

Continuous-variable probing of structured environments

MATTEO BINA, FEDERICO GRASSELLI AND MATTEO PARIS



UNIVERSITÀ DEGLI STUDI DI MILANO

Aim of the work

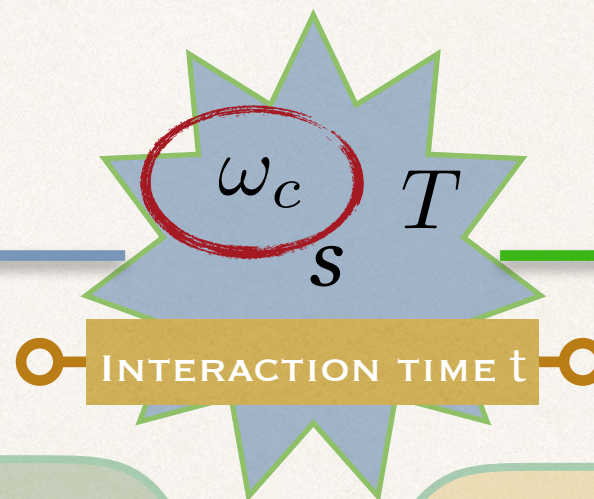
$$\varrho_0 = S(\xi)\nu(n_{\text{th}})S^\dagger(\xi)$$

$$\xi \quad n_{\text{th}}$$

PROBE STATE

$$\varrho_0$$

STRUCTURED
ENVIRONMENT



EVOLVED STATE

$$\varrho_{\omega_c}(t)$$

MEASUREMENT

SPECTRAL DENSITY

$$J(\omega) = \omega_c \left(\frac{\omega}{\omega_c} \right)^s e^{-\frac{\omega}{\omega_c}}$$

✦ OHMIC $s = 1$

✦ SUB-OHMIC $s = \frac{1}{2}$

✦ SUPER-OHMIC $s = 3$

✦ OPTIMAL POVM \rightarrow QFI

✦ FEASIBLE MEASUREMENT $\mathcal{O} \rightarrow$ FI

QUANTUM CRAMER-RAO BOUND

$$\text{Var}_{\omega_c} \geq \frac{1}{F_{\mathcal{O}}(\omega_c)} \geq \frac{1}{H(\omega_c)}$$

The Model

SYSTEM-ENVIRONMENT HAMILTONIAN

$$H = \frac{\hbar\omega_0}{2} (P^2 + X^2) + \sum_n \frac{\hbar\omega_n}{2} (P_n^2 + X_n^2) - \alpha X \otimes \sum_n \hbar\gamma_n X_n$$

- ♦ Hu, Paz, Zhang, PRD **45**, 2843 (1992)
- ♦ Vasile *et al.*, PRA **80**, 062324 (2009)

SYSTEM MASTER EQUATION

Born (weak-coupling) approximation only! $\alpha \ll 1$

$$\frac{d}{dt}\varrho(t) = -\frac{i}{\hbar}[H_0, \varrho(t)] + i r(t)[X^2, \varrho(t)] - i\gamma(t)[X, \{P, \varrho(t)\}] - \Delta(t)[X, [X, \varrho(t)]] + \Pi(t)[X, [P, \varrho(t)]]$$

RESERVOIR SPECTRAL DENSITY

$$J(\omega) = \omega_c \left(\frac{\omega}{\omega_c} \right)^s e^{-\frac{\omega}{\omega_c}}$$

$$r(t) = \alpha^2 \int_0^t d\tau \cos(\omega_0\tau) \int_0^\infty d\omega J(\omega) \sin(\omega\tau)$$

$$\gamma(t) = \alpha^2 \int_0^t d\tau \sin(\omega_0\tau) \int_0^\infty d\omega J(\omega) \sin(\omega\tau)$$

$$\Delta(t) = \alpha^2 \int_0^t d\tau \cos(\omega_0\tau) \int_0^\infty d\omega \coth\left(\frac{\hbar\omega}{2k_B T}\right) J(\omega) \cos(\omega\tau)$$

$$\Pi(t) = \alpha^2 \int_0^t d\tau \sin(\omega_0\tau) \int_0^\infty d\omega \coth\left(\frac{\hbar\omega}{2k_B T}\right) J(\omega) \cos(\omega\tau)$$

The Model

PROBE: SINGLE-MODE GAUSSIAN STATE

$$\varrho_0 = S(\xi)\nu(n_{\text{th}})S^\dagger(\xi) \quad \longrightarrow \quad \chi[\vec{z}] = \exp \left[-\frac{1}{2} \vec{z}^T \Omega \sigma \Omega^T \vec{z} - i \vec{z}^T \Omega \vec{\delta} \right]$$

characteristic function

COVARIANCE MATRIX

Null first-moment vector $\vec{\delta} = 0$

$$\sigma_0 = \left(\frac{1}{2} + n_{\text{th}} \right) \begin{pmatrix} e^{2\xi} & 0 \\ 0 & e^{-2\xi} \end{pmatrix} \quad \longrightarrow \quad \sigma(t) = \Delta_\Gamma(t) \mathbb{I} + e^{-\Gamma(t)} R(t) \sigma_0 R^{-1}(t)$$

Secular approximation: neglect fast rotating terms $2\omega_0 t$

Evolved probe state depends on:

squeezing

ξ

initial thermal photons

n_{th}

interaction time

t

reservoir temperature

T

$$\Delta_\Gamma(t) \equiv e^{-\Gamma(t)} \int_0^t e^{\Gamma(\tau)} \Delta(\tau) d\tau$$

