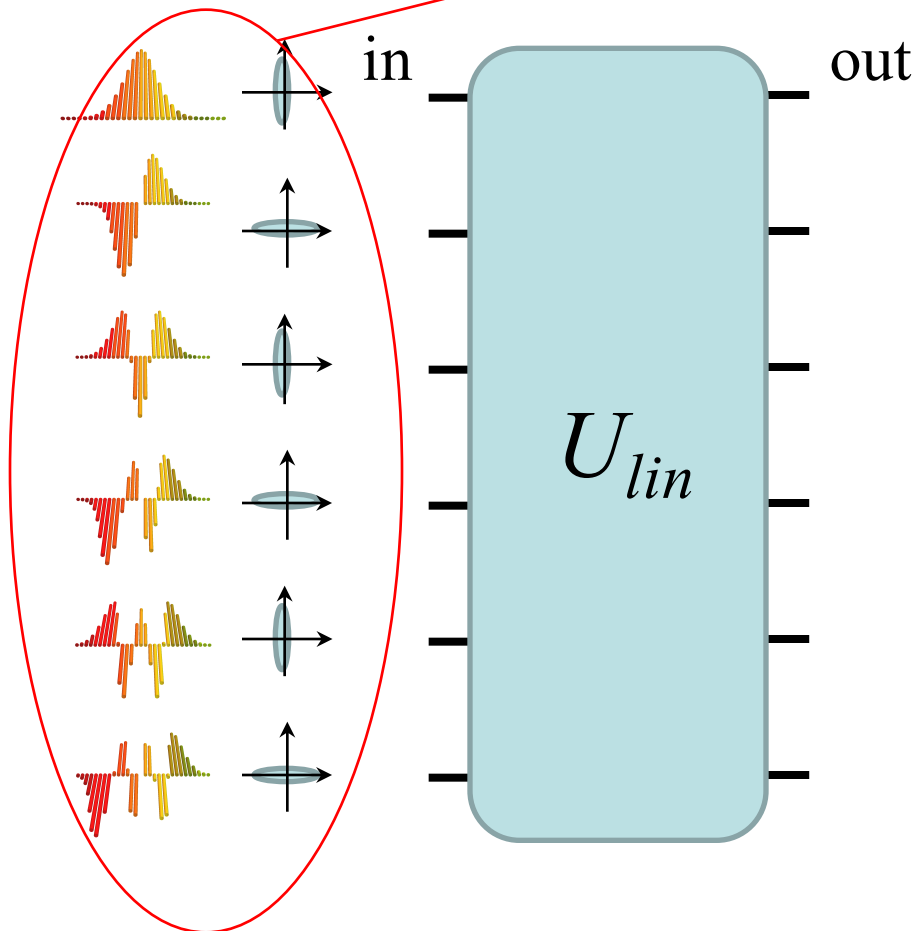
$$\vec{a}_{out} = V^\dagger S U \vec{a}$$

\nearrow basis change
 \nearrow squeezing

if \vec{a} collection of vacuum states

$$\vec{a}_{out} = V^\dagger S \vec{a}$$

Supermodes



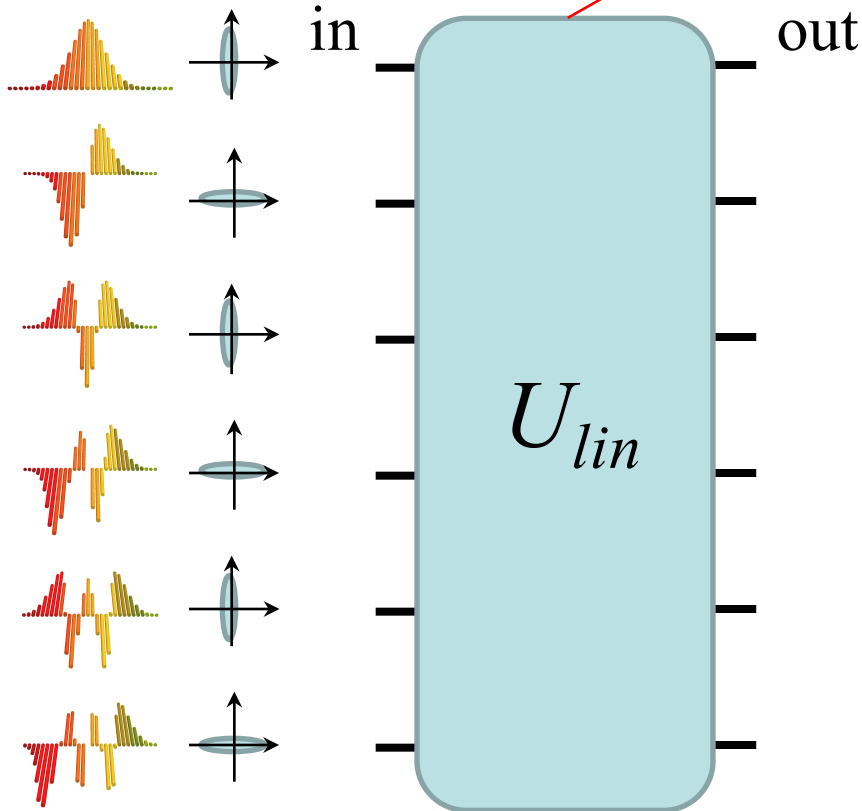
$$\vec{a}_{out} = V^\dagger S U \vec{a}$$

\nearrow basis change
 \nearrow squeezing

if \vec{a} collection of vacuum states

$$\vec{a}_{out} = V^\dagger S \vec{a}$$

Supermodes



$$\vec{a}_{out} = V^\dagger S U \vec{a}$$

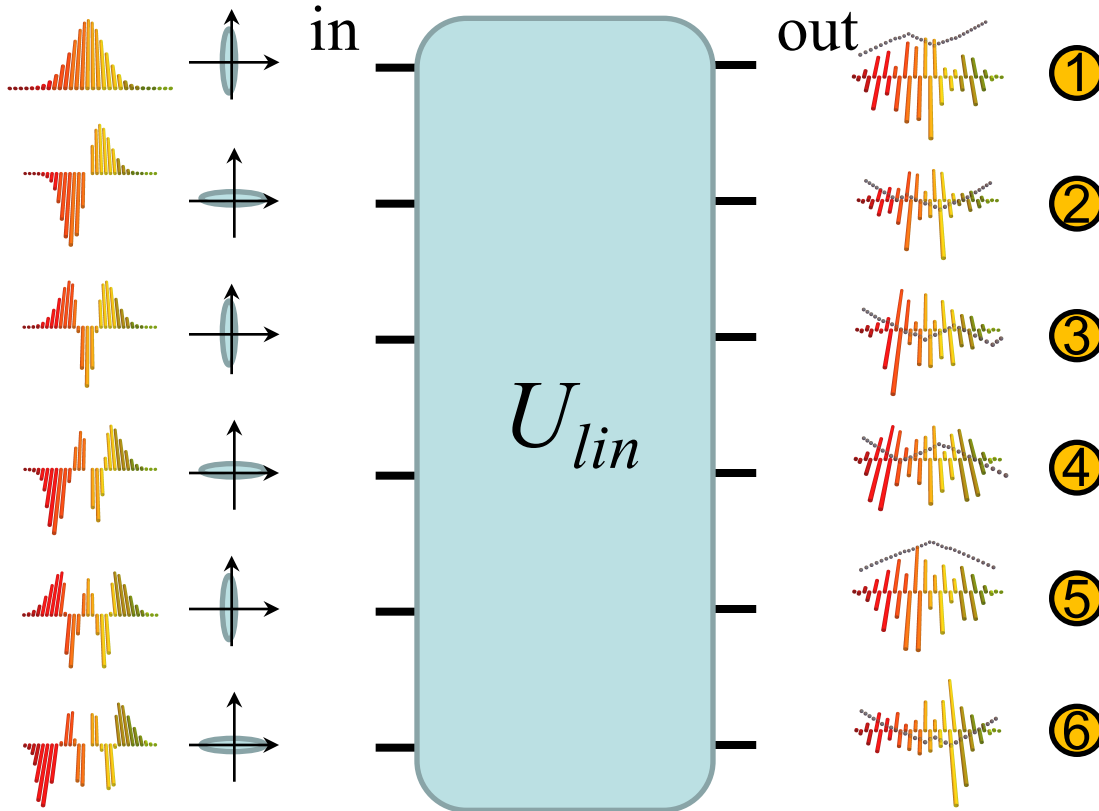
\nearrow basis change
 \nwarrow squeezing

if \vec{a} collection of vacuum states

$$\vec{a}_{out} = V^\dagger S \vec{a}$$

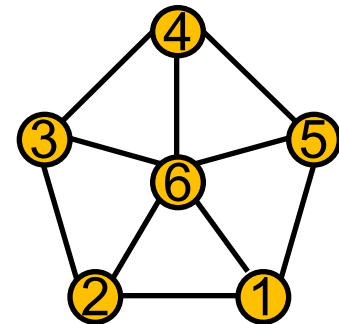
\nwarrow Supermodes

Cluster states



$$\Delta \left(\hat{p}_n - \sum_b \hat{x}_b \right) \rightarrow 0$$

nullifier



if \vec{a} collection of vacuum states

$$\vec{a}_{out} = V^\dagger \underbrace{S}_{\text{Supermodes}} \vec{a}$$

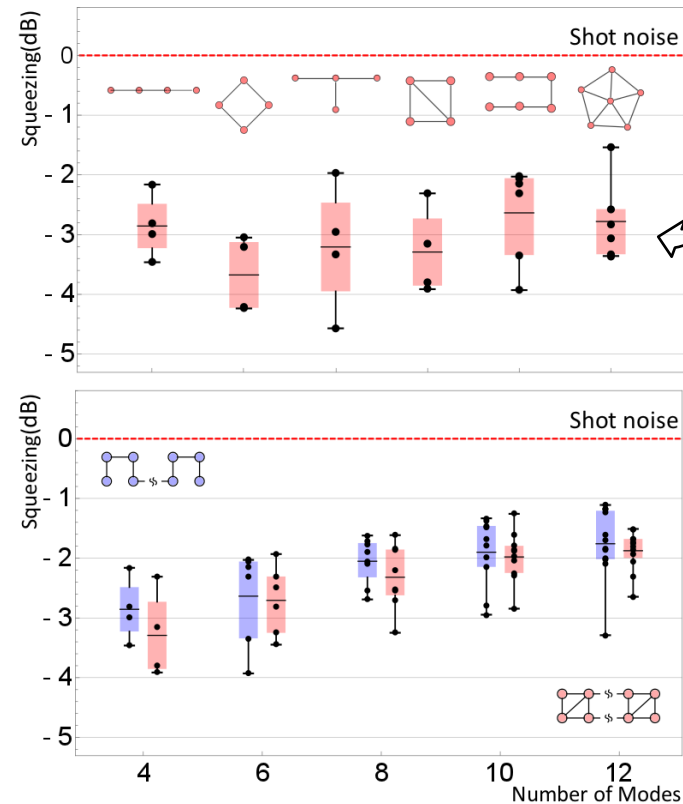
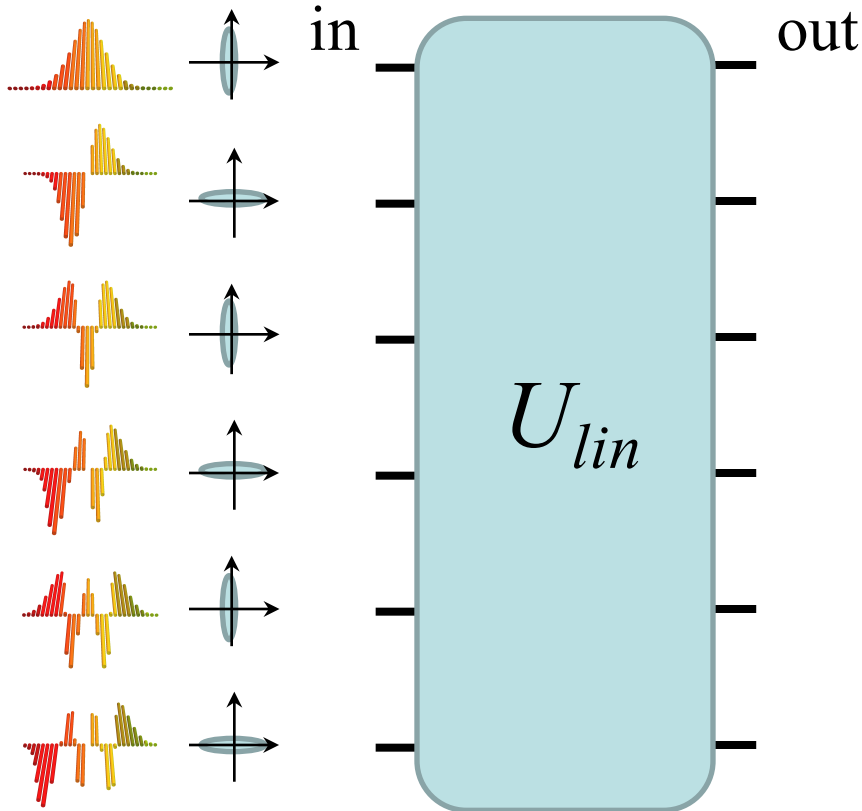
Cluster states

$$\Delta \left(\hat{p}_n - \sum_b \hat{x}_b \right) \rightarrow 0$$

nullifiers

$$\vec{a}_{out} = V^\dagger \underbrace{SU}_{\text{basis change}} \vec{a}$$

squeezing



$$\vec{a}_{out} = V^\dagger S U \vec{a}$$

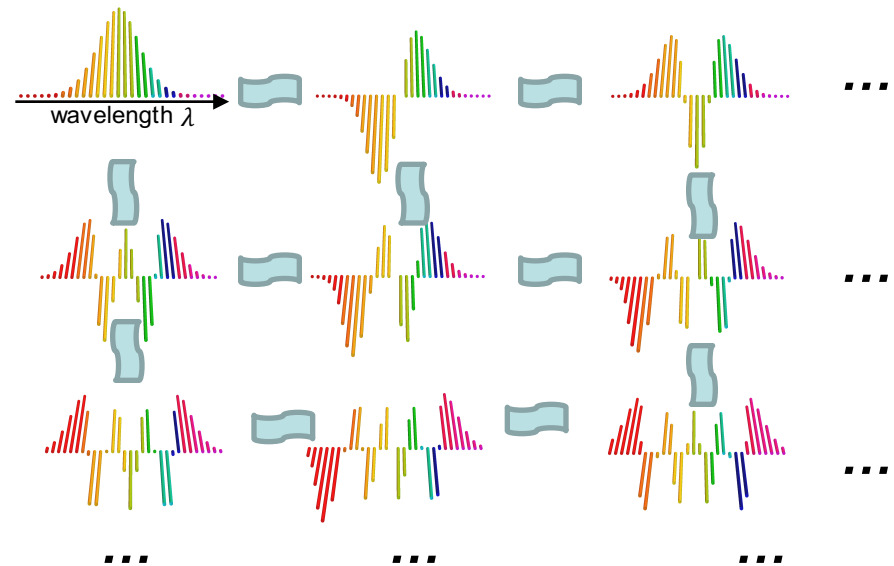
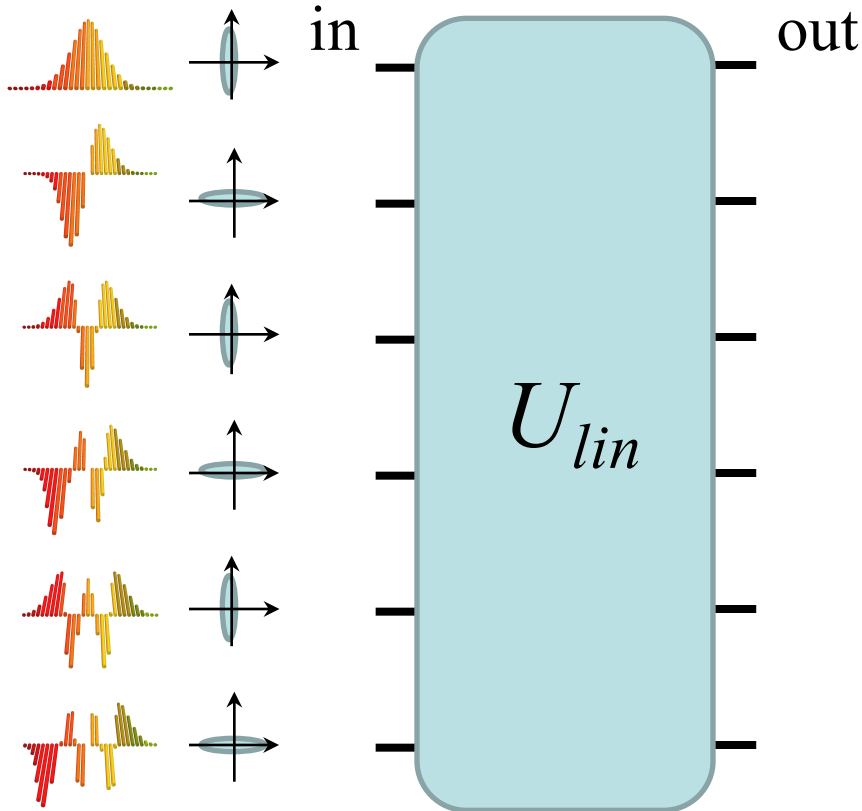
\nearrow basis change
 \nwarrow squeezing

