

Momentum Resolved Spectroscopy Using Atomic Quantum Probes



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arXiv:1511.00833



OPEN QUANTUM SYSTEMS & ENTANGLEMENT

Outline



What to probe?

Outline



What to probe?



How is it measured so far?

Outline



What to probe?



How is it measured so far?



New method

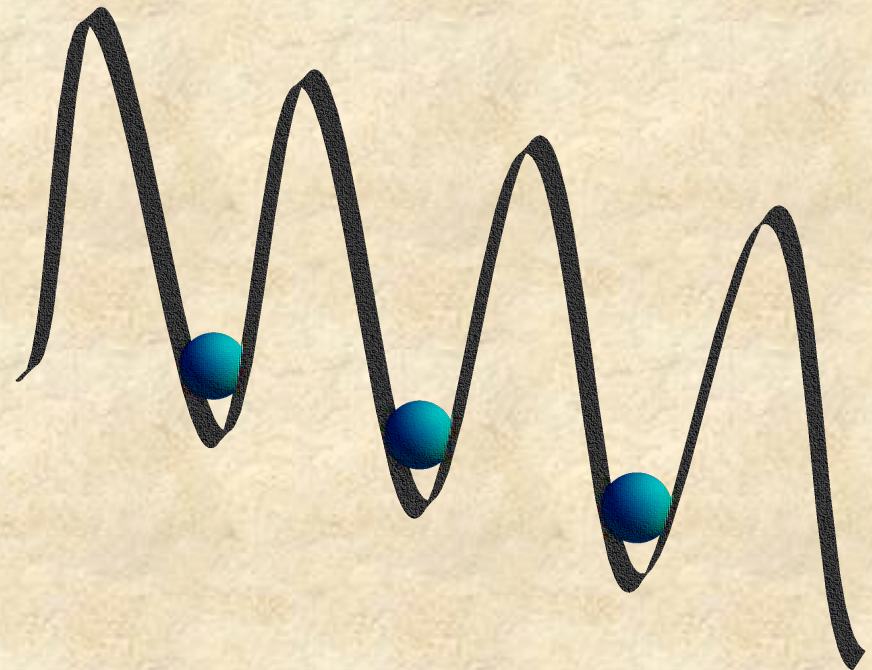
Bosons on optical lattice (1D)

Bose Hubbard

$$\hat{H} = -J \sum_{\langle i,j \rangle} \hat{c}_i^\dagger \hat{c}_j + \frac{U}{2} \sum_i \hat{c}_i^\dagger \hat{c}_i^\dagger \hat{c}_i \hat{c}_i - \mu \sum_i \hat{c}_i^\dagger \hat{c}_i$$

Superfluid Phase

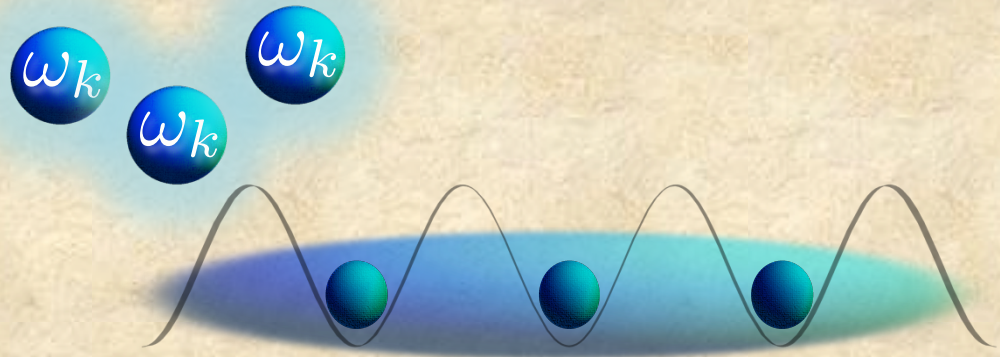
$$\hat{H} = \sum_k \omega_k \hat{b}_k^\dagger \hat{b}_k$$



Superfluid

Bogoliubov modes

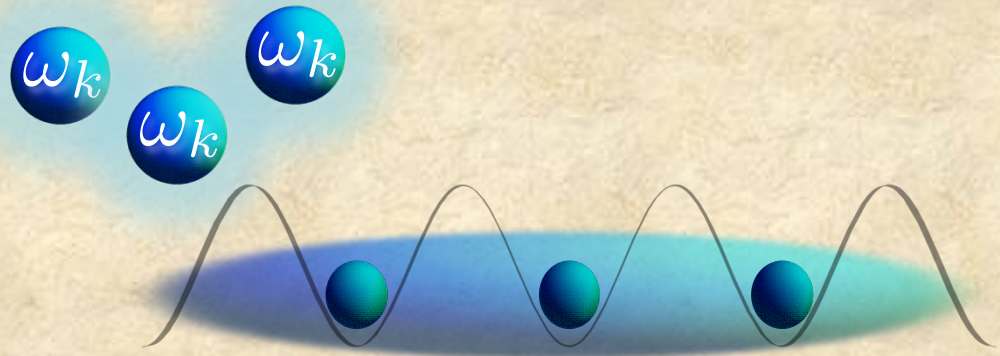
$$\begin{cases} \hat{b}_{\mathbf{k}} = u_{\mathbf{k}} \hat{c}_{\mathbf{k}} - v_{\mathbf{k}} \hat{c}_{-\mathbf{k}}^{\dagger} \\ \hat{b}_{\mathbf{k}}^{\dagger} = u_{\mathbf{k}} \hat{c}_{\mathbf{k}}^{\dagger} - v_{\mathbf{k}} \hat{c}_{-\mathbf{k}} \end{cases}$$