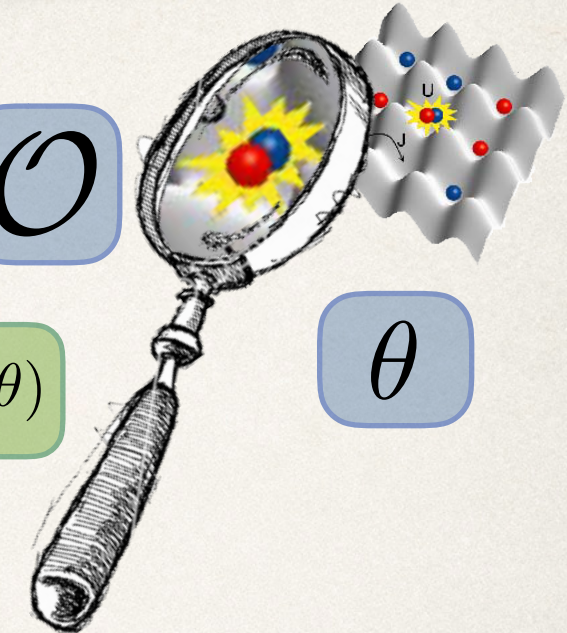
$$\Gamma(t) = 2 \int_0^t \gamma(\tau) d\tau$$

$$R(t) \approx \begin{pmatrix} \cos(\omega_0 t) & \sin(\omega_0 t) \\ -\sin(\omega_0 t) & \cos(\omega_0 t) \end{pmatrix}$$



Q.E.T.**Quantum Estimation Theory**

Estimate a non-directly observable parameter θ :

 \mathcal{O}  θ

Observable $\mathcal{O} \longrightarrow$ Set of data $x = \{x_1, \dots, x_M\} \longrightarrow$ Sample probability $p(x|\theta)$

Mean value $E_{\theta}[\hat{\theta}] = \int dx_1 \dots dx_M p(x_1, \dots, x_M|\theta) \hat{\theta}(x_1, \dots, x_M)$

Variance $\text{Var}[\hat{\theta}] = E \left[\left(\hat{\theta} - E_{\theta}[\hat{\theta}] \right)^2 \right] = \int dx_1 \dots dx_M p(x_1, \dots, x_M|\theta) \left(\hat{\theta} - E_{\theta}[\hat{\theta}] \right)^2$

FISHER INFORMATION (FI)

$$F_{\mathcal{O}}(\theta) = \int dx \frac{[\partial_{\theta} p(x|\theta)]^2}{p(x|\theta)}$$

CRAMER-RAO BOUND

$$\text{Var}_{\theta} \geq \frac{1}{F_{\mathcal{O}}(\theta)}$$

QUANTUM FISHER INFORMATION (QFI)

$H(\theta)$ **Optimized over all measurements**

QUANTUM CRAMER-RAO BOUND

$$\text{Var}_{\theta} \geq \frac{1}{F_{\mathcal{O}}(\theta)} \geq \frac{1}{H(\theta)}$$



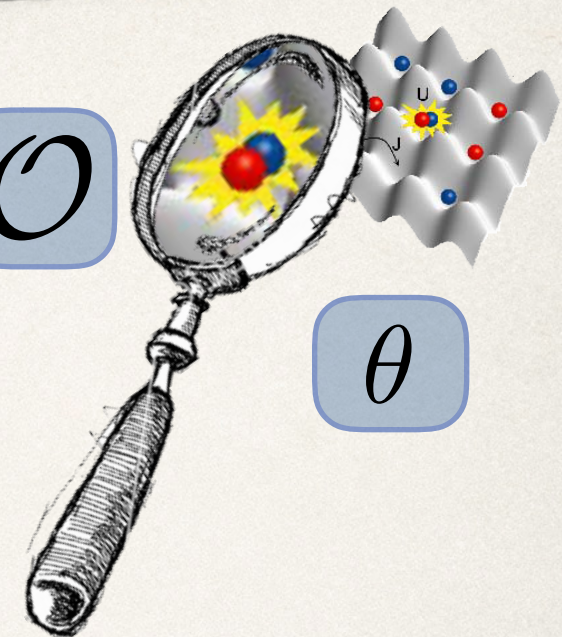
QUProCS MEETING II - PALMA DE MALLORCA 6-7 APRIL 2017

♦ Z. Jiang, Phys. Rev. A **89**, 032128 (2014)

GAUSSIAN STATES WITH $\vec{\delta} = 0$

$$H(\theta) = \frac{1}{2\lambda^4 - 1/8} \left\{ \lambda^4 \text{Tr}[(\sigma^{-1} \dot{\sigma})^2] - \frac{1}{4} \text{Tr}[(\Omega \dot{\sigma})^2] \right\}$$

\mathcal{O}



θ

The parameter of interest is the CUTOFF FREQUENCY

θ



ω_c

FISHER INFORMATION (FI)

$$F_{\mathcal{O}}(\theta) = \int dx \frac{[\partial_{\theta} p(x|\theta)]^2}{p(x|\theta)}$$

CRAMER-RAO BOUND

$$\text{Var}_{\theta} \geq \frac{1}{F_{\mathcal{O}}(\theta)}$$

QUANTUM FISHER INFORMATION (QFI)

$H(\theta)$ Optimized over all measurements

QUANTUM CRAMER-RAO BOUND

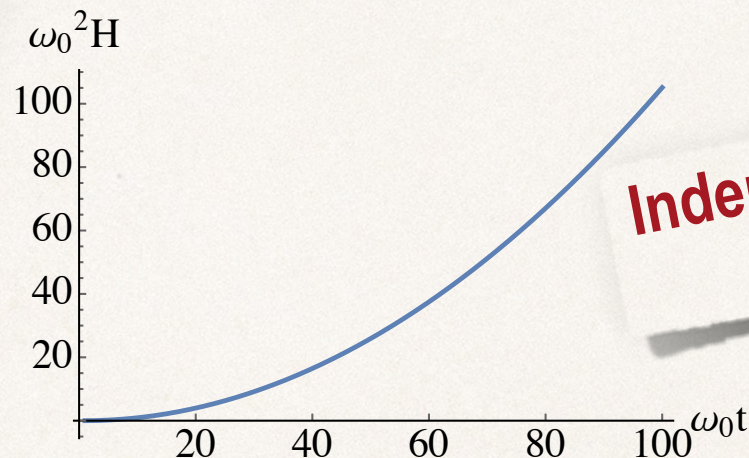
$$\text{Var}_{\theta} \geq \frac{1}{H(\theta)} \geq \frac{1}{F_{\mathcal{O}}(\theta)}$$