if \vec{a} collection of vacuum states

$$\vec{a}_{out} = V^\dagger S \vec{a}$$

\nwarrow Supermodes

Cluster states



if \vec{a} collection of vacuum states

$$\vec{a}_{out} = V^\dagger S \vec{a}$$

Supermodes

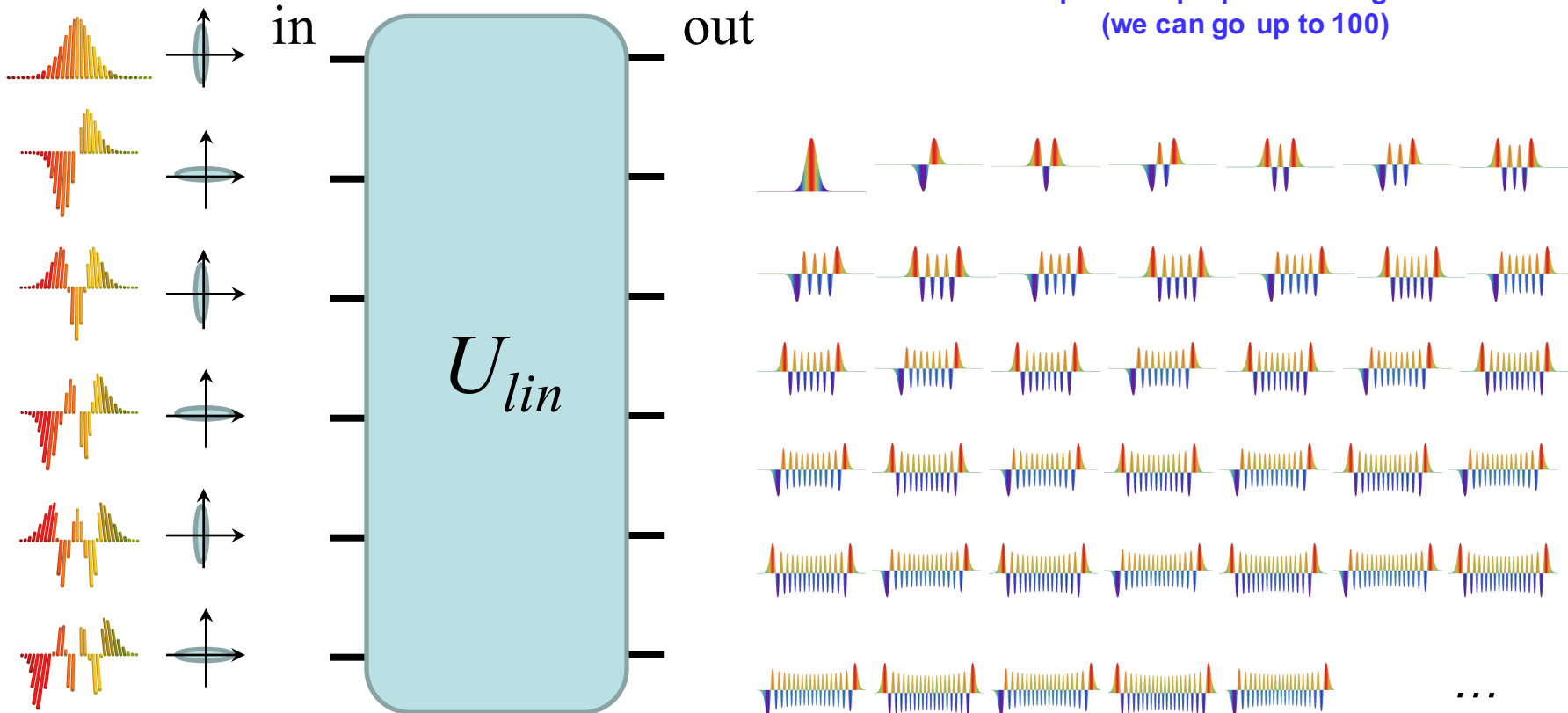
Cluster states

$$\vec{a}_{out} = V^\dagger S U \vec{a}$$

basis change squeezing

Current setup

At least 40 Hermite-Gauss modes
with quantum properties are generated
(we can go up to 100)



$$\vec{a}_{out} = V^\dagger S U \vec{a}$$

\nearrow basis change
 \nearrow squeezing

if \vec{a} collection of vacuum states

$$\vec{a}_{out} = V^\dagger S \vec{a}$$

\nwarrow Supermodes

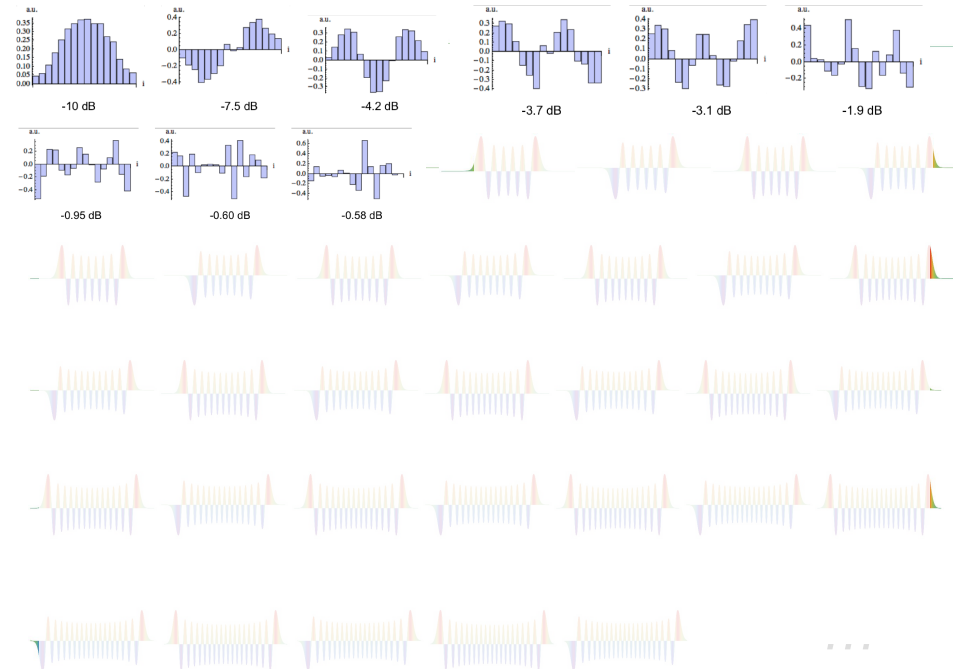
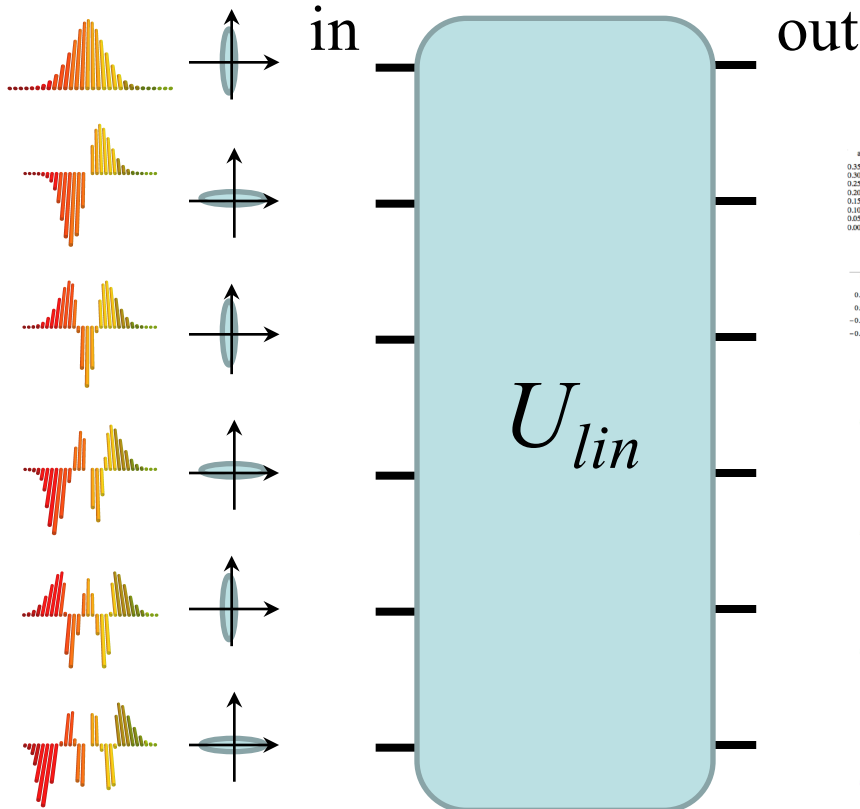
Cluster states

Current setup

**At least 40 Hermite-Gauss modes
with quantum properties are generated**

But

We can address them only up to $\sim 16..$



$$\vec{a}_{out} = V^\dagger S U \vec{a}$$

basis change squeezing

if \vec{a} collection of vacuum states

$$\vec{a}_{out} = V^\dagger S \vec{a}$$

Supermodes

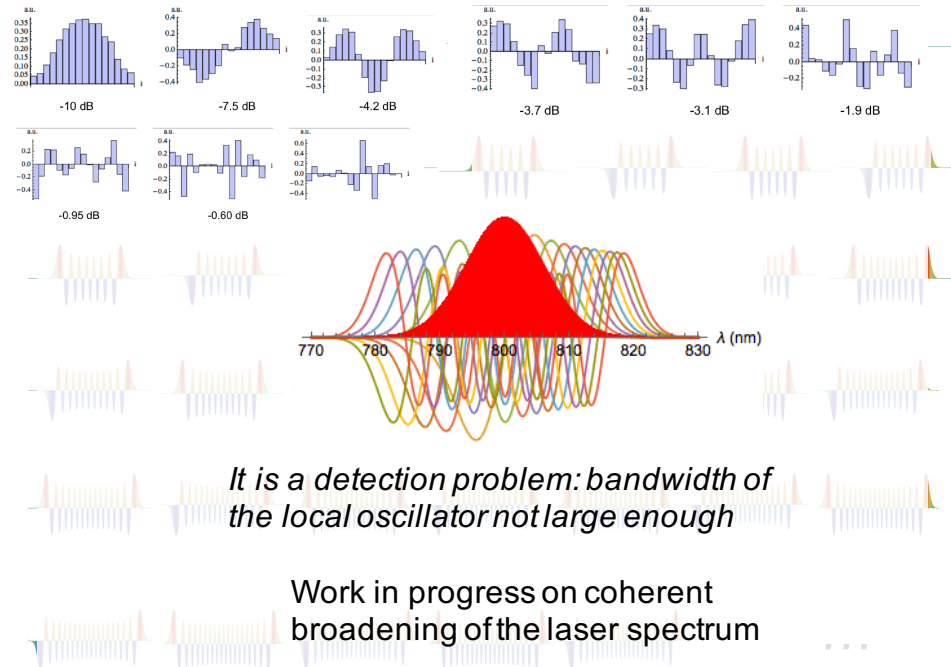
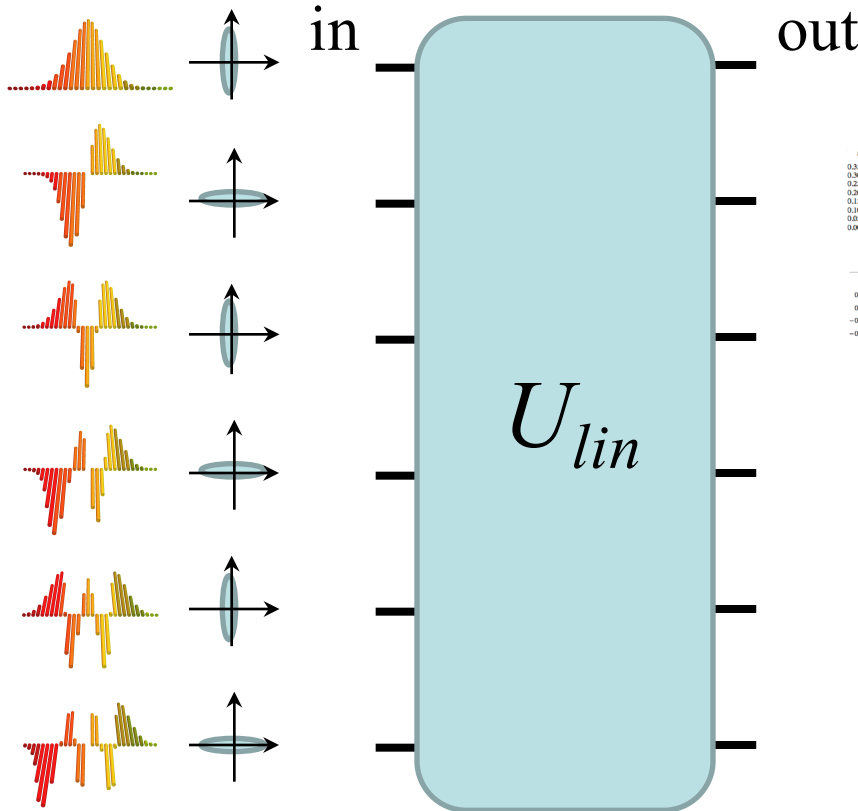
Cluster states