Superfluid

Bogoliubov modes

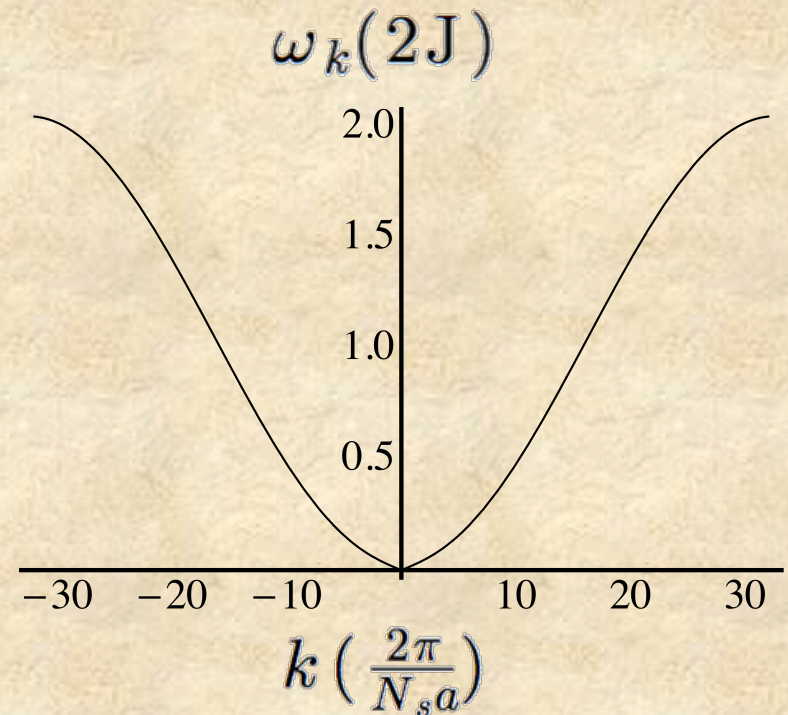
$$\begin{cases} \hat{b}_{\mathbf{k}} = u_{\mathbf{k}} \hat{c}_{\mathbf{k}} - v_{\mathbf{k}} \hat{c}_{-\mathbf{k}}^\dagger \\ \hat{b}_{\mathbf{k}}^\dagger = u_{\mathbf{k}} \hat{c}_{\mathbf{k}}^\dagger - v_{\mathbf{k}} \hat{c}_{-\mathbf{k}} \end{cases}$$



Spectrum of Excitations

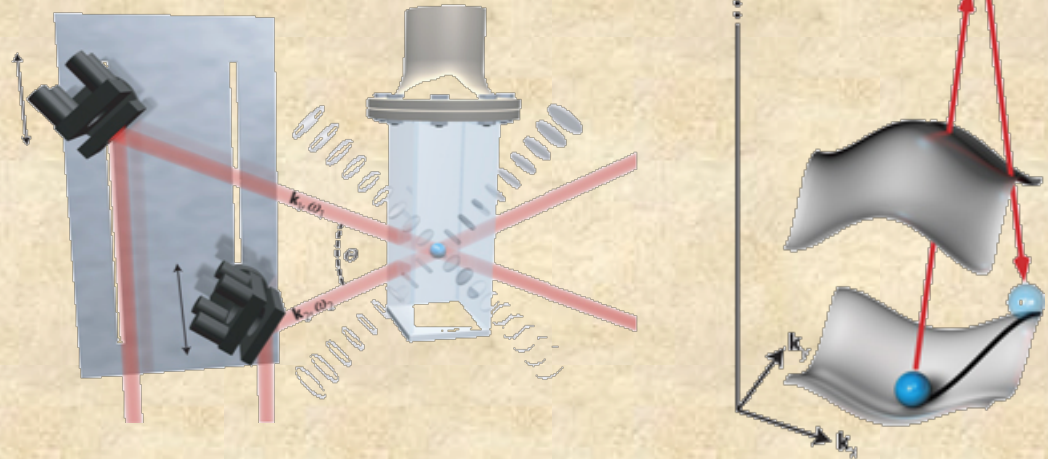
$$\omega_k = \sqrt{\epsilon_k^2 + 2U n_0 \epsilon_k}$$