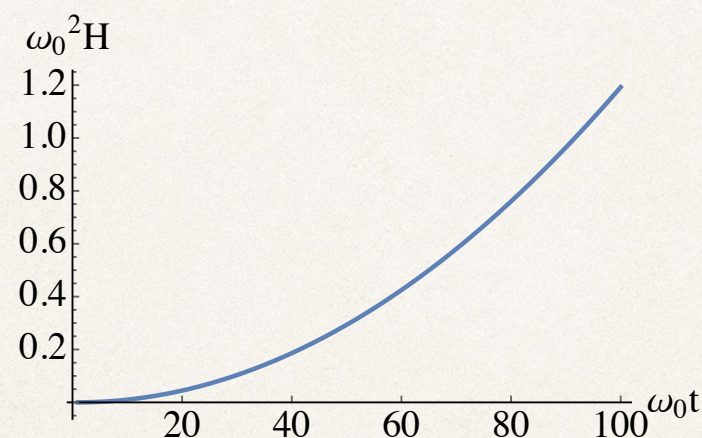
Results: QFI

RESERVOIR TEMPERATURE

HIGH - T



LOW - T



QFI : TIME

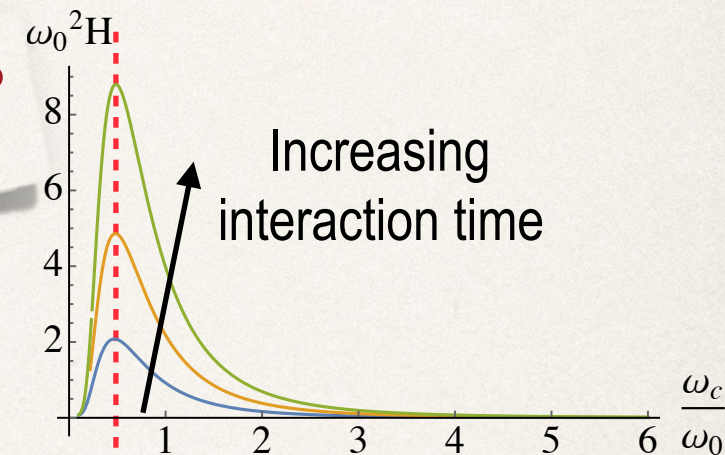
QFI : CUTOFF/PROBE FREQUENCY

Independent on reservoir s and time t !!

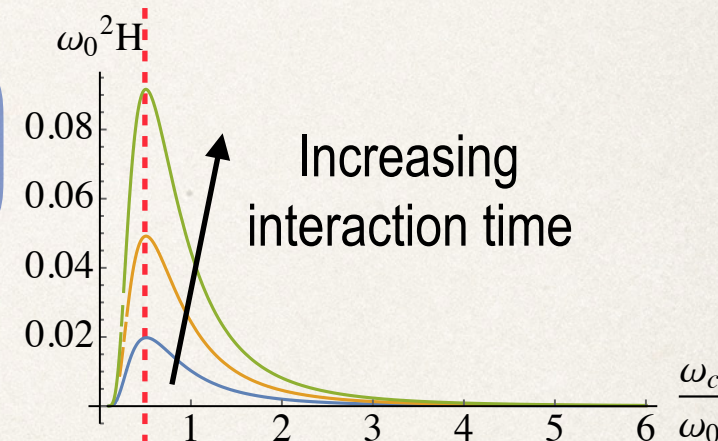
Scaling factor

$$\frac{\text{QFI}^H}{\text{QFI}^L} = f(\xi, n_{th}, T)$$

Independent on temperature T and time t !!



Increasing interaction time



Increasing interaction time

$$\left(\frac{\omega_c}{\omega_0}\right)^{(\max)} = f(s) = \begin{cases} (s \pm \sqrt{s})^{-1} & \text{for } s \neq 0, s \neq 1 \\ 1/2 & \text{for } s = 1. \end{cases}$$