Current setup

**At least 40 Hermite-Gauss modes
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But

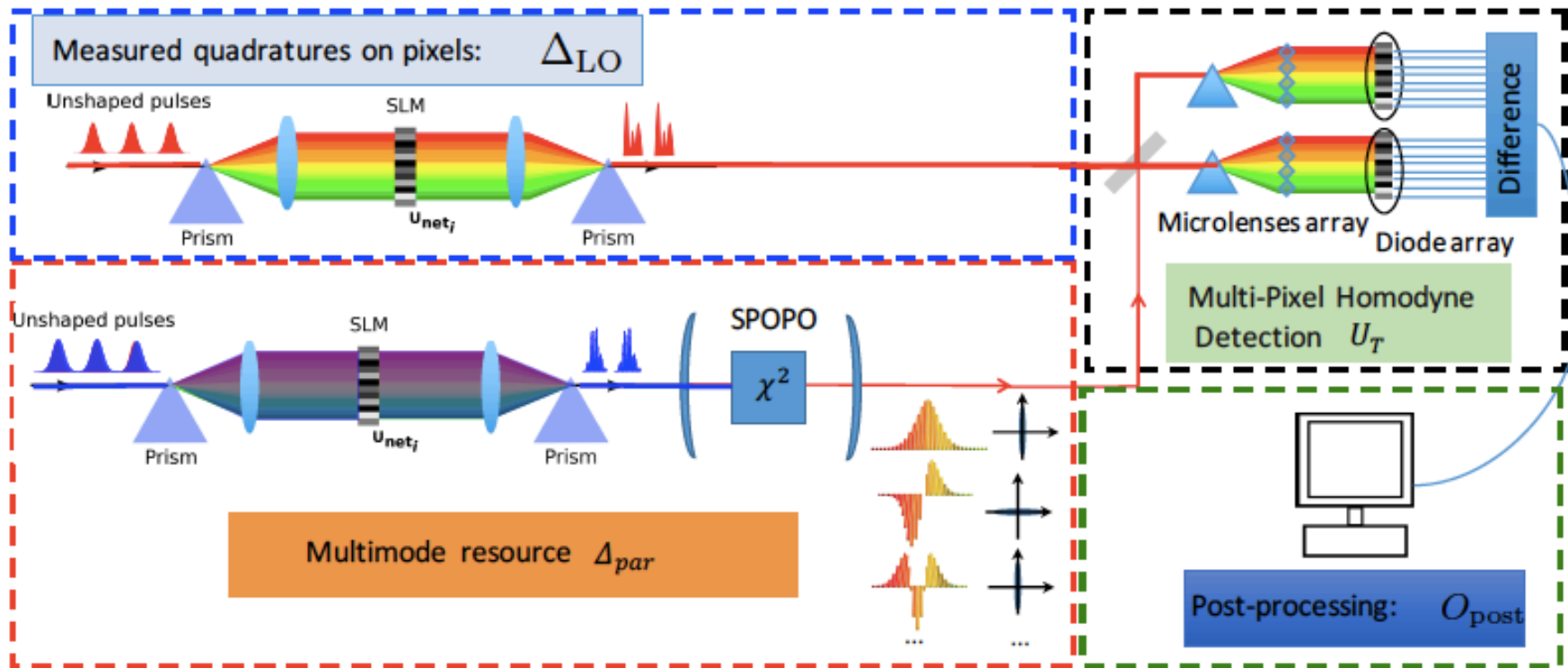
We can address them only up to ~ 16..



*It is a detection problem: bandwidth of
the local oscillator not large enough*

Work in progress on coherent
broadening of the laser spectrum

...

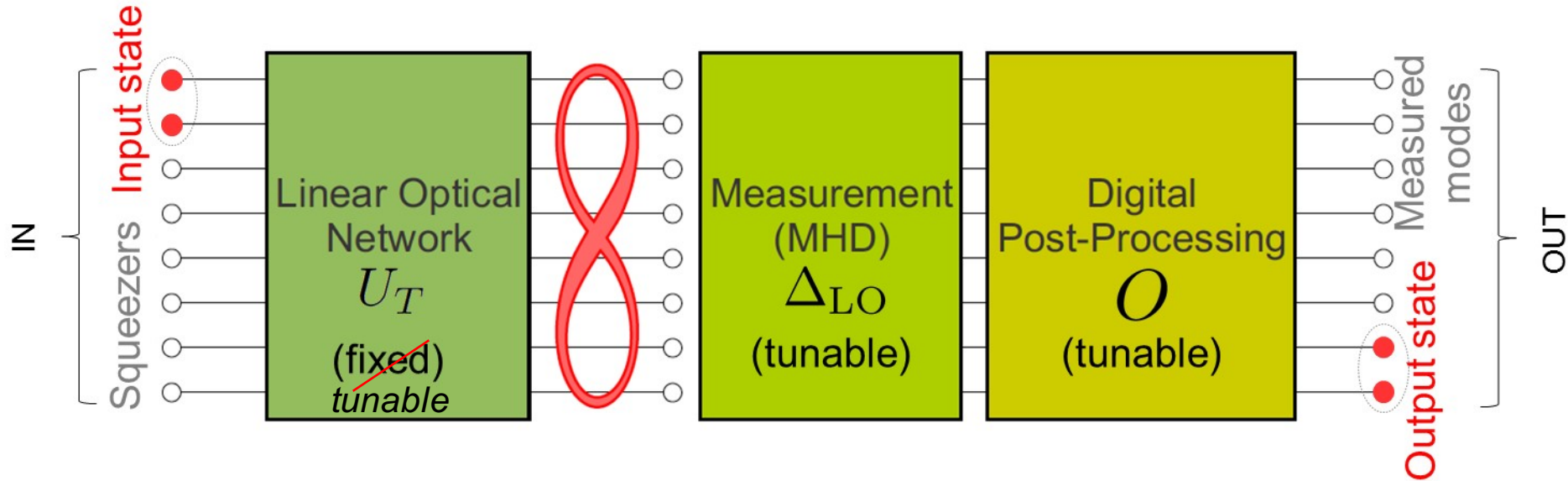


$$\vec{a}_{out} = V^\dagger S \vec{a} = O_{post} \Delta_{LO} G \vec{a} = O_{post} \Delta_{LO} U_T \Delta_{par} \vec{a}$$

$G = U_T \Delta_{par}$ parametric process, Δ_{par} squeezing (basis of supermodes), U_T basis change to pixels

Δ_{LO} phase of LO components

O_{post} final basis change given by the post-processing



System is solved to perform the proper evolution

$$\begin{pmatrix} \vec{x}^{\text{out}} \\ \vec{p}^{\text{out}} \end{pmatrix} = \begin{pmatrix} A & B \\ C & D \end{pmatrix} \begin{pmatrix} \vec{x}^{\text{in}} \\ \vec{p}^{\text{in}} \end{pmatrix} + \begin{pmatrix} \vec{\delta}_x \\ \vec{\delta}_p \end{pmatrix} + \begin{pmatrix} \vec{\eta}_x \\ \vec{\eta}_p \end{pmatrix}$$

G. Ferrini *et al*, New J Phys 15, 093015 (2013)

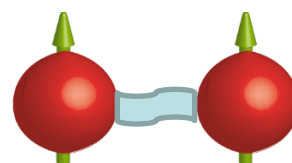
G. Ferrini *et al.*, PRA **94**, 062332 (2016)

Excess noise due to
finite squeezing

Correction factors
from post processing

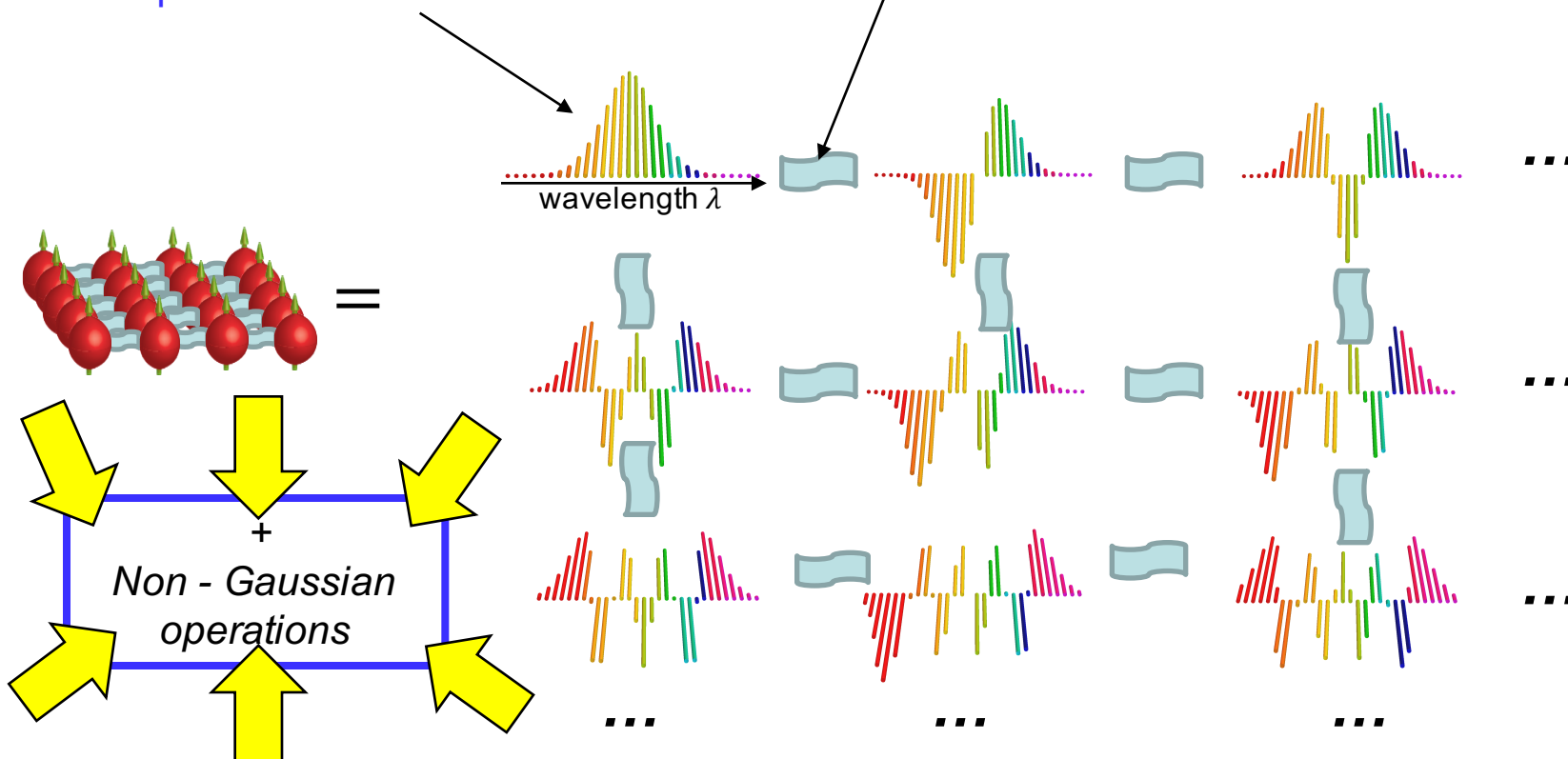


+

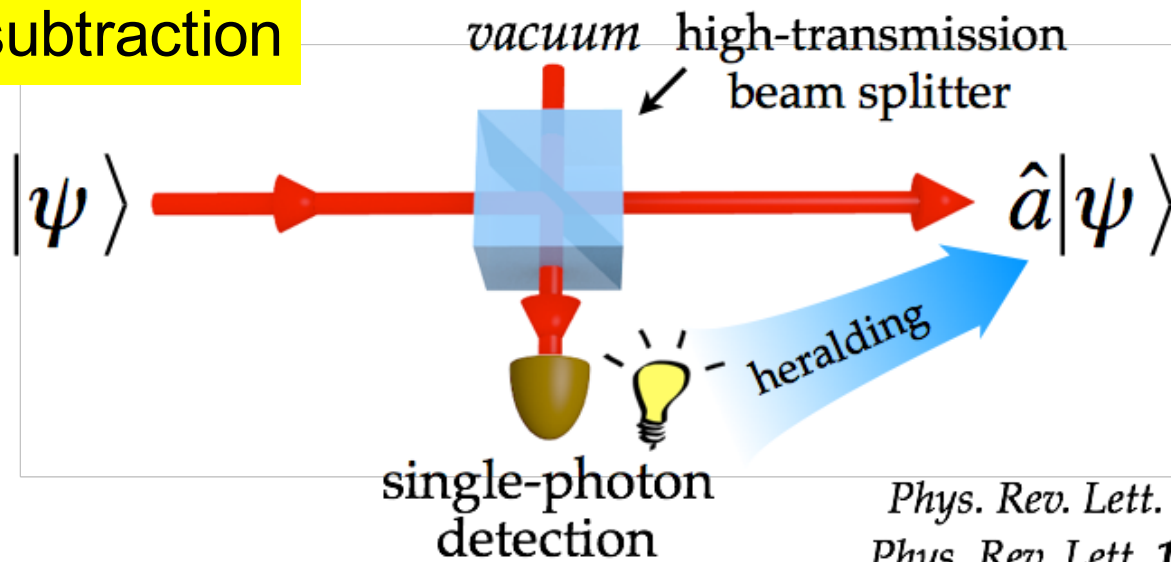


Quantum Correlations (entanglement)

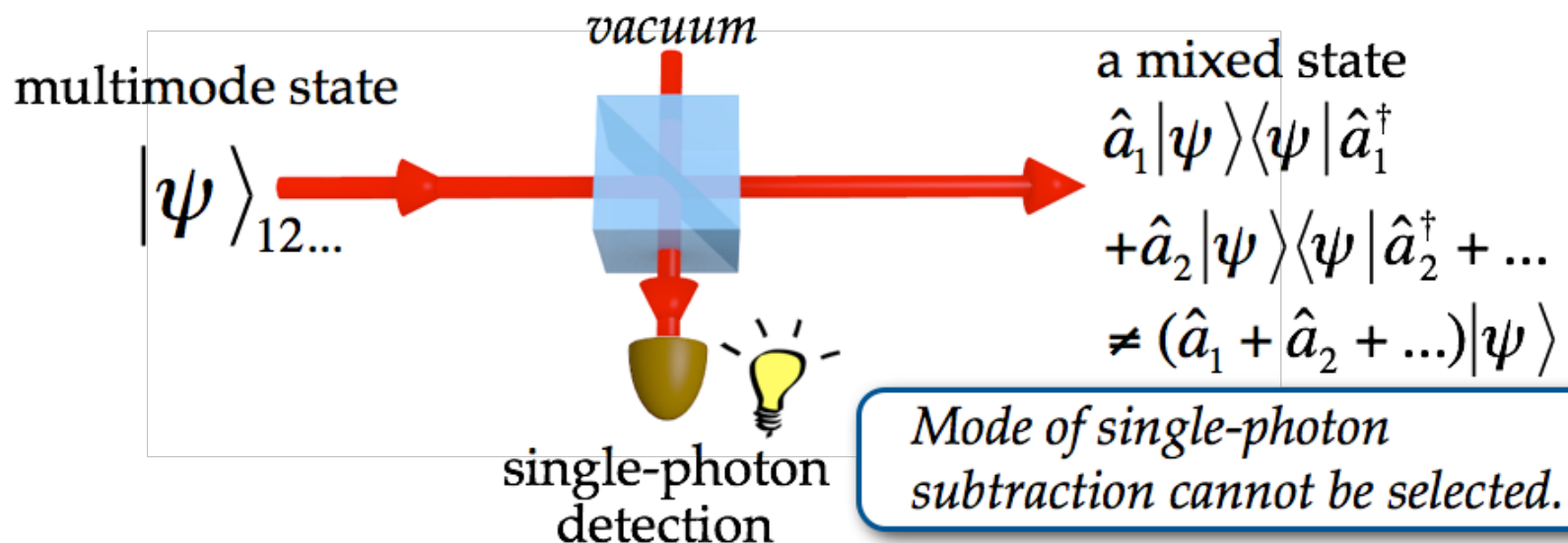
Entanglement between quadratures
/ (i.e. amplitude and phase)