$$\epsilon_k = 2J(1 - \cos(ka))$$

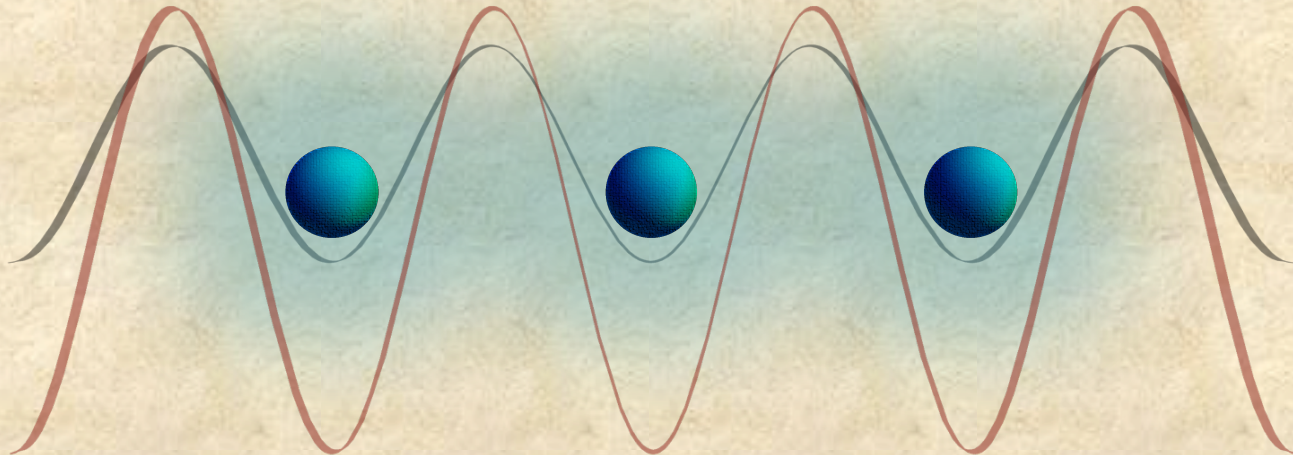


Experimental techniques

Bragg Spectroscopy¹



Modulation of the depth of the optical lattice²



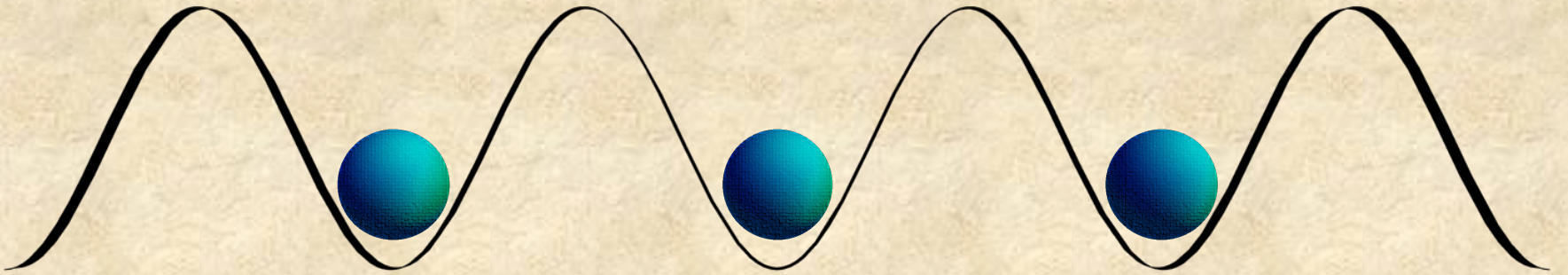
1) P. T. Ernst et al., Nature Physics **6**, 56 - 61 (2010)

2) C. Schori et al., Phys. Rev. Lett **93**, 240402 (2004)

Experimental techniques

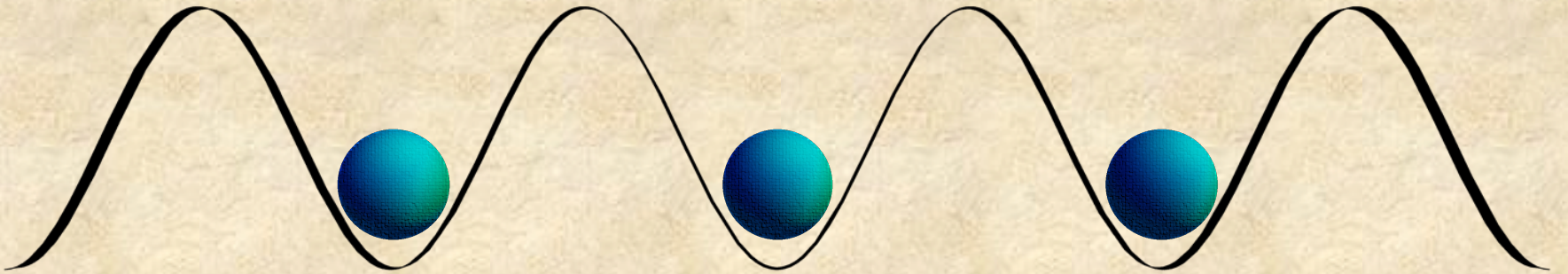
Effective but destructive!