$$\left(\frac{\omega_c}{\omega_0}\right)^{(\max)}$$





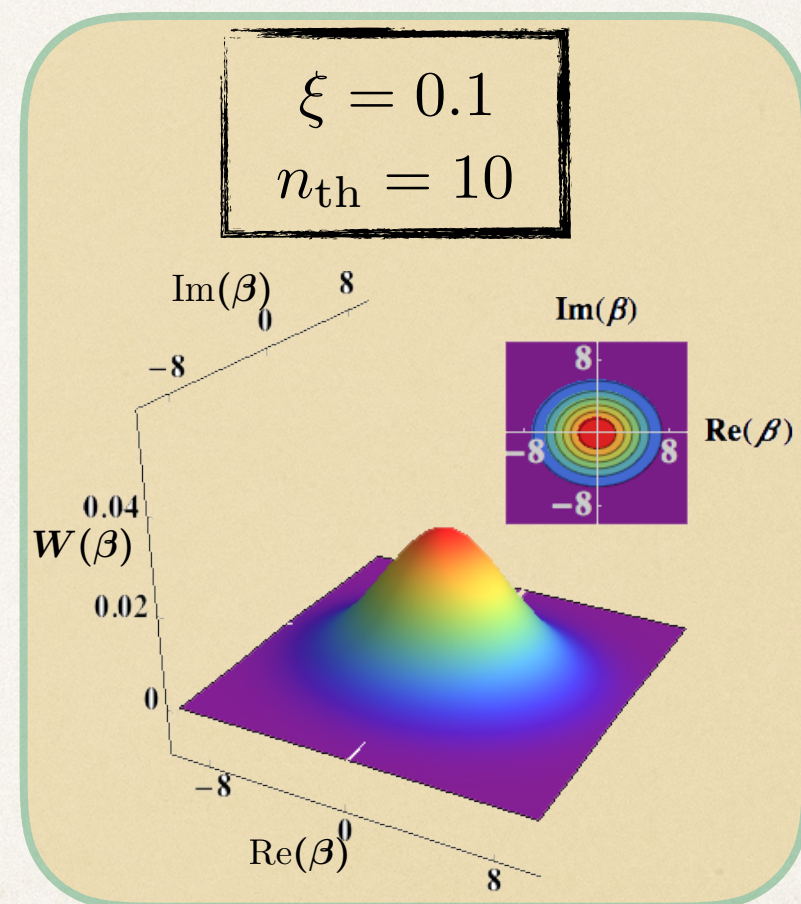
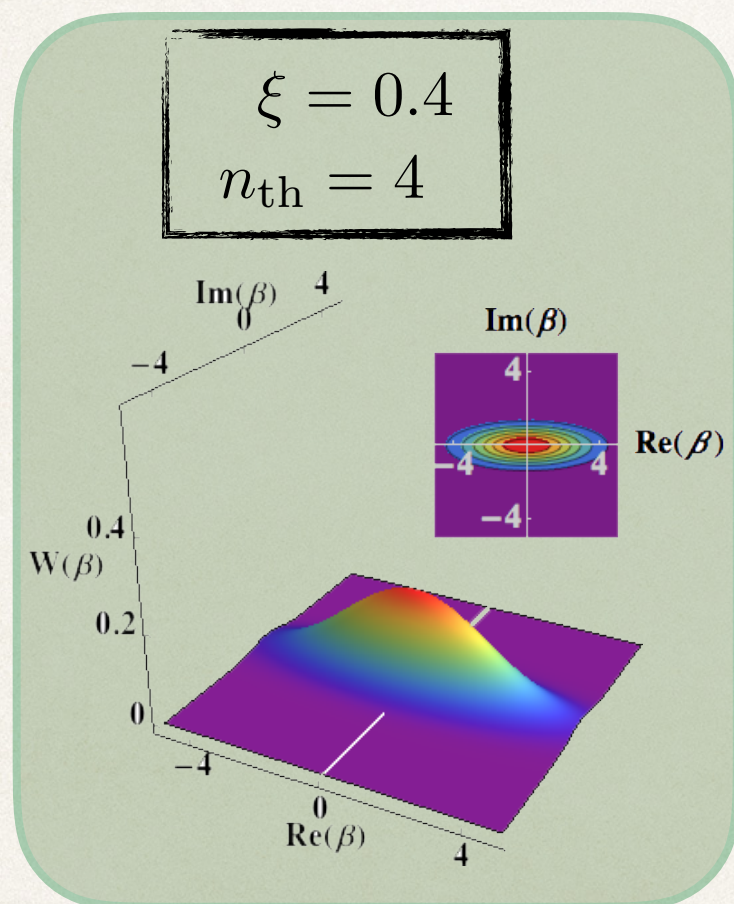
Results: QFI

QFI : INITIAL STATE PARAMETERS

ξ n_{th}

SQUEEZED THERMAL STATES :

$$\rho_0 = S(\xi)\nu(n_{\text{th}})S^\dagger(\xi)$$





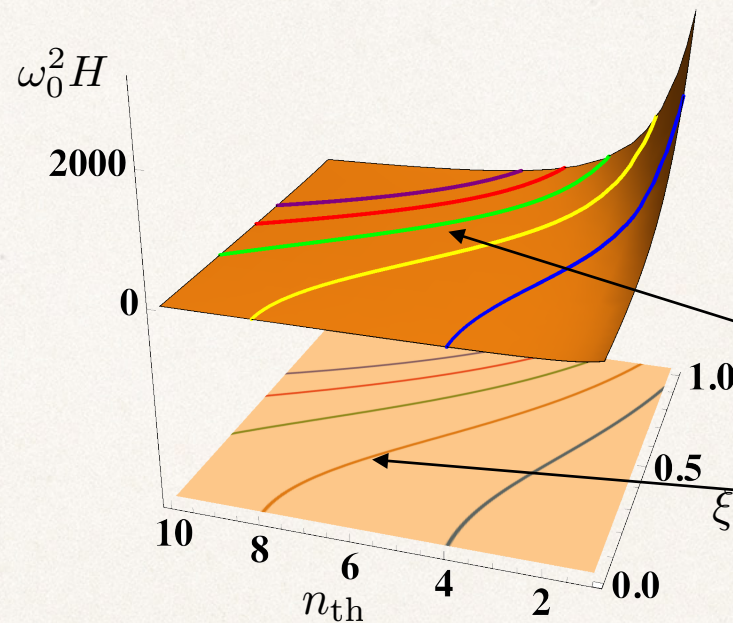
Results: QFI

$$\rho_0 = S(\xi) \nu(n_{\text{th}}) S^\dagger(\xi)$$

QFI : INITIAL STATE PARAMETERS

ξ n_{th}

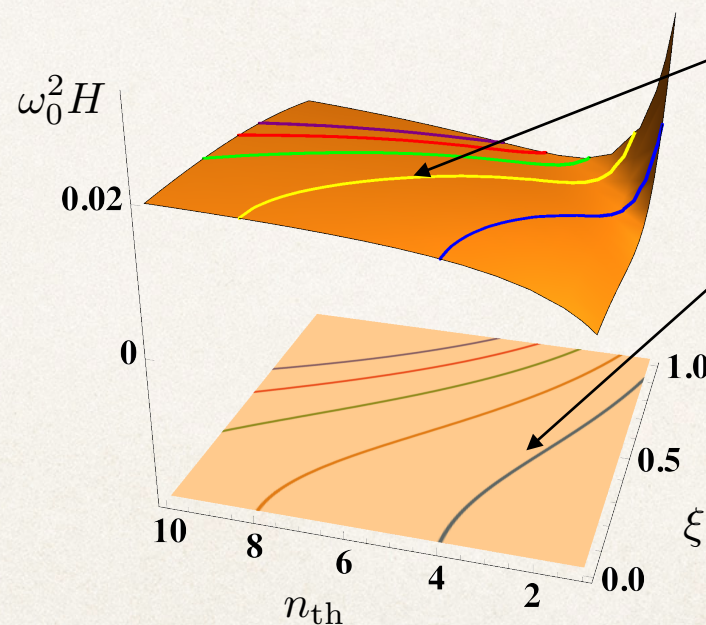
HIGH - T



TOTAL ENERGY

$$\mathcal{N} = \langle a^\dagger a \rangle_{\rho_0} = \sinh^2(\xi) + n_{\text{th}} [1 + 2 \sinh^2(\xi)]$$