Photon subtraction

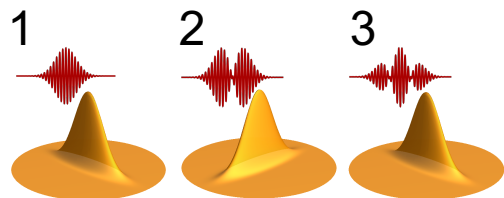


Phys. Rev. Lett. **92**, 153601 (2004)
Phys. Rev. Lett. **110**, 130403 (2013)



Sum Frequency Generation

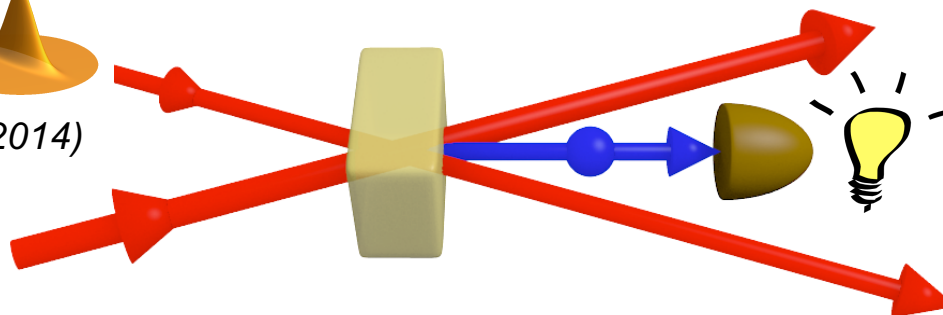
Signal



Nature Photon. **8**, 109 (2014)

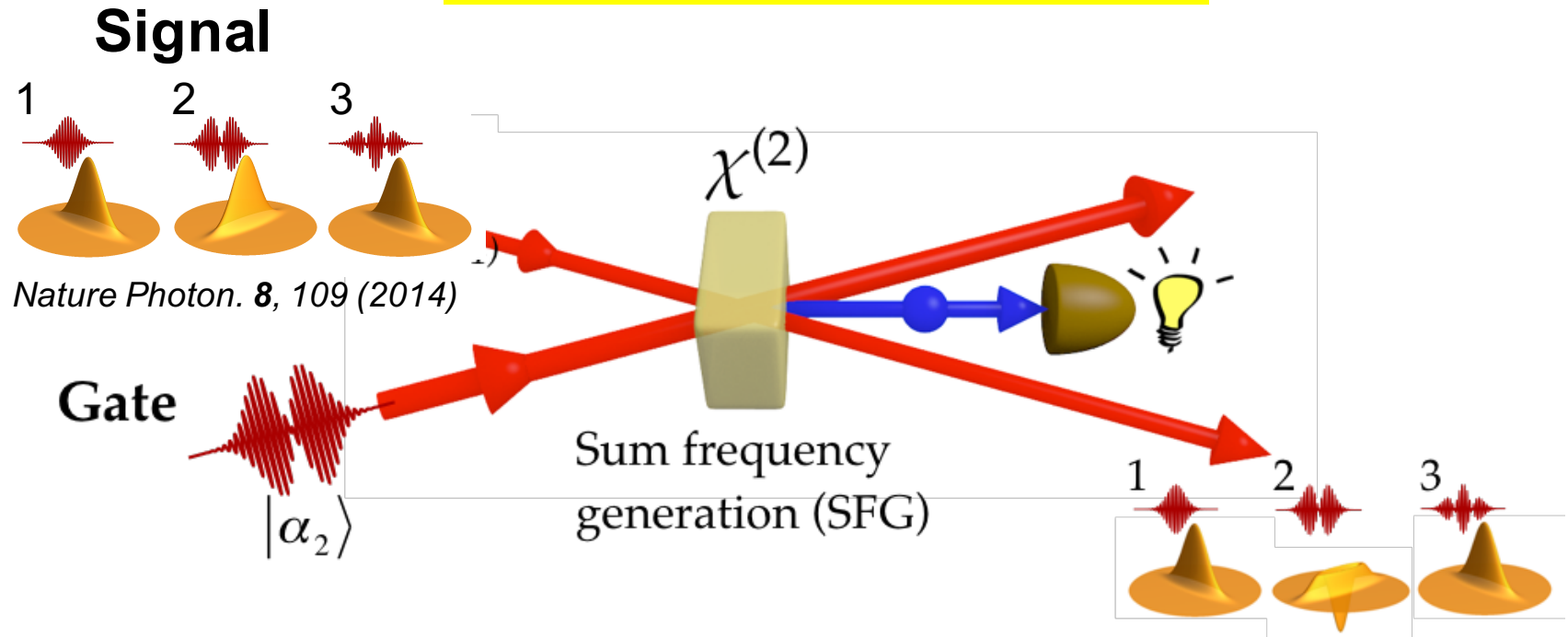
Gate

$$|\alpha_1, \alpha_2, \dots\rangle$$



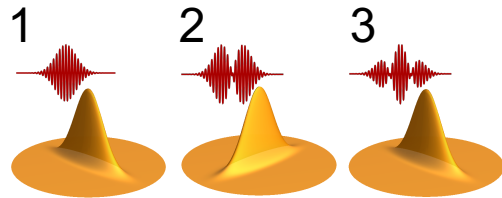
$$\propto (\alpha_1 \hat{a}_1 + \alpha_2 \hat{a}_2 + \dots) |\psi_s\rangle$$

Sum Frequency Generation



Sum Frequency Generation

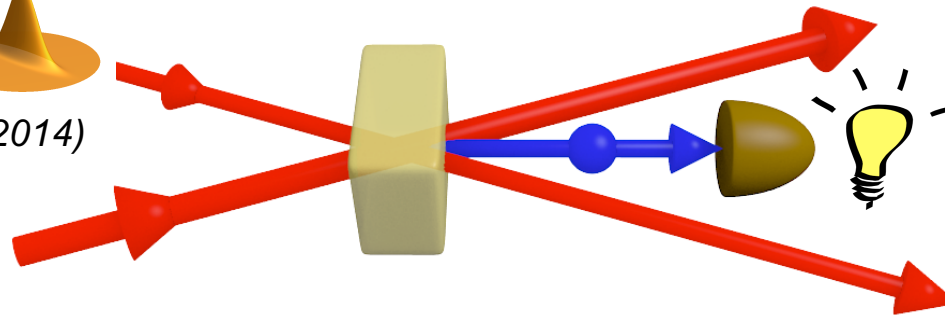
Signal



Nature Photon. **8**, 109 (2014)

Gate

$$|\alpha_1, \alpha_2, \dots\rangle$$

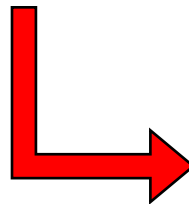


$$\propto (\alpha_1 \hat{a}_1 + \alpha_2 \hat{a}_2 + \dots) |\psi_s\rangle$$

$$\hat{\rho}^{(out)} = \sum_{i,j} \chi_{i,j} \hat{a}_i \rho^{(in)} \hat{a}_j^\dagger$$

Eigenmodes of the process

- > Schmidt number
- > purity



$$\hat{\rho}^{(out)} \propto \sum_k p_k \hat{A}_k \rho^{(in)} \hat{A}_k^\dagger$$

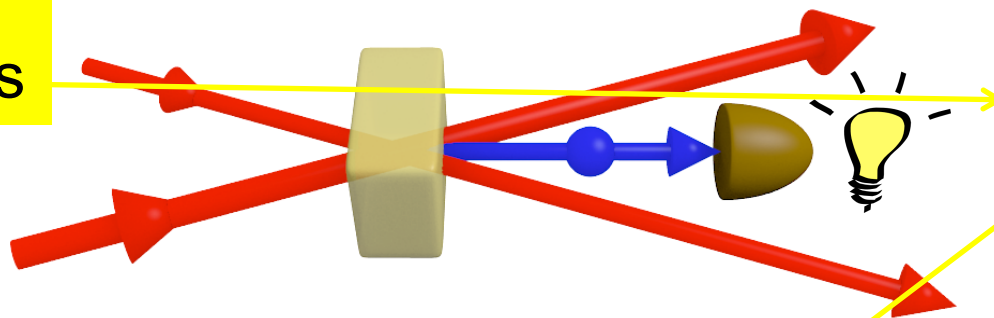
Sum Frequency Generation

Signal

Probe with
coherent states

Gate

$$|\alpha_1, \alpha_2, \dots\rangle$$



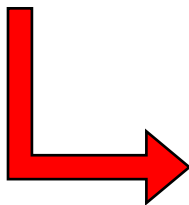
Compute
success rate

$$\propto (\alpha_1 \hat{a}_1 + \alpha_2 \hat{a}_2 + \dots) |\psi_s\rangle$$

$$\hat{\rho}^{(out)} = \sum_{i,j} \chi_{i,j} \hat{a}_i \rho^{(in)} \hat{a}_j^\dagger$$

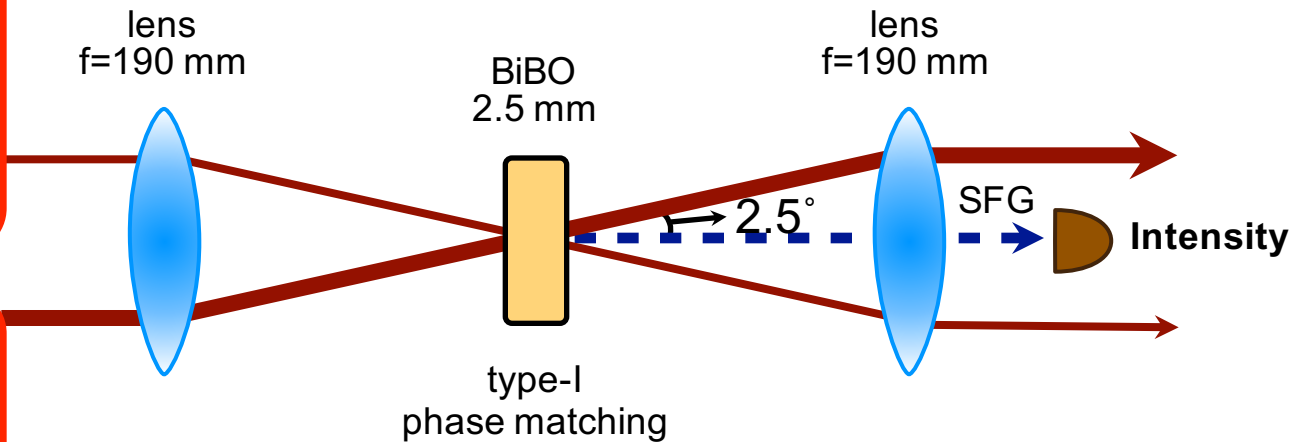
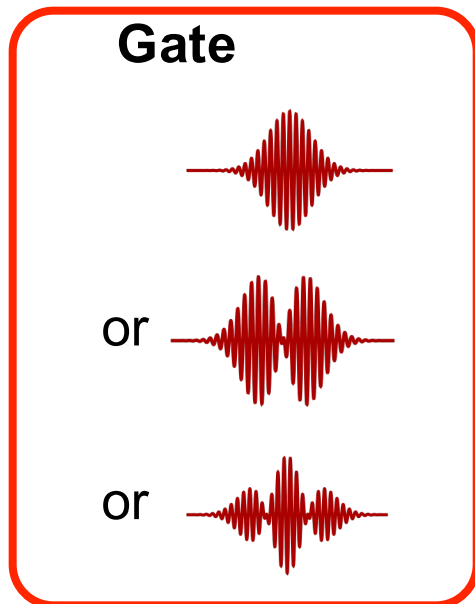
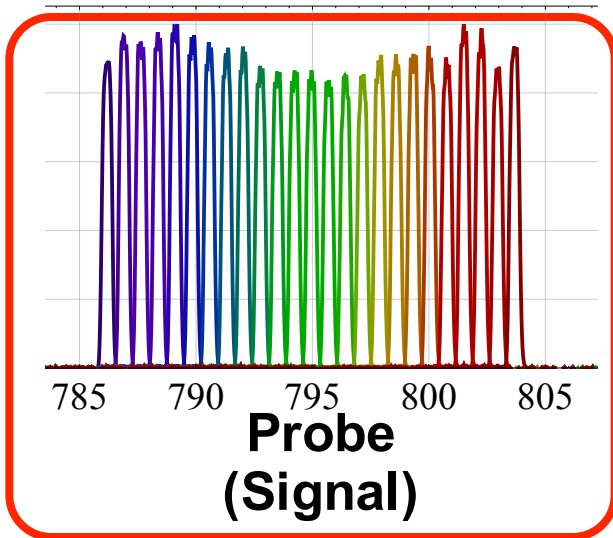
Eigenmodes of the process

- > Schmidt number
- > purity



$$\hat{\rho}^{(out)} \propto \sum_k p_k \hat{A}_k \rho^{(in)} \hat{A}_k^\dagger$$

Probe basis: wavelength bands



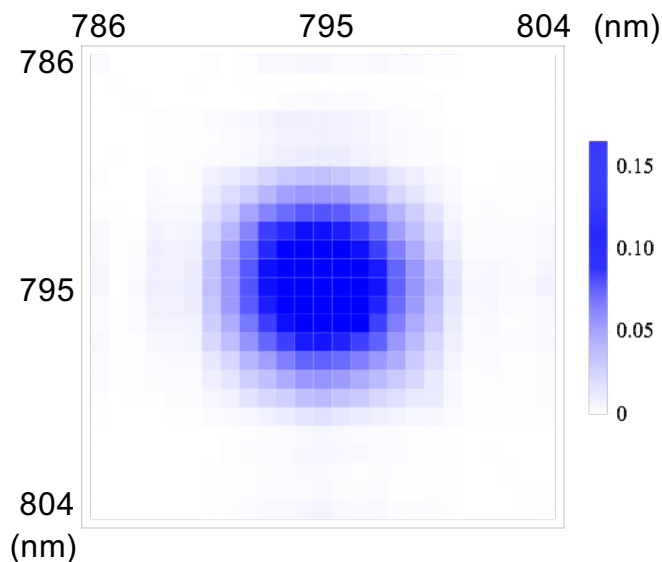
Gate Pulse shape:
Tune the single-photon subtractors.