New Probing Scheme



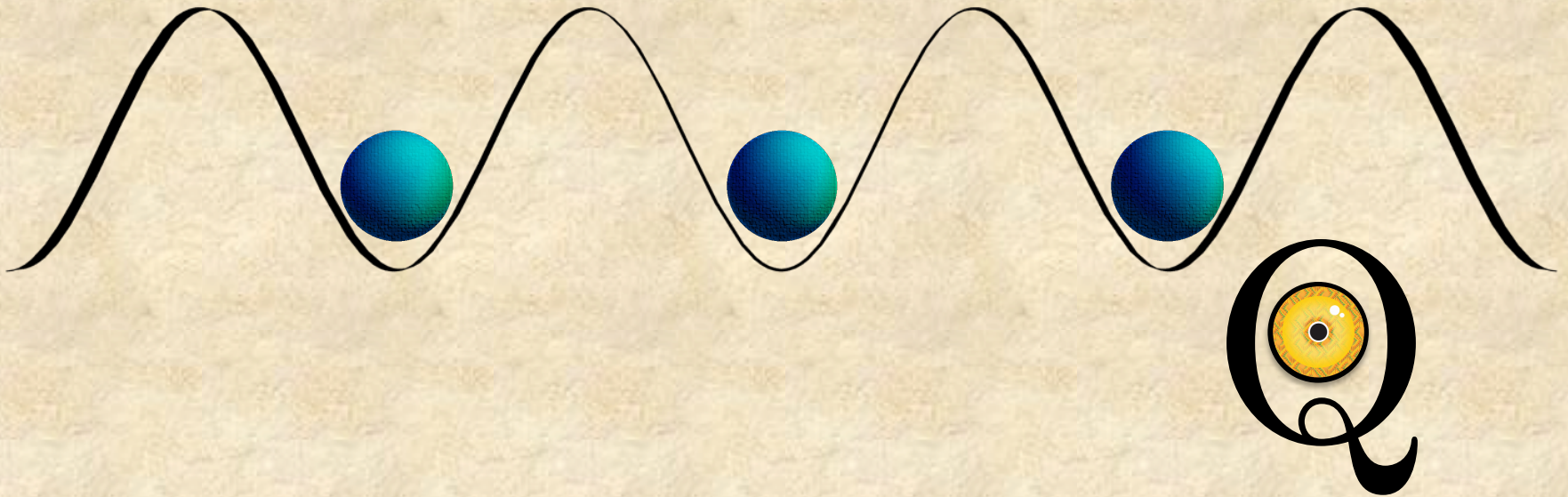
New Probing Scheme

Less invasive



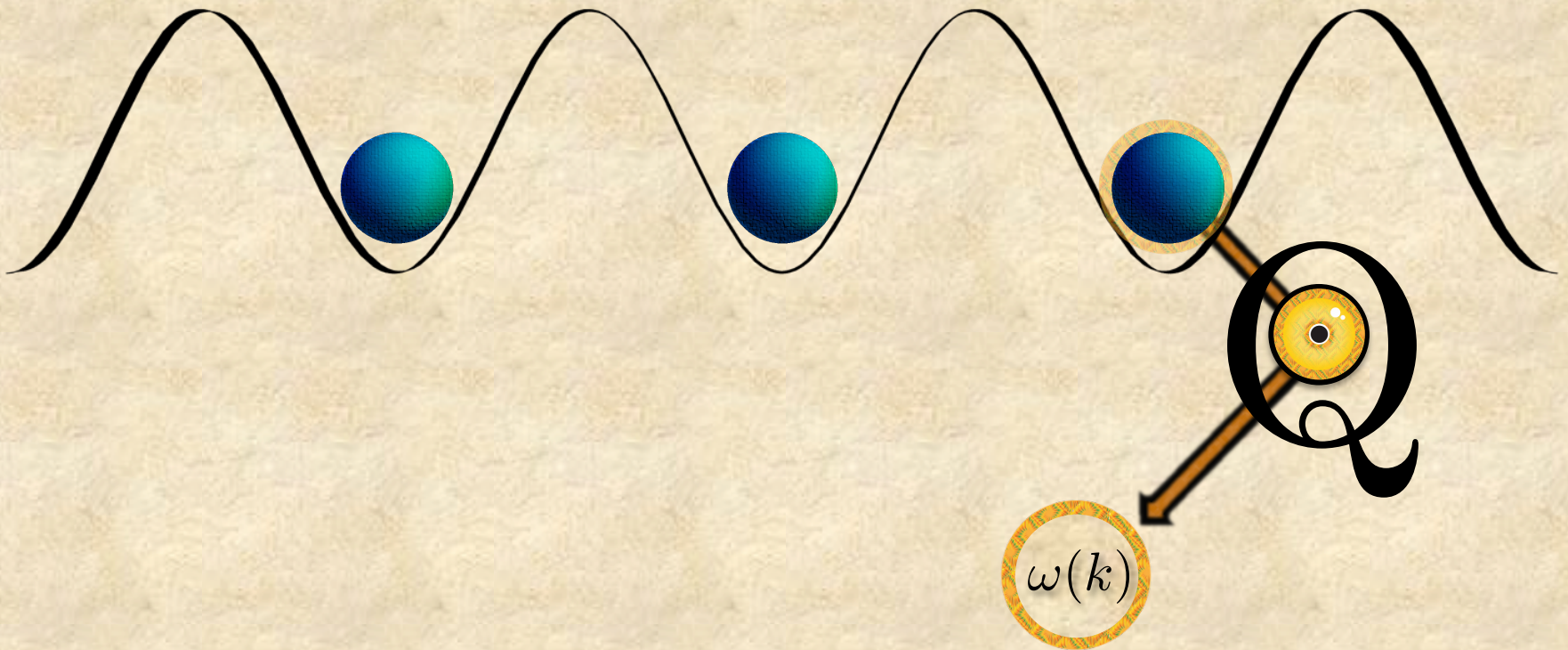
New Probing Scheme

Less invasive



New Probing Scheme

Less invasive



New Probing Scheme

Less invasive



Weak interaction

Locality



Probe

$$\hat{H}_p = \sum_{\bar{n}} \nu_{\bar{n}} |\bar{n}\rangle \langle \bar{n}|$$

Probe

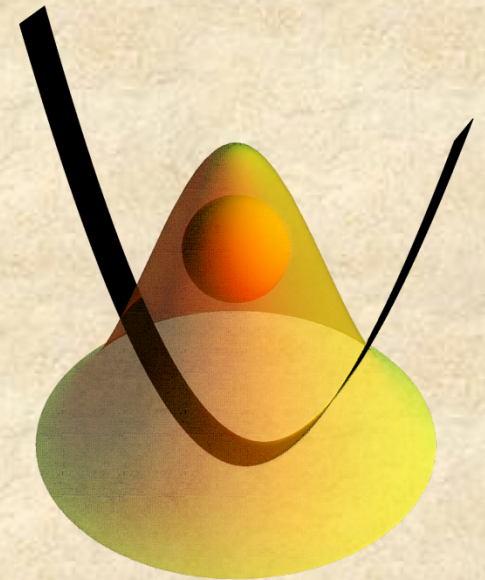
$$\hat{H}_p = \sum_{\bar{n}} \nu_{\bar{n}} |\bar{n}\rangle \langle \bar{n}|$$

Impurity atom in a 3D Harmonic trap

$$\psi_0(x) \simeq e^{-\frac{x^2}{2x_0^2}} e^{-\frac{y^2}{2y_0^2}} e^{-\frac{z^2}{2z_0^2}}$$