LOW - T



Fixed probe total energy \mathcal{N}

QFI



ξ



n_{th}



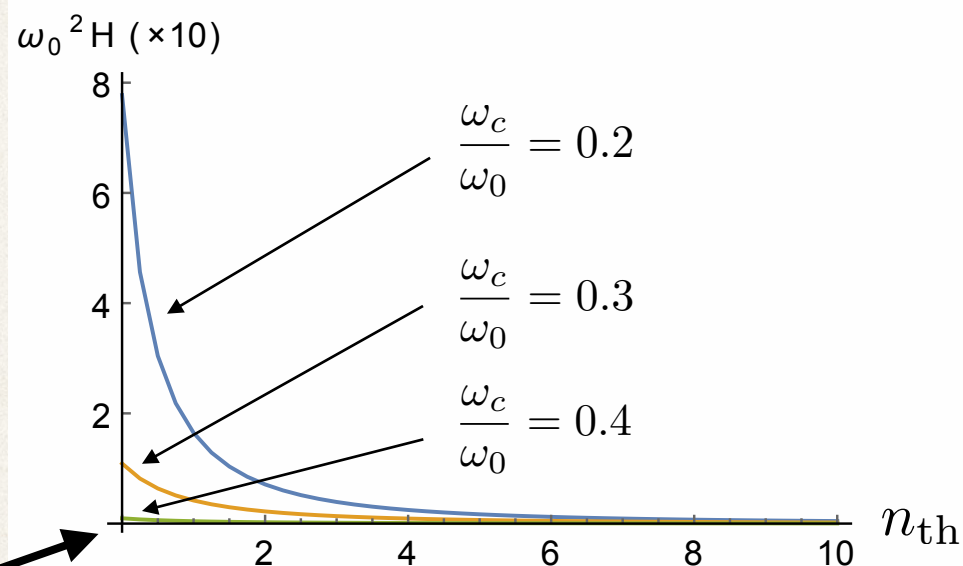
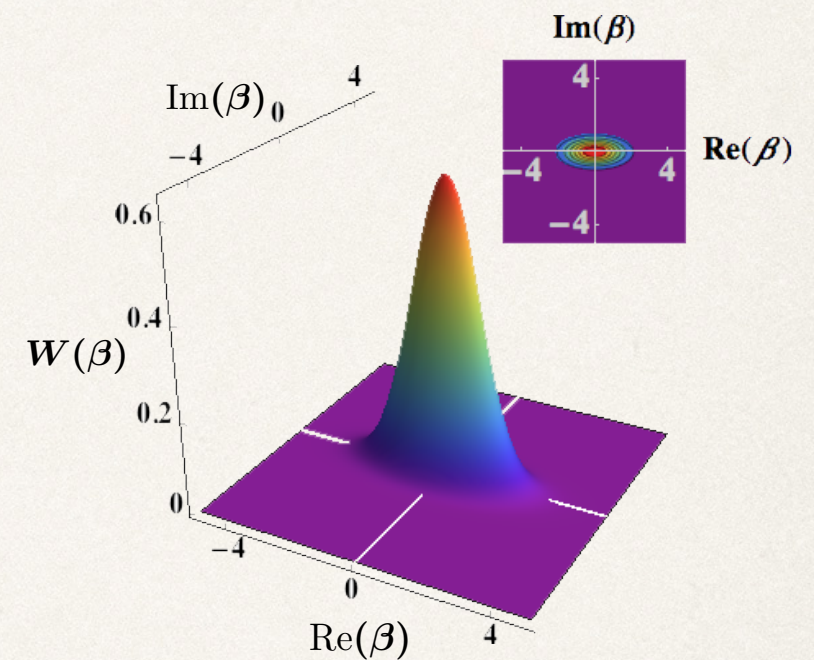
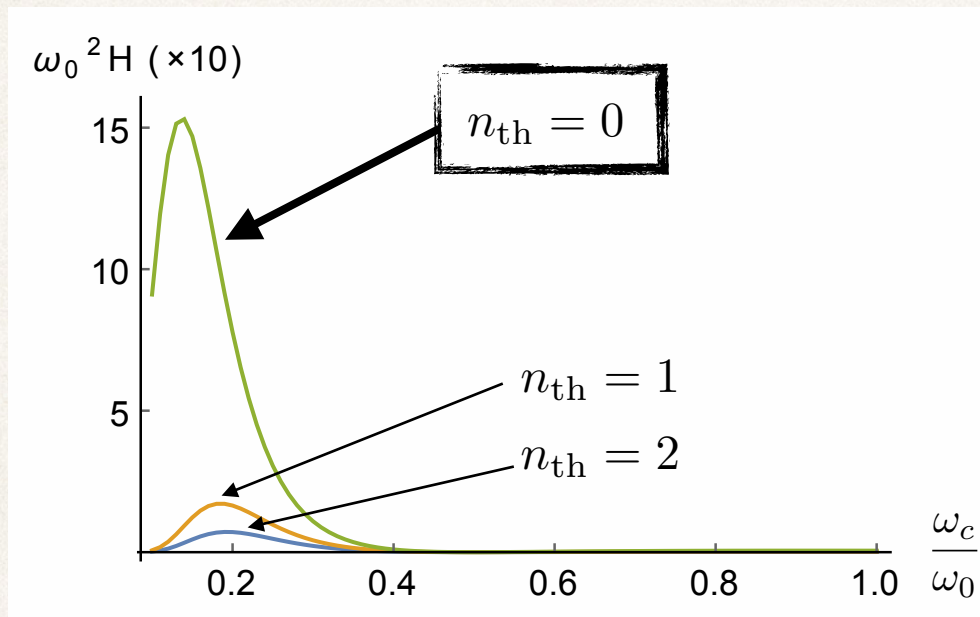