Reconstruction using Maximum Likelihood Estimation

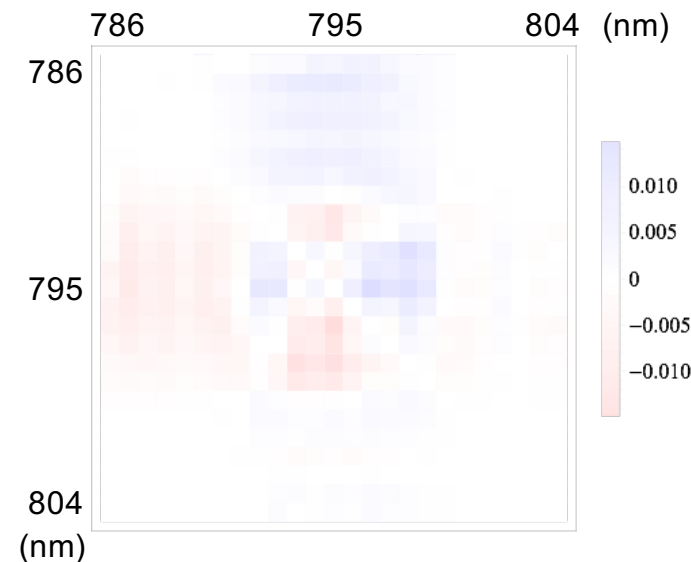
Gate



$\text{Re}[\chi]$



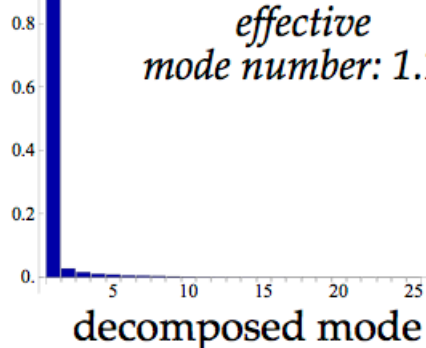
$\text{Im}[\chi]$



Diagonalisation: mode eigenvalues and shapes

contribution, P_k

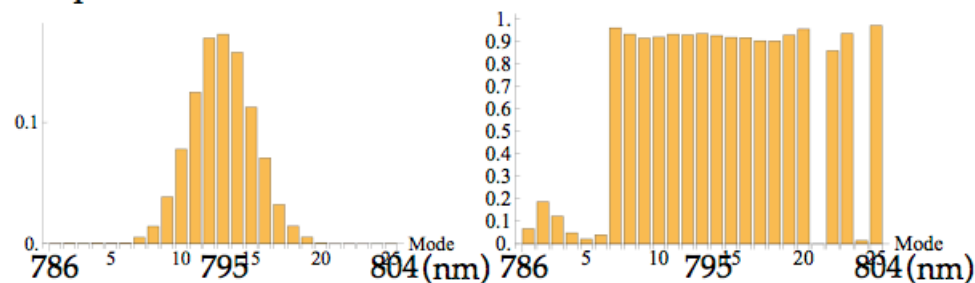
*effective
mode number: 1.18*



the first subtraction mode \hat{A}_1

Amplitude

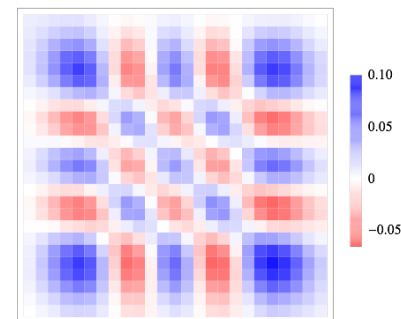
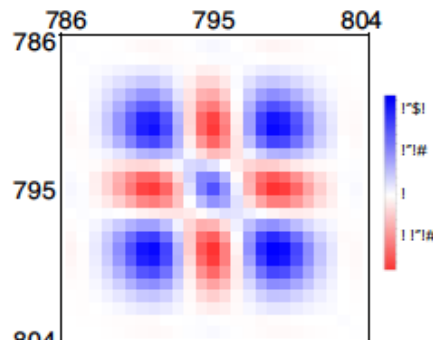
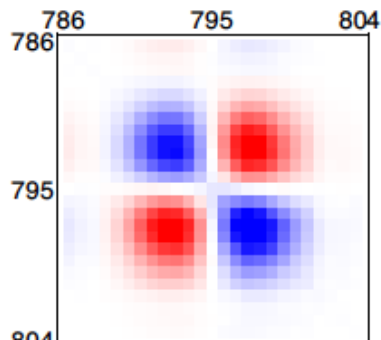
Phase



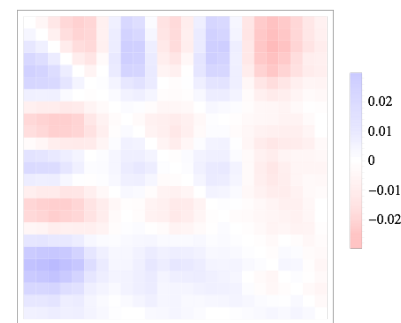
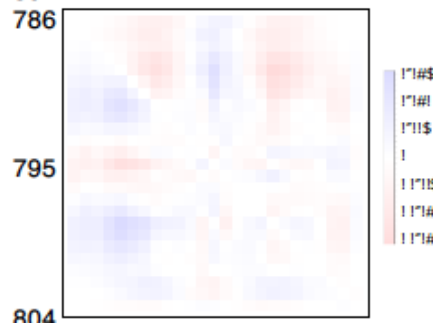
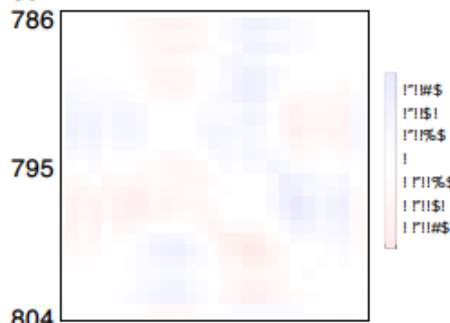
Gate



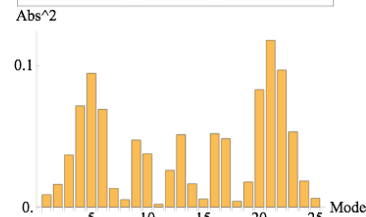
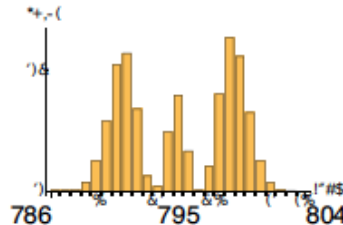
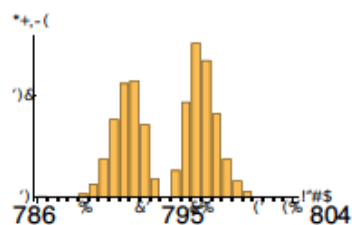
$\text{Re}[\chi]$



$\text{Im}[\chi]$



Subtraction mode



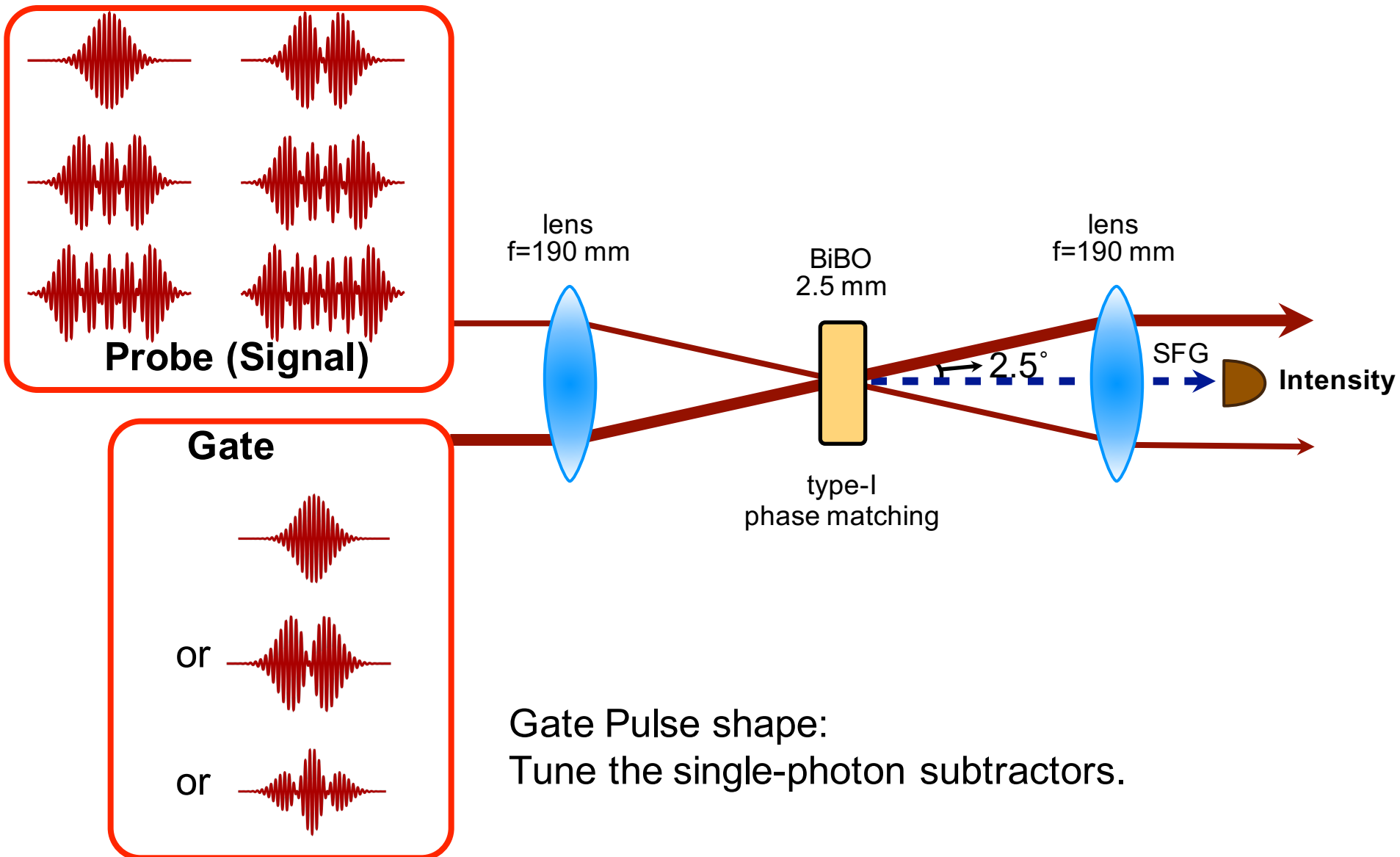
Number of modes

1.2 mode

1.24 mode

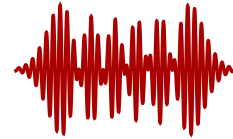
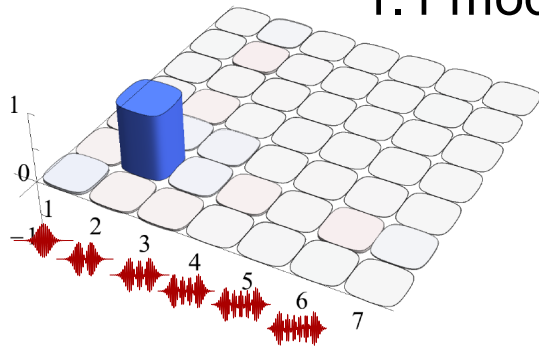
1.36 mode

Probe basis: Hermite Gauss modes

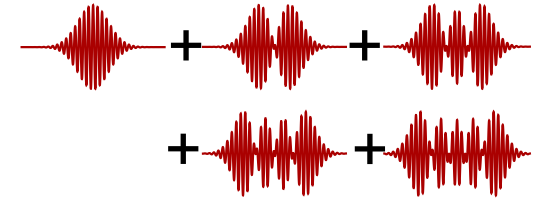
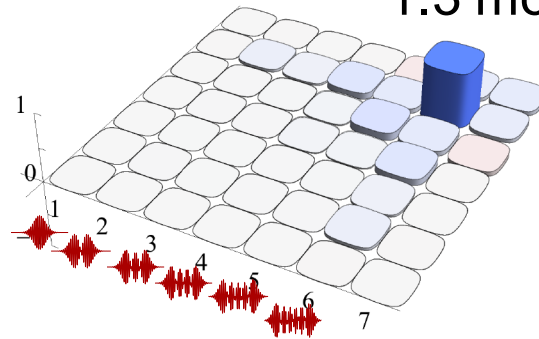




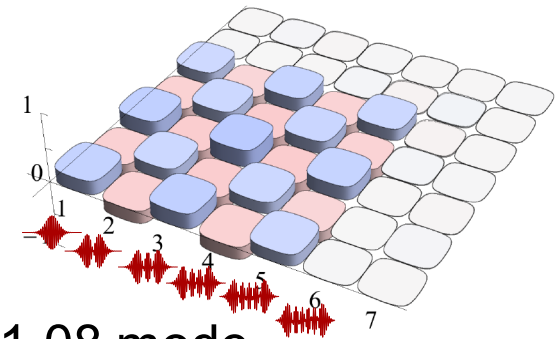
1.1 mode



1.3 mode



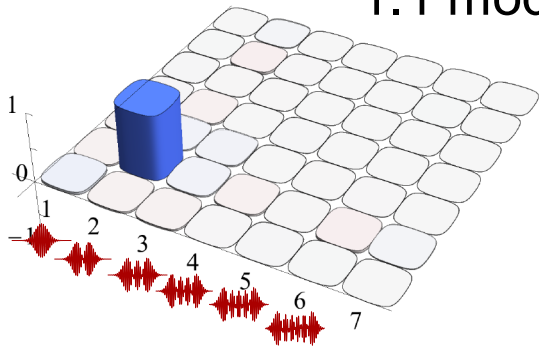
1.08 mode



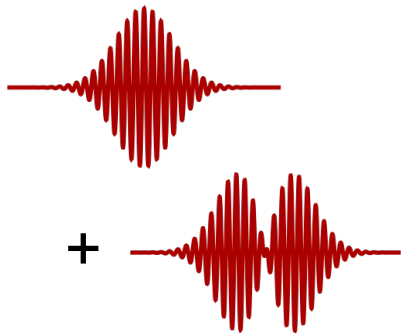
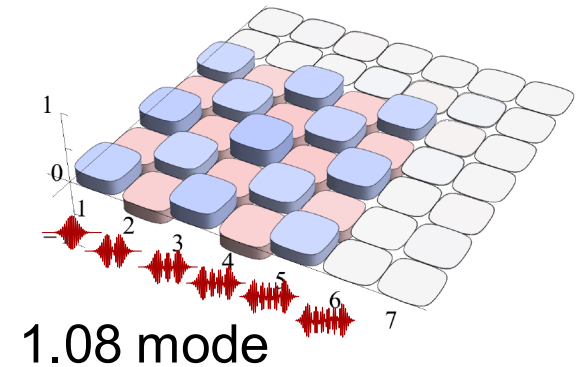
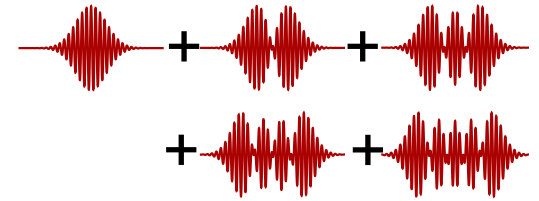
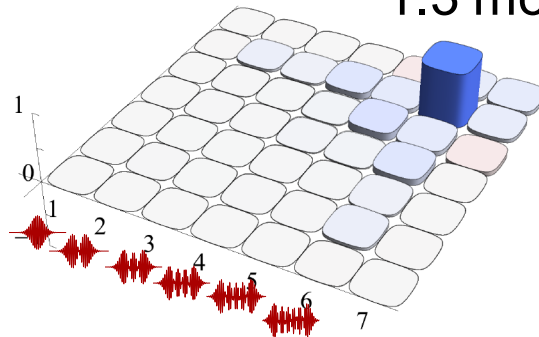
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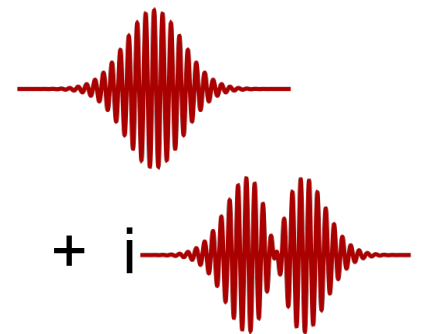
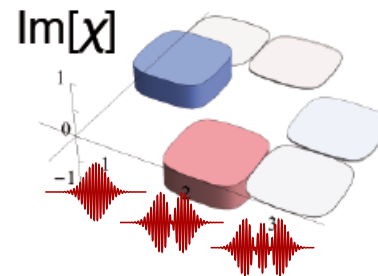
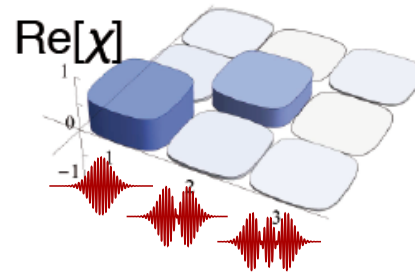
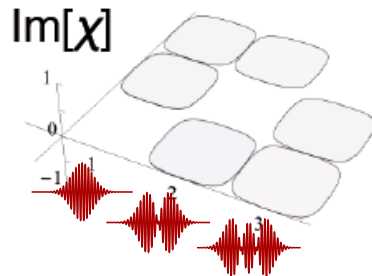
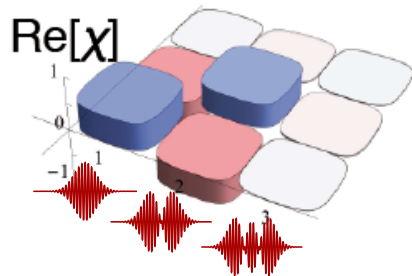
1.1 mode