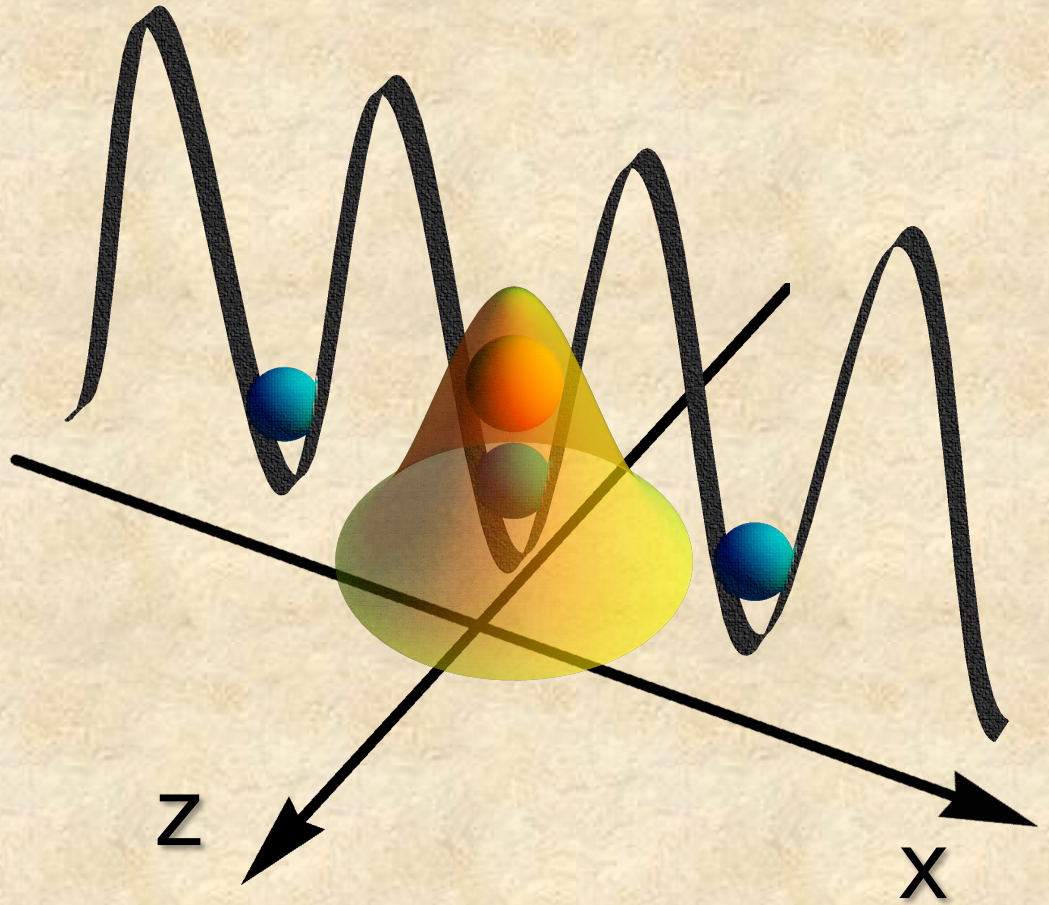


Controllable (ν_x, ν_y, ν_z)