Results: QFI

$$\rho_0 = S(\xi)|0\rangle\langle 0|S^\dagger(\xi)$$

QFI : INITIAL SQUEEZED VACUUM STATE

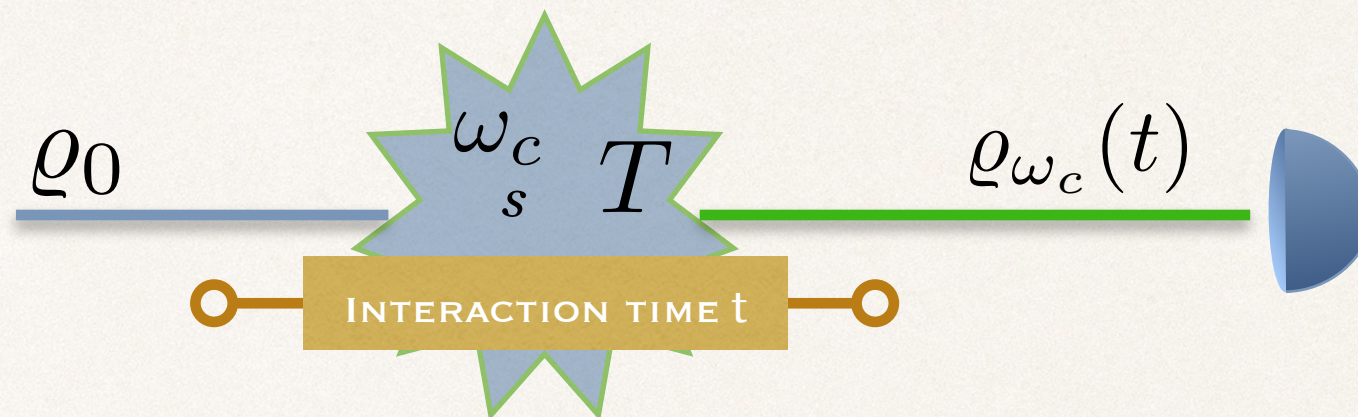
$$n_{\text{th}} = 0$$



These plots are obtained for:

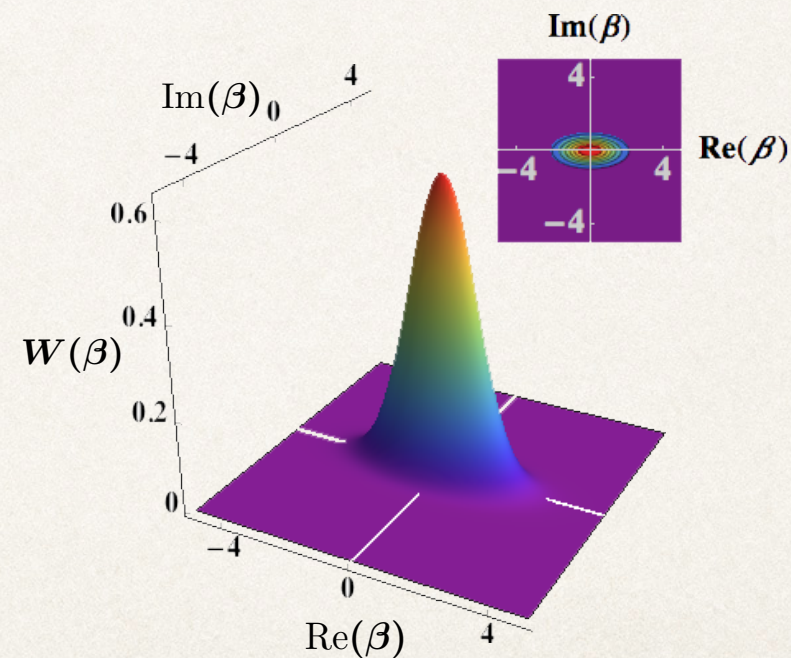
- ◆ Super-Ohmic reservoir ($s = 3$)
- ◆ Squeezing $\xi = 0.1$
- ◆ High Temperature limit

Results: QFI



QFI : BEST STRATEGY IS TO PROBE
with a single-mode **SQUEEZED VACUUM** state

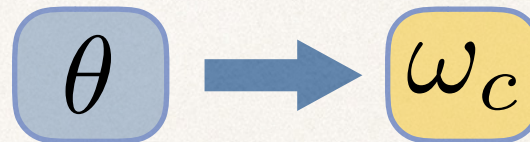
$$\rho_0 = S(\xi)|0\rangle\langle 0|S^\dagger(\xi)$$



Results: FI

FISHER INFORMATION (FI)

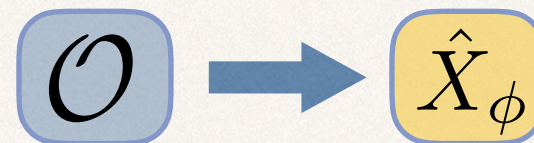
$$F_{\mathcal{O}}(\theta) = \int dx \frac{[\partial_{\theta} p(x|\theta)]^2}{p(x|\theta)}$$