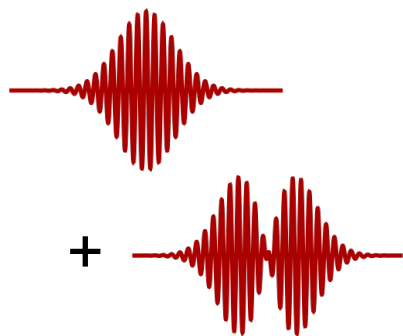
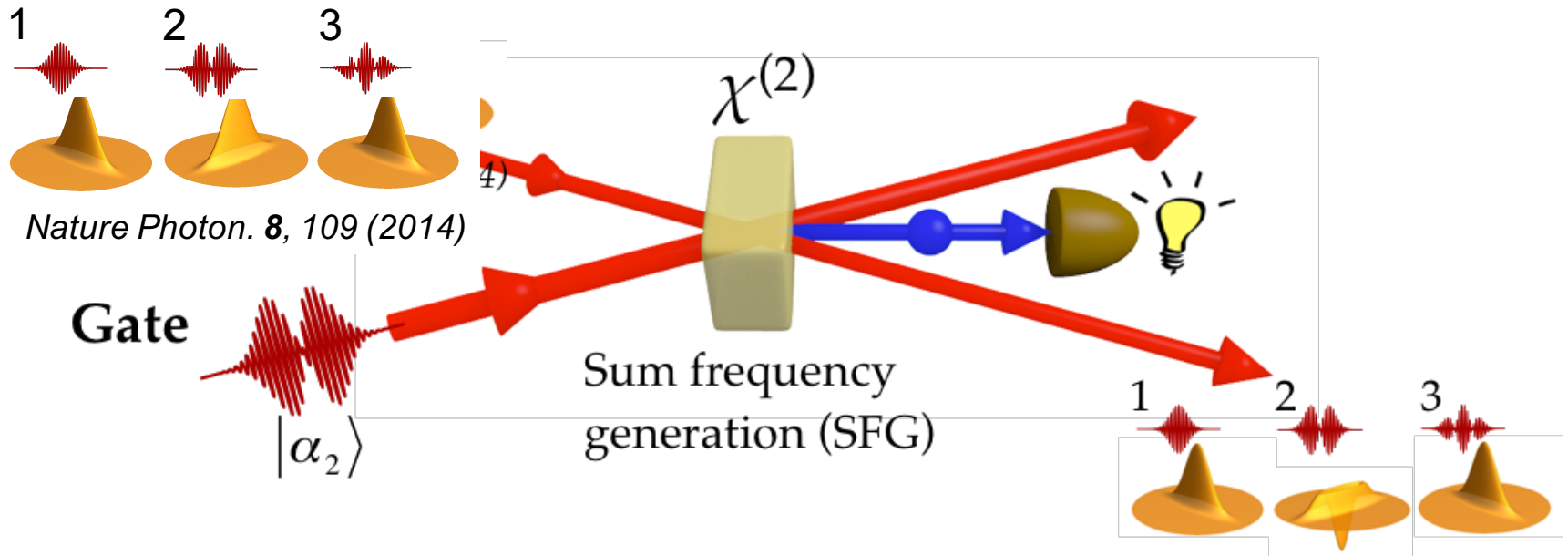
1.3 mode



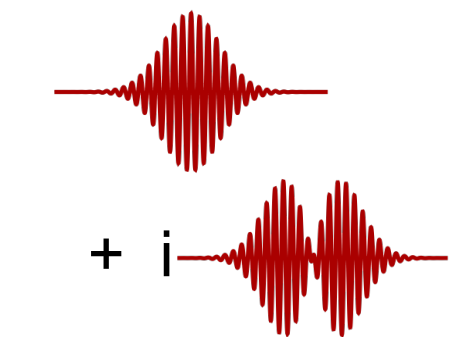
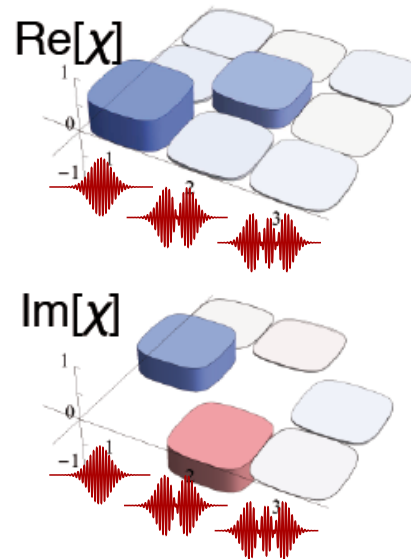
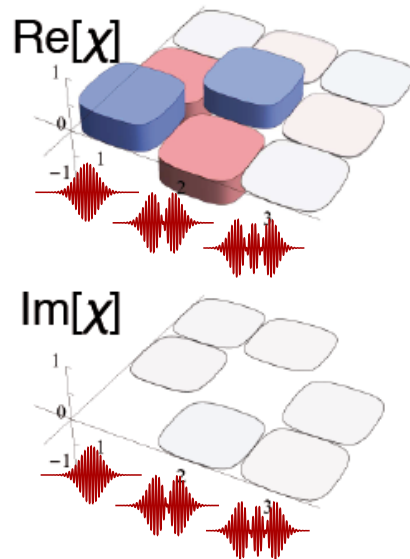
1.08 mode



1.09 mode



1.08 mode



1.09 mode

Outline

Multimode quantum optics in Quantum Information technologies

Introduction

Multicolor entanglement

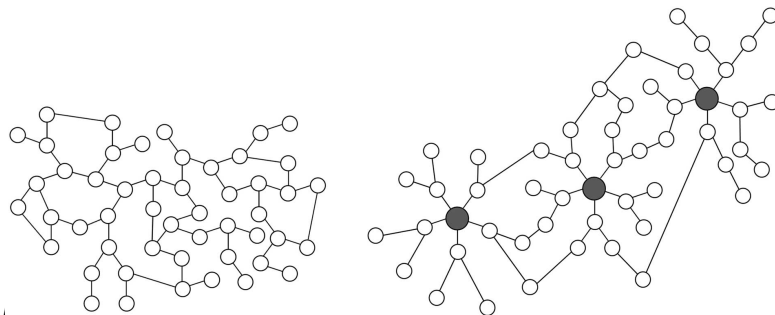
Towards measurement based quantum computing

Simulation of complex quantum networks

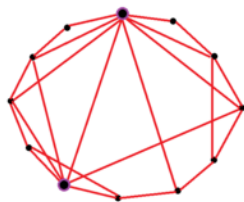
Probing a structured environment

Energy transport: some ideas

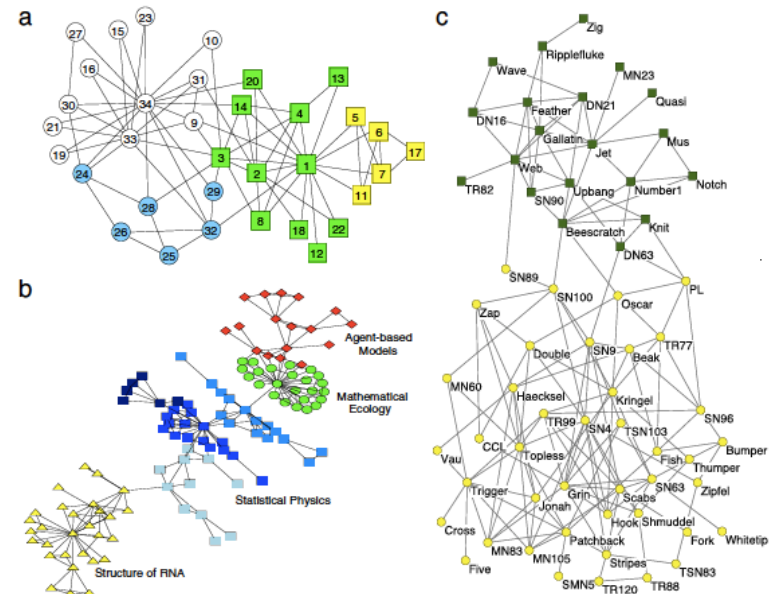
in this talk, complex network = not regular lattice



International Journal of Oncology, 10.3892/ijo.2013.2114



small world network
wikipedia



Community structures

S. Fortunato Physics Report 486

Why?

Bosonic networks

- **Fundamental reasons:**
understanding the interplay between their complex structure and their quantum properties

G. Bianconi and C. Rahmede Phys. Rev. E **93**, 032315 (2016)

- **Quantum networks: at the base of quantum information protocols.**
*Futures quantum technologies : quantum complex networks,
ex: quantum communication in a complex www*

G. D. Paparo & M. A. Martin-Delgado, Scientific Reports 2 , 444 (2012)

- **Open systems: environment described by complex quantum networks**

J. Nokkala, F. Galve, R. Zambrini, S. Maniscalco and J. Piilo, Scientific Reports 6, 26861 (2016)

- **Energy transfer in quantum complex structures, ex: light harvesting**

M. Walschaers, F. Schlawin, T. Wellens, and A. Buchleitner, Annu. Rev. Condens. Matter Phys. 2016. 7:223–48

M. Faccin, P. Migdal, T. H. Johnson, V. Bergholm, and J. D. Biamonte, Phys. Rev. X **4**, 041012 (2014).

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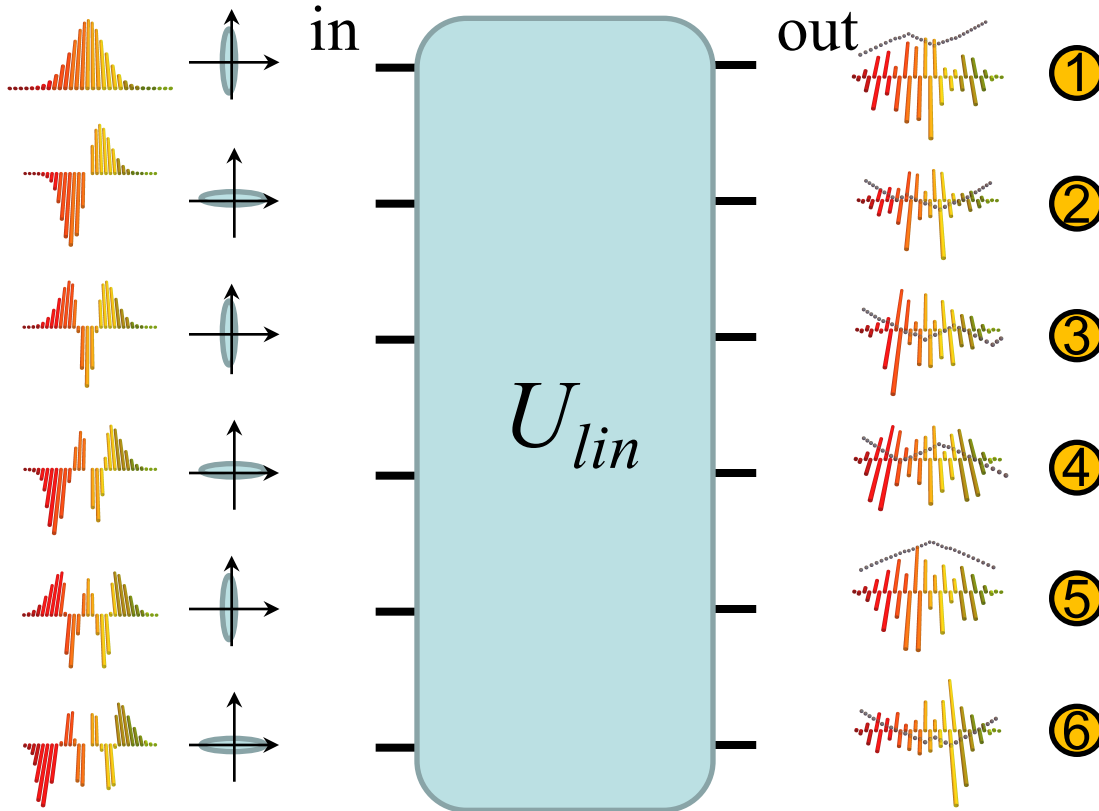
\nearrow basis change
 \nwarrow squeezing

if \vec{a} collection of vacuum states

$$\vec{a}_{out} = V^\dagger S \vec{a}$$

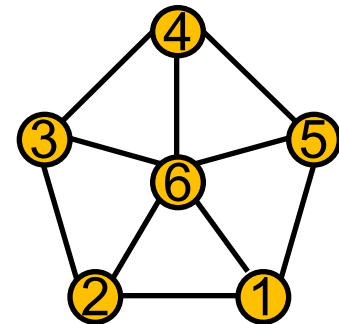
\nwarrow Supermodes

Cluster states

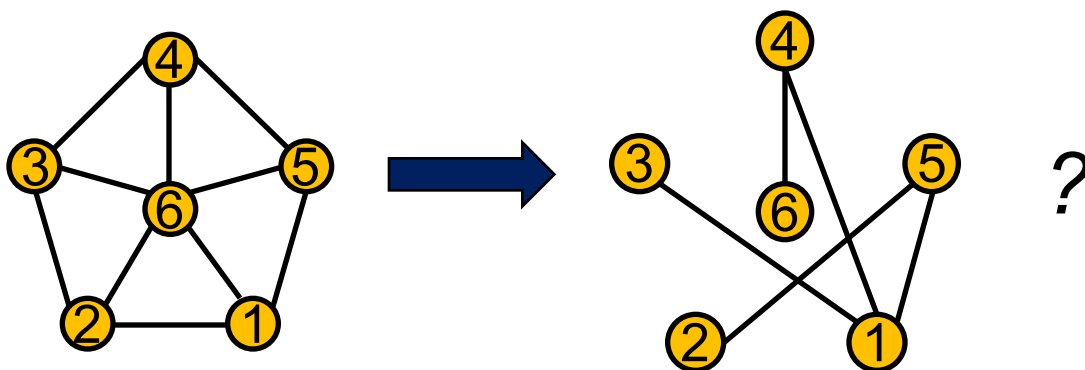


$$\Delta \left(\hat{p}_n - \sum_b \hat{x}_b \right) \rightarrow 0$$

nullifier



Simulation of complex quantum networks



$$\vec{a}_{out} = V^\dagger S \vec{a}$$

***Find the Bloch-Messiah decomposition
of its evolution and experimentally
implement V and S***

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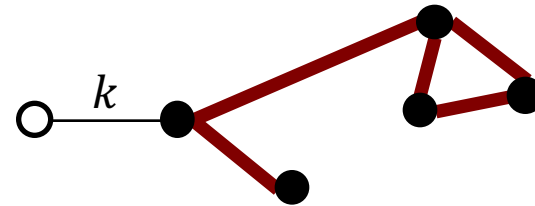
M. Faccin, P. Migdal, T. H. Johnson, V. Bergholm, and J. D. Biamonte, Phys. Rev. X **4**, 041012 (2014).