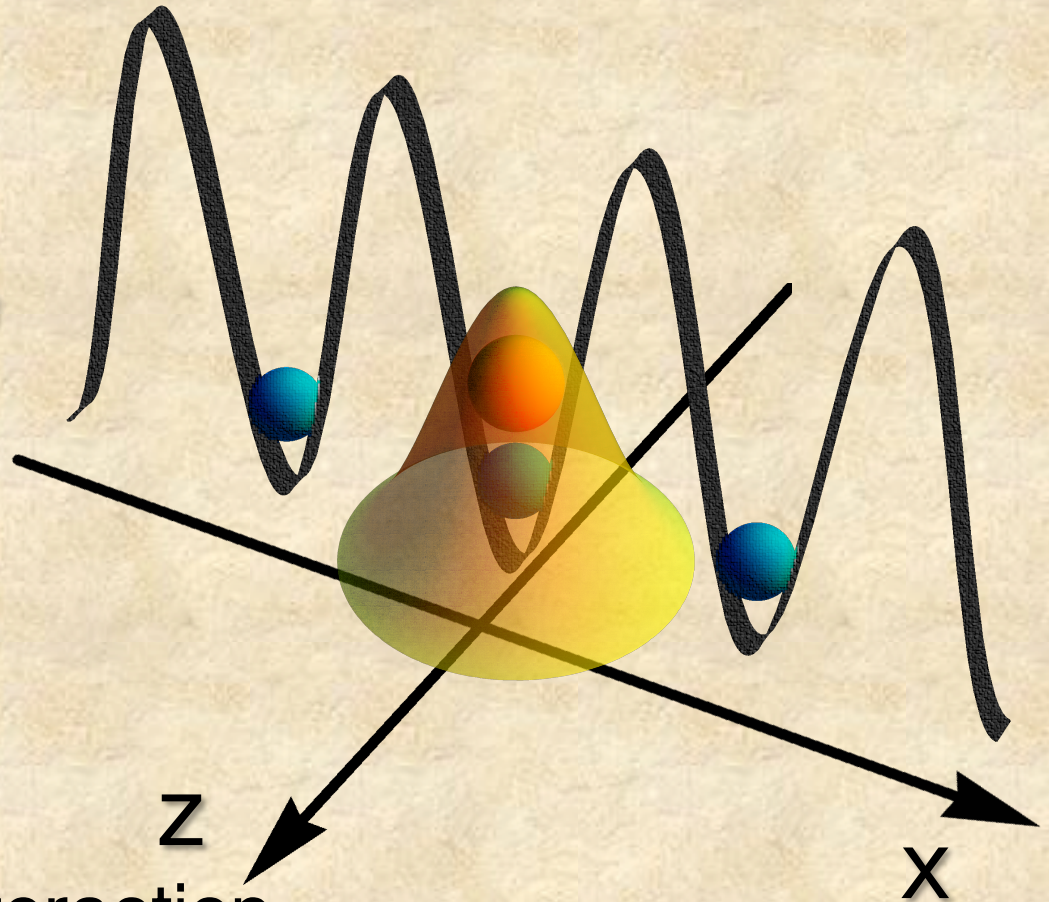
Step (I): Preparation

$$\rho(0) = |\bar{0}\rangle \langle \bar{0}| \otimes \rho_\beta$$



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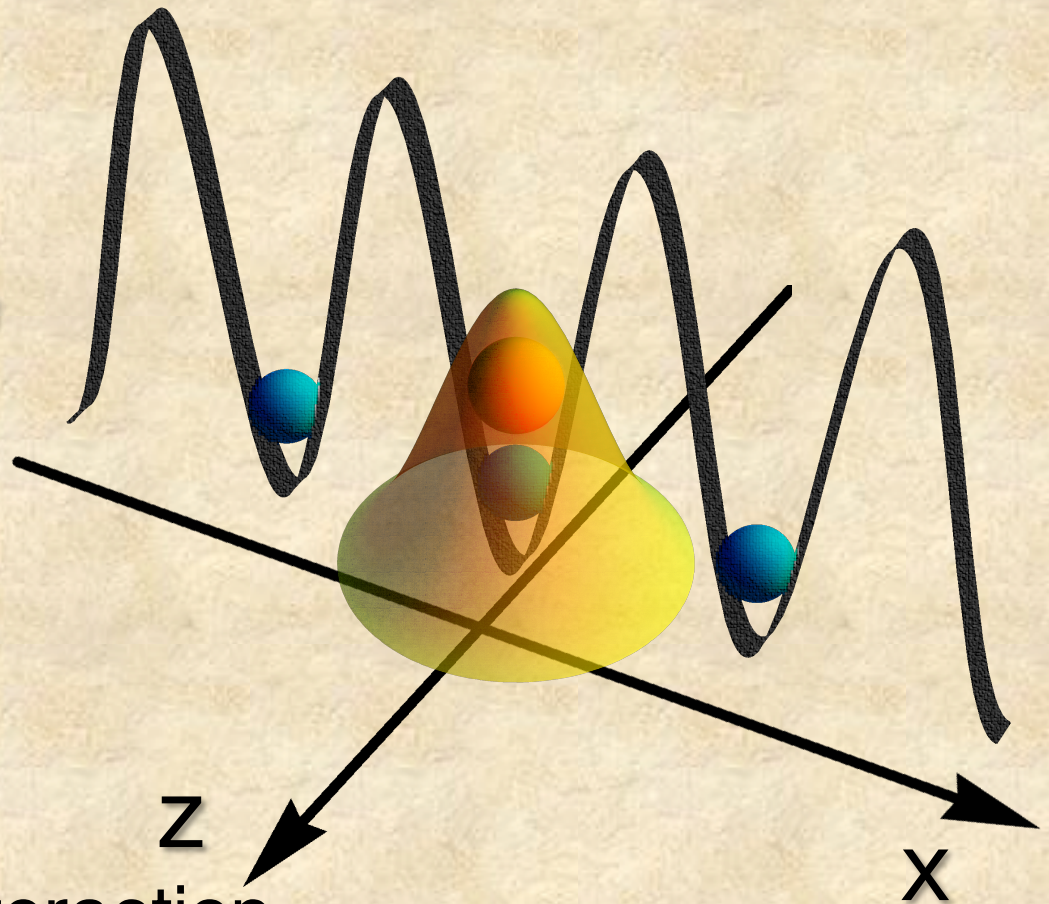


Density-density interaction

$$\hat{H}_{int} = g \sum_{\substack{\bar{n}, \bar{m} \\ i, j}} \int d\mathbf{x} \, \psi_{\bar{n}}^*(\mathbf{x}) \psi_{\bar{m}}(\mathbf{x}) \omega_i^*(x) \omega_j(x) |\bar{n}\rangle \langle \bar{m}| \otimes \hat{c}_i^\dagger \hat{c}_j$$

Step (I): Preparation

$$\rho(0) = |\bar{0}\rangle \langle \bar{0}| \otimes \rho_\beta$$



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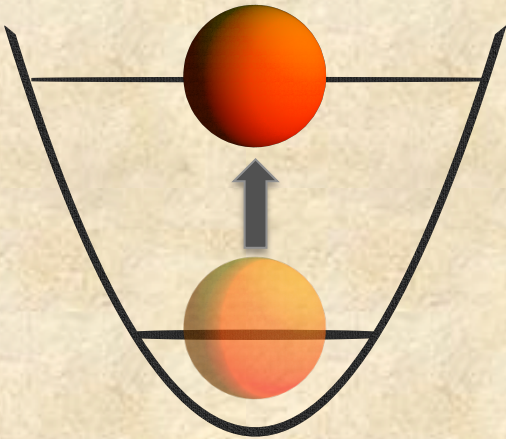
Superfluid

$$\hat{c}_i^\dagger \hat{c}_j \simeq n_0 + \sqrt{\frac{n_0}{N_s}} \sum_{k \neq 0} \left[\left(u_k e^{ikx_i} + v_k e^{ikx_j} \right) \hat{b}_k^\dagger + \left(u_k e^{-ikx_j} + v_k e^{-ikx_i} \right) \hat{b}_k \right]$$

Step (II):

Measure the transition probability

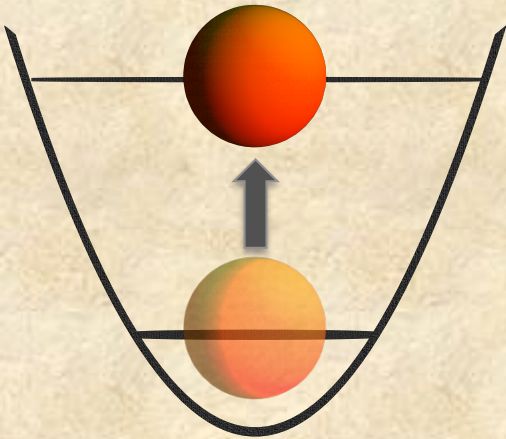
$$\Gamma_{\bar{0} \rightarrow (0,0,n_z)}^{\mathbf{I}} = \text{Tr}_B \langle 0, 0, n_z | \rho(t) | 0, 0, n_z \rangle$$



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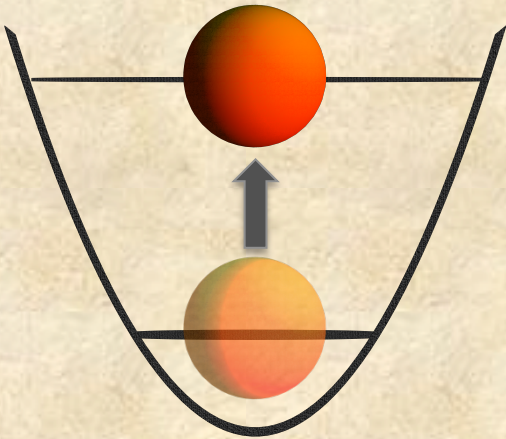
Perturbative approach

$$\hat{U}(t) = \hat{\mathbb{I}} + g\hat{U}^{(1)}(t) + g^2\hat{U}^{(2)}(t) + \dots$$