QUANTUM CRAMER-RAO BOUND

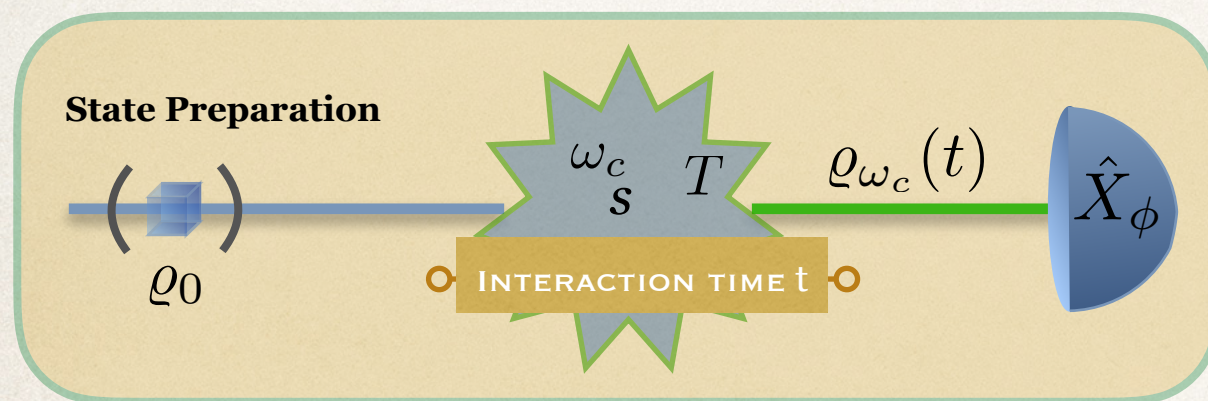
$$\text{Var}_{\theta} \geq \frac{1}{F_{\mathcal{O}}(\theta)} \geq \frac{1}{H(\theta)}$$