Simulation of complex structured environment +1 system/probe connected to one node

$$H_E = \frac{\mathbf{p}^T \mathbf{p}}{2} + \mathbf{q}^T \mathbf{A} \mathbf{q}$$

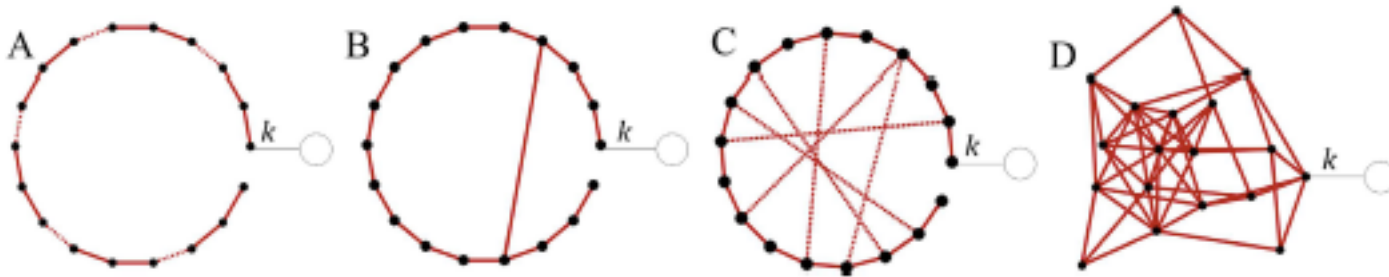
$$H_S = (p_S^2 + \omega_S^2 q_S^2)/2,$$

$$H_I = -k q_S q_i$$



$$A_{ij} = \delta_{ij} \omega_i^2 / 2 - (1 - \delta_{ij}) h_{ij} / 2$$

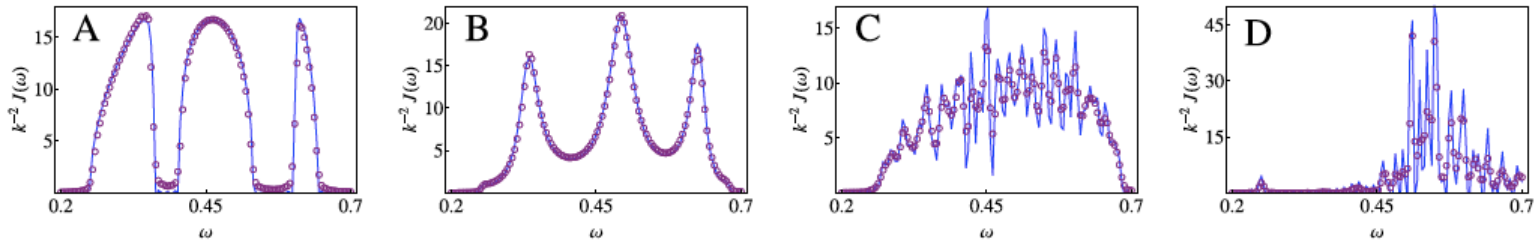
Evolution given by $H_E + H_S + H_I \longrightarrow$ Measure the state of the probe and recover the structure




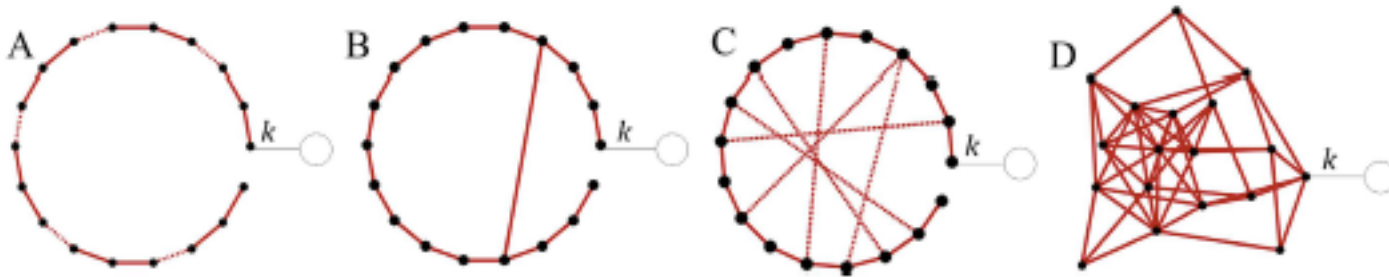
$$J(\omega) = \frac{\pi}{2} \sum_i \frac{k^2 g_i^2}{\Omega_i} \delta(\omega - \Omega_i)$$

spectral density of environmental coupling

$$J(\omega) = \omega \int_0^{t_{\max}} \gamma(t) \cos(\omega t) dt$$

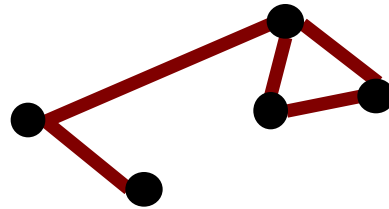


Evolution given by $H_E + H_S + H_I$  Measure the state of the probe and recover the structure



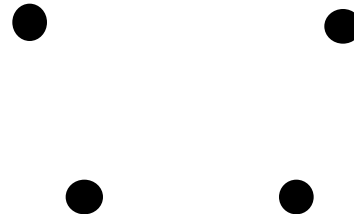
Simulation of complex structured environment

$$H_E = \frac{\mathbf{p}^T \mathbf{p}}{2} + \mathbf{q}^T \mathbf{A} \mathbf{q}$$



$$H_E = \frac{\mathbf{P}^T \mathbf{P}}{2} + \mathbf{Q}^T \mathbf{D} \mathbf{Q}$$

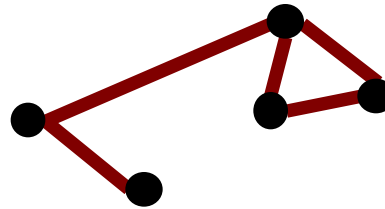
$$\bullet \quad \Omega_i = \sqrt{2D_{ii}}$$



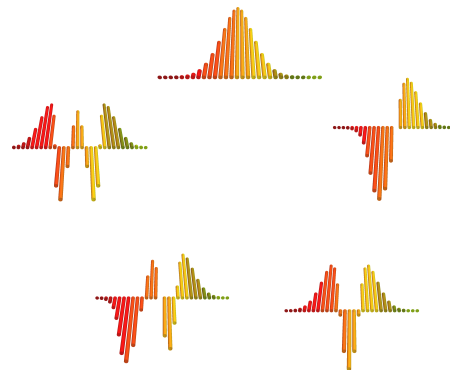
eigenmodes

Simulation of complex structured environment

$$H_E = \frac{p^T p}{2} + q^T A q$$



$$H_E = \frac{P^T P}{2} + Q^T D Q$$



*super
eigenmodes*

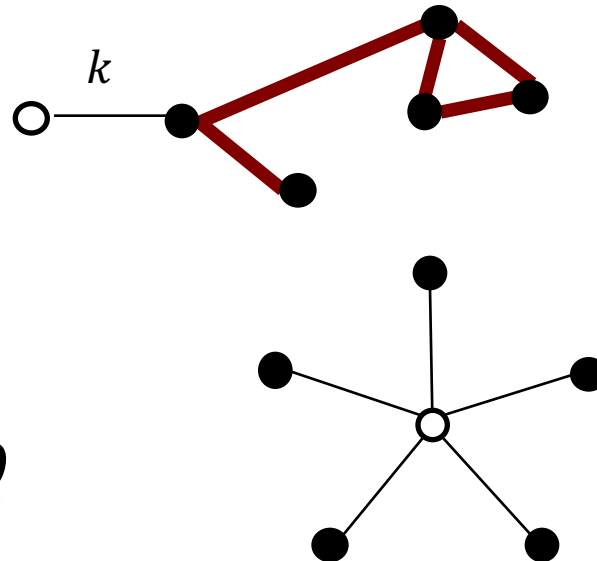
Simulation of complex structured environment +1 system/probe connected to one node

$$H_E = \frac{\mathbf{p}^T \mathbf{p}}{2} + \mathbf{q}^T \mathbf{A} \mathbf{q}$$

$$H_I = -k q_S q_i$$

$$H_E = \frac{\mathbf{P}^T \mathbf{P}}{2} + \mathbf{Q}^T \mathbf{D} \mathbf{Q}$$

$$H_I = -k q_S \mathbf{K}_i \cdot \mathbf{Q}$$



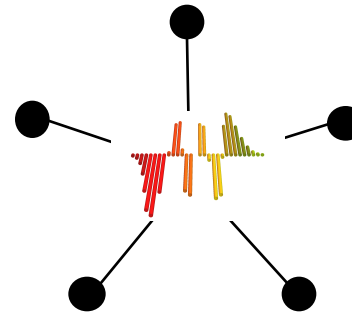
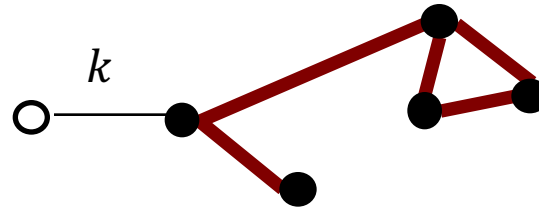
Simulation of complex structured environment +1 system/probe connected to one node

$$H_E = \frac{p^T p}{2} + q^T A q$$

$$H_I = -k q_S q_i$$

$$H_E = \frac{P^T P}{2} + Q^T D Q$$

$$H_I = -k q_S K_i \cdot Q$$



one supplementary
supermode

Simulation of complex structured environment +1 system/probe connected to one node

$$\begin{pmatrix} Q(t) \\ q_s(t) \\ P(t) \\ p_s(t) \end{pmatrix} = \mathcal{S} \begin{pmatrix} Q(0) \\ q_s(0) \\ P(0) \\ p_s(0) \end{pmatrix}$$

time evolution of the system + the network
= symplectic matrix acting on the quadratures

Simulation of complex structured environment +1 system/probe connected to one node

$$\begin{pmatrix} Q(t) \\ q_s(t) \\ P(t) \\ p_s(t) \end{pmatrix} = S \begin{pmatrix} Q(0) \\ q_s(0) \\ P(0) \\ p_s(0) \end{pmatrix}$$

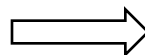
time evolution of the system + the network
= symplectic matrix acting on the quadratures

Bloch-Messiah decomposition

$$S = V K \text{ when possible}$$

*rotation =
basis change*

squeezing



Implement experimentally

Simulation of complex structured environment +1 system/probe connected to one node

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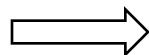
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Implement experimentally

measure the probe

Simulation of complex structured environment +1 system/probe connected to one node

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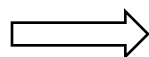
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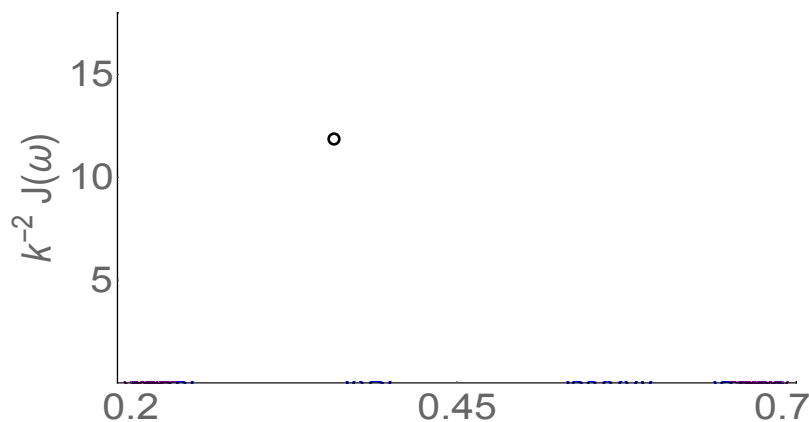
squeezing



Implement experimentally

measure the probe

get a point



Simulation of complex structured environment +1 system/probe connected to one node

$$\begin{pmatrix} Q(t) \\ q_s(t) \\ P(t) \\ p_s(t) \end{pmatrix} = \mathcal{S} \begin{pmatrix} Q(0) \\ q_s(0) \\ P(0) \\ p_s(0) \end{pmatrix}$$

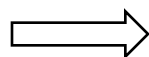
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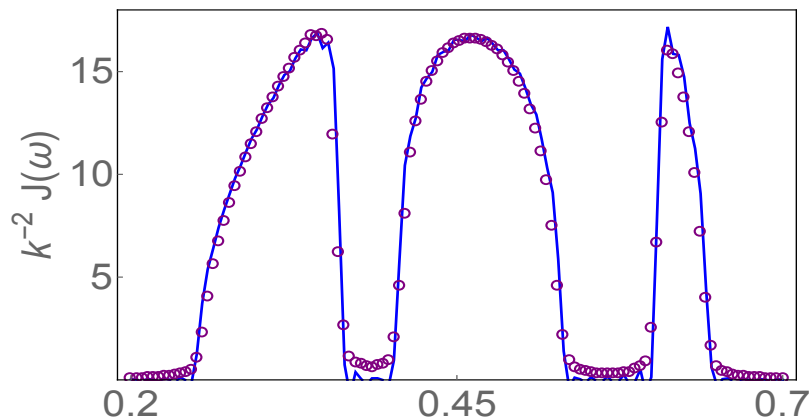


Implement experimentally

measure the probe

get a point

repeat



Simulation of complex structured environment +1 system/probe connected to one node

$$\begin{pmatrix} Q(t) \\ q_s(t) \\ P(t) \\ p_s(t) \end{pmatrix} = \mathcal{S} \begin{pmatrix} Q(0) \\ q_s(0) \\ P(0) \\ p_s(0) \end{pmatrix}$$

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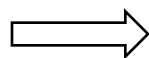
$$\mathcal{S} = \mathcal{V} \mathcal{K} \text{ when possible}$$

rotation=
basis change

squeezing

network structure
encoded in the rotation \mathcal{V}

energy parameters $\{k, g, \omega_0, \omega_s\}$
encoded in the squeezing \mathcal{K}