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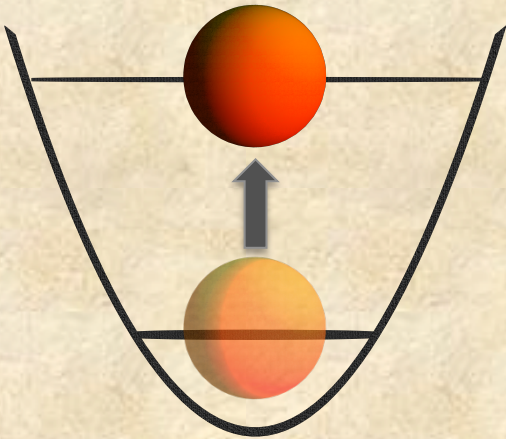
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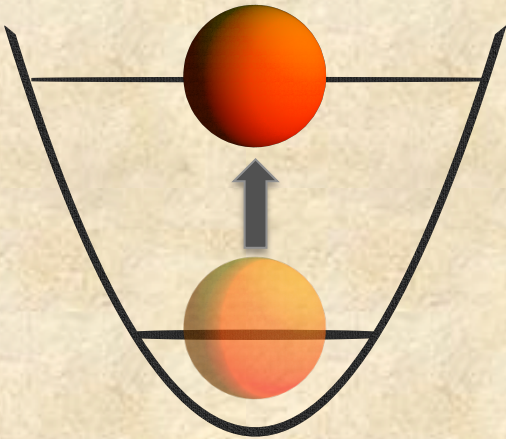
Relevant

$$\langle 0, 0, n_z | \hat{H}_{int} | \bar{0} \rangle = g \sum_{i,j} \int d\mathbf{x} \psi_{(0,0,n_z)}^*(\mathbf{x}) \psi_{(0,0,0)}(\mathbf{x}) \omega_i^*(x) \omega_j(x) \hat{c}_i^\dagger \hat{c}_j$$

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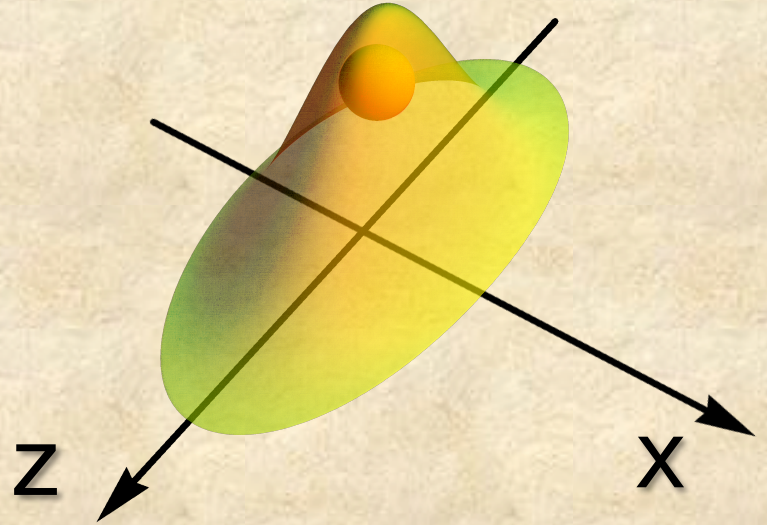
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Relevant

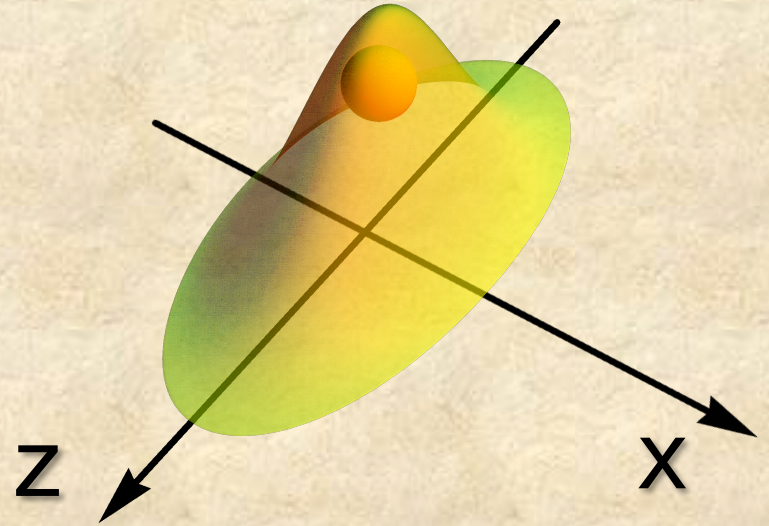
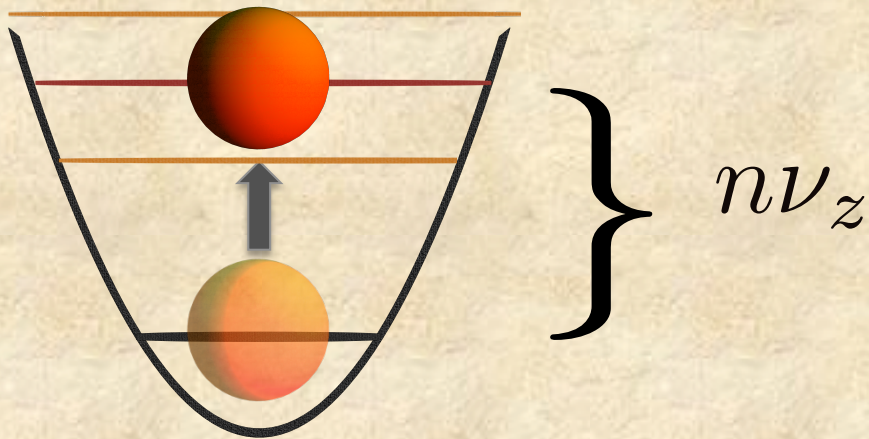
$$\langle 0, 0, n_z | \hat{H}_{int} | \bar{0} \rangle \simeq g \int d\mathbf{x} \psi_{(0,0,n_z)}^*(\mathbf{x}) \psi_{(0,0,0)}(\mathbf{x}) \omega_0^*(x) \omega_0(x) \hat{c}_0^\dagger \hat{c}_0$$