HOMODYNE DETECTION



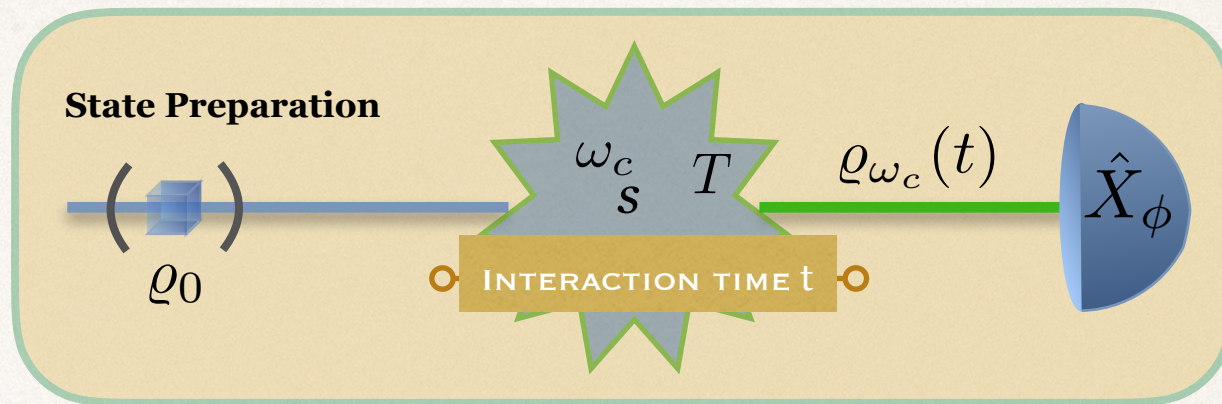
\hat{X}_0 **Position**

$\hat{X}_{\frac{\pi}{2}} = \hat{P}$ **Momentum**

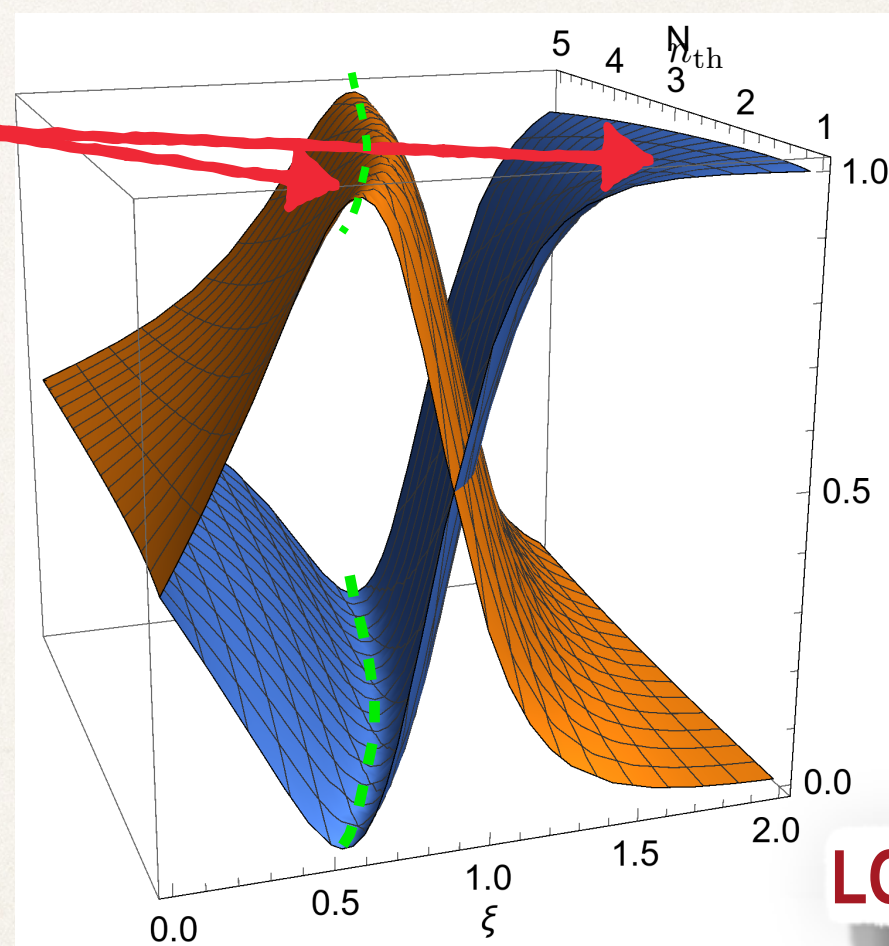
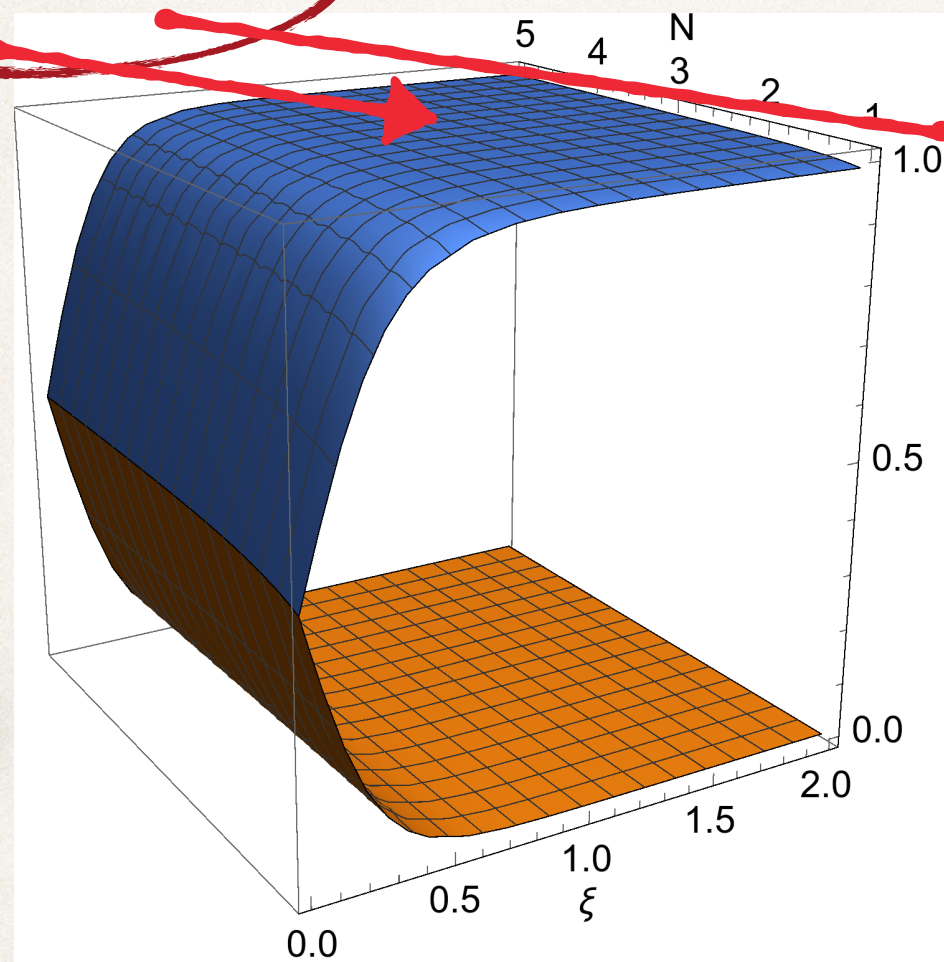


$$\hat{X}_{\phi} = \frac{e^{-i\phi} a + e^{i\phi} a^{\dagger}}{\sqrt{2}}$$

Results: FI



OPTIMAL MEASUREMENT



Orange square: $\frac{F_{\hat{X}_0}(\omega_c)}{H(\omega_c)}$

Blue square: $\frac{F_{\hat{X}_{\pi/2}}(\omega_c)}{H(\omega_c)}$

Conclusions

**THANKS FOR
YOUR ATTENTION!**

HIGHEST QFI FOR THE CUTOFF FREQUENCY?

$$\rho_0 = S(\xi) \nu(n_{\text{th}}) S^\dagger(\xi)$$



SQUEEZED VACUUM

$$\rho_0 = S(\xi) |0\rangle \langle 0| S^\dagger(\xi)$$

INVARIANTS AND SCALING FACTORS OF THE QFI?

$$\left(\frac{\omega_c}{\omega_0} \right)^{(\max)}$$

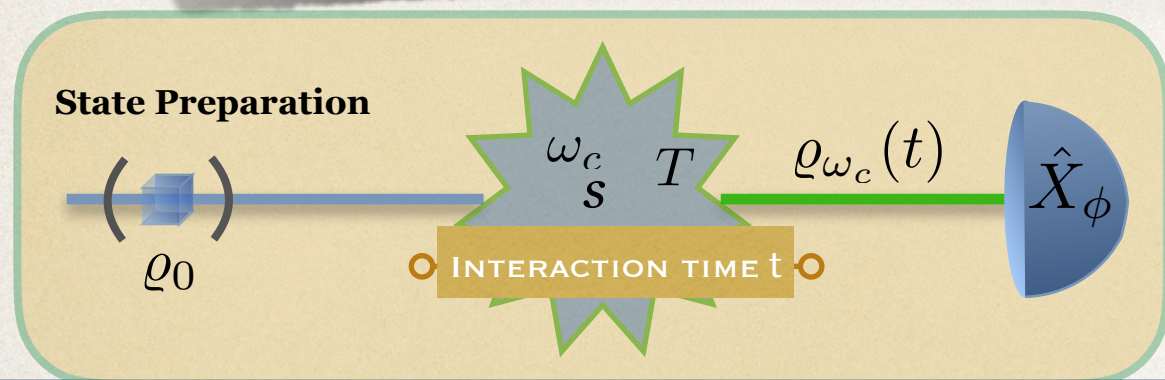
**Independent on
temperature T and time t!!**

$$\frac{\text{QFI}^H}{\text{QFI}^L} = f(\xi, n_{\text{th}}, T)$$

**Independent on reservoir s
and time t !!**

OPTIMAL MEASUREMENT SCHEMES?

HOMODYNE DETECTION



$$\text{Var}_{\omega_c} \geq \frac{1}{F_O(\omega_c)} \geq \frac{1}{H(\omega_c)}$$