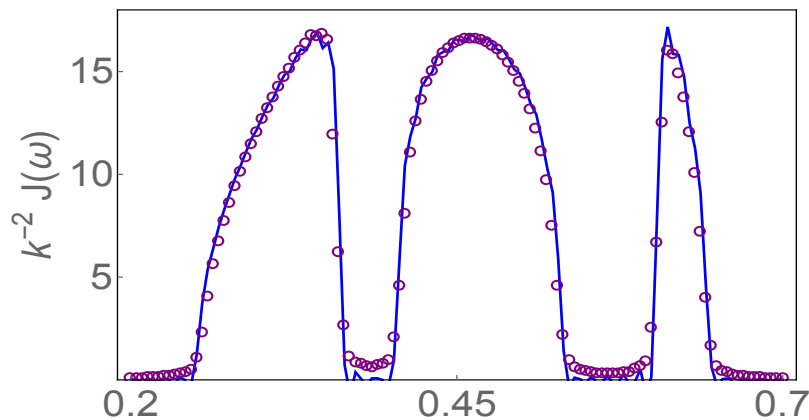


Implement experimentally

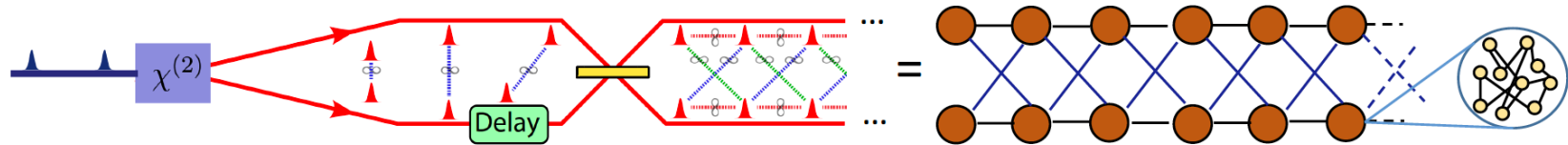
measure the probe

get a point

repeat

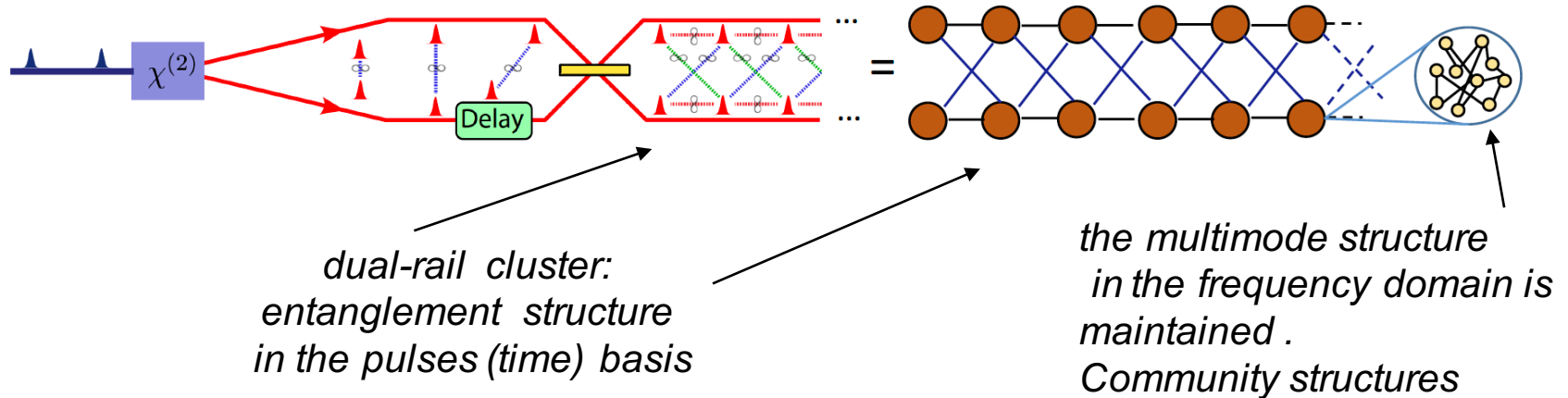


- **Number of modes :16** at the moment but they are only limited by the measurement procedure -> they could be around **40** if we increase the spectral bandwidth of LO
- Exploring new experimental setup -> “big-states” **10^5** modes

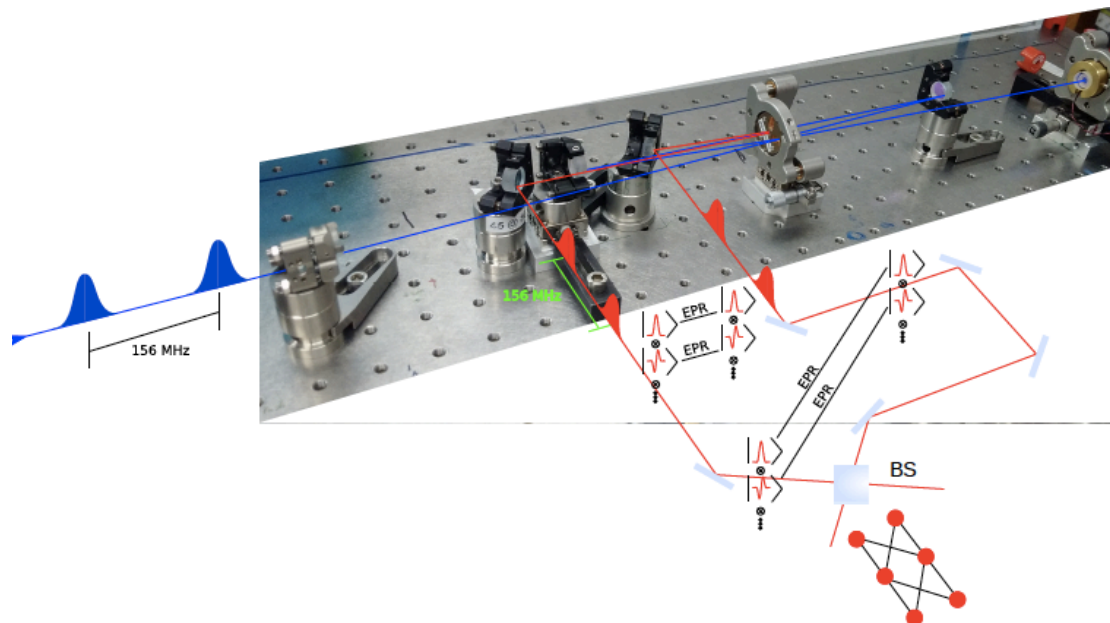
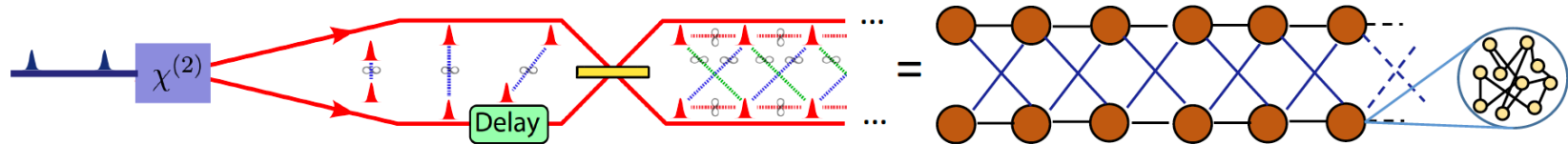


dual-rail cluster:
entanglement structure
in the pulses (time) basis

- **Number of modes :16** at the moment but they are only limited by the measurement procedure -> they could be around **40** if we increase the spectral bandwidth of LO
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- Exploring new experimental setup -> “big-states” **10^5** modes



Outline

Multimode quantum optics in Quantum Information technologies

Introduction

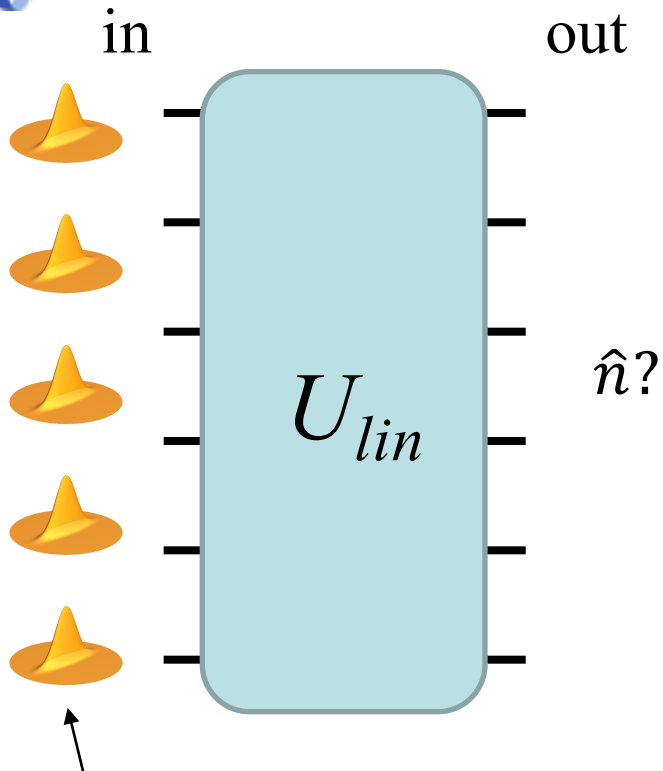
Multicolor entanglement

Towards measurement based quantum computing

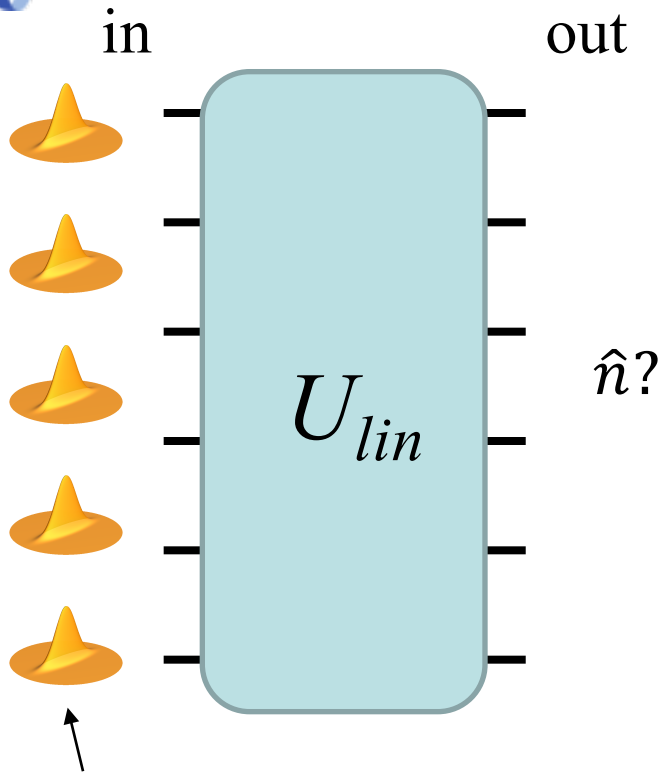
Simulation of complex quantum networks

Probing a structured environment

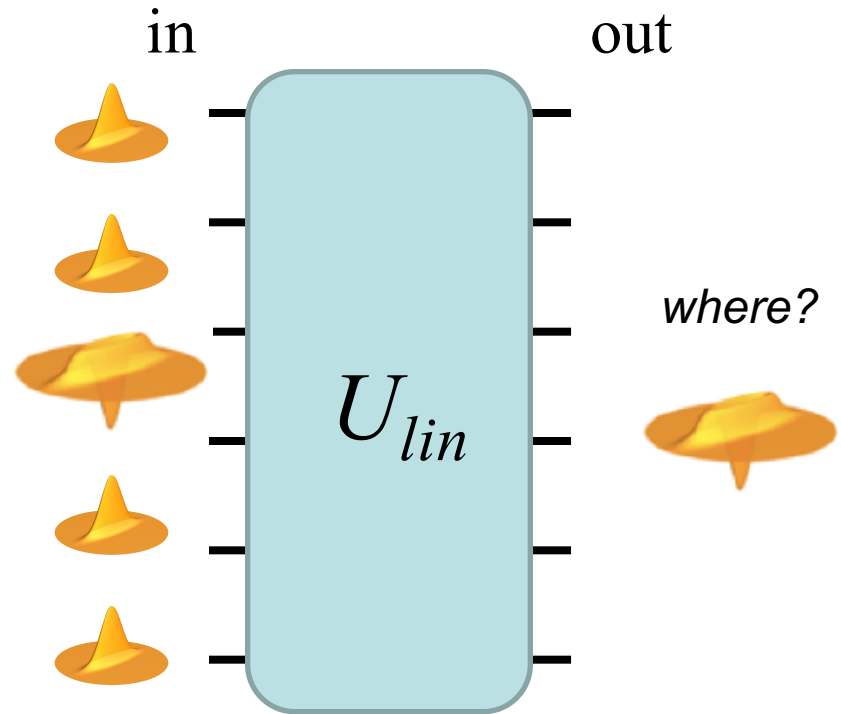
Energy transport: some ideas

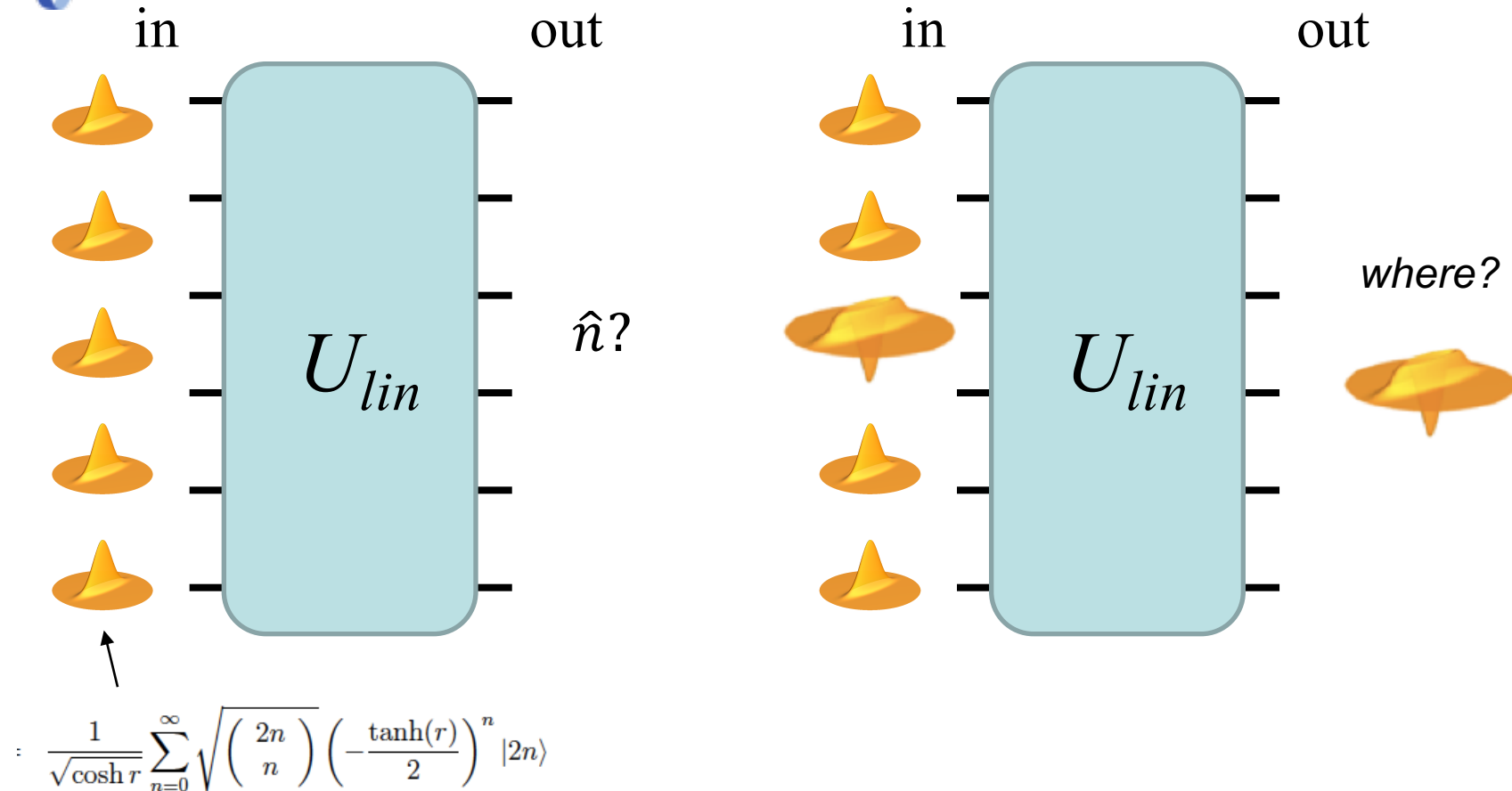


$$: \frac{1}{\sqrt{\cosh r}} \sum_{n=0}^{\infty} \sqrt{\binom{2n}{n}} \left(-\frac{\tanh(r)}{2} \right)^n |2n\rangle$$



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Connection with boson sampling problems

C S. Hamilton, R. Kruse, L. Sansoni, S. Barkhofen, C. Silberhorn, I. Jex , arXiv:1612.01199
 R. N. Alexander, N. C. Gabay, P. P. Rohde, and N. C. Menicucci,
 Phys. Rev. Lett. **118**, 110503 (2017)

Thank you!