Step (III):
Vary ν_z

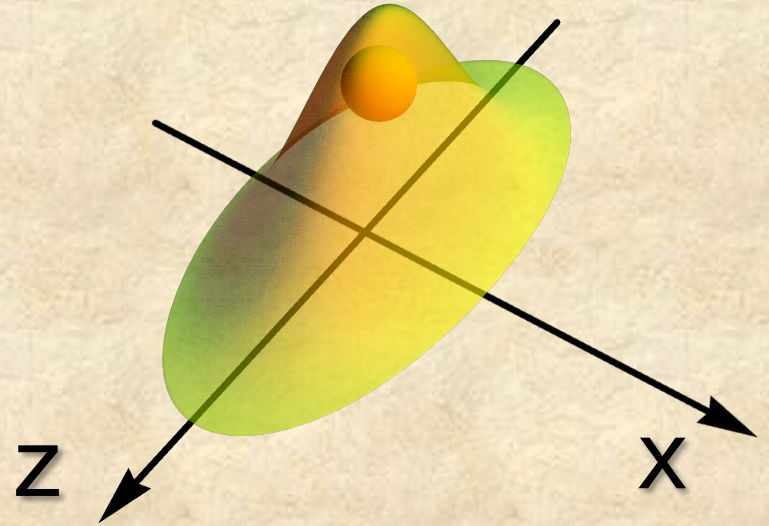


Step (III):
Vary ν_z

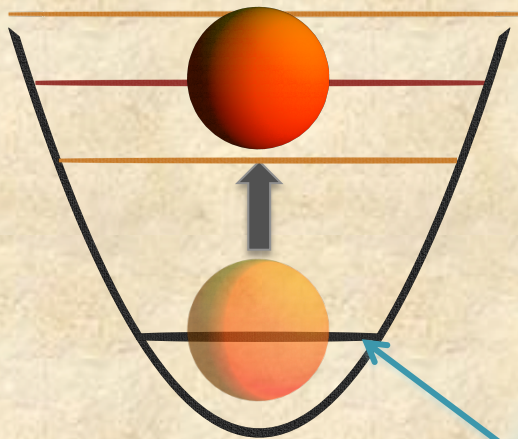
Probe Energy levels



Step (III):
Vary ν_z

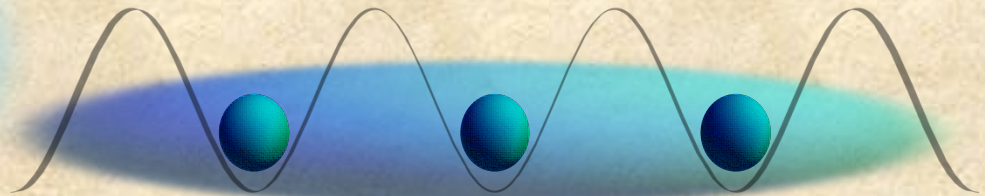


Probe Energy levels



$$n\nu_z = \omega_k$$

Resonance!

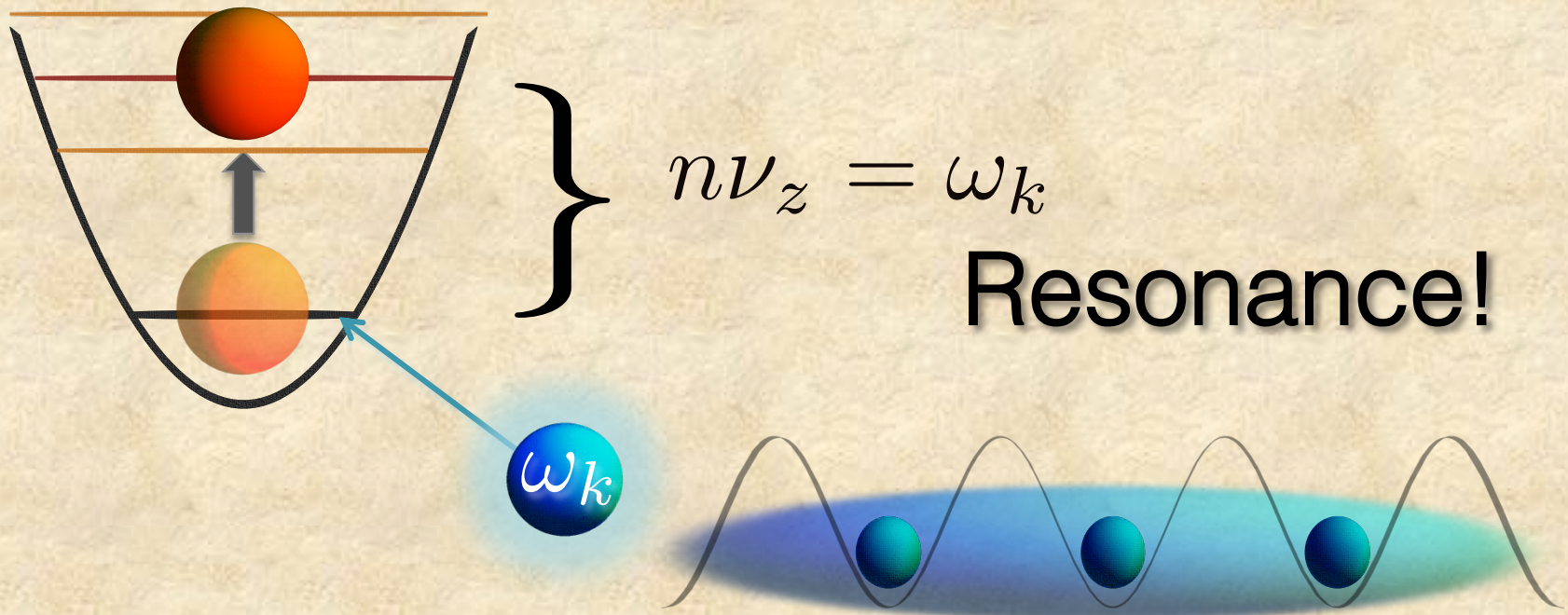


Transition Probability

$$\Gamma_{\vec{0} \rightarrow (0,0,n_z)}^{\mathbf{I}} = g_{\mathbf{I},n}^2 \nu \left[2n_0^2 \frac{1 - \cos(n\nu t)}{n^2 \nu^2} + \sum_k (\Gamma_k^+(n\nu, t) + \Gamma_k^-(n\nu, t)) \right]$$

$$\Gamma_k^{\pm}(n\nu, t) = 2\beta_k^2 \frac{1 - \cos[(n\nu \mp \omega_k)t]}{(n\nu \mp \omega_k)^2} \left[\frac{1 \mp 1}{2} + n(\omega_k) \right]$$

Probe Energy levels



Transition Probability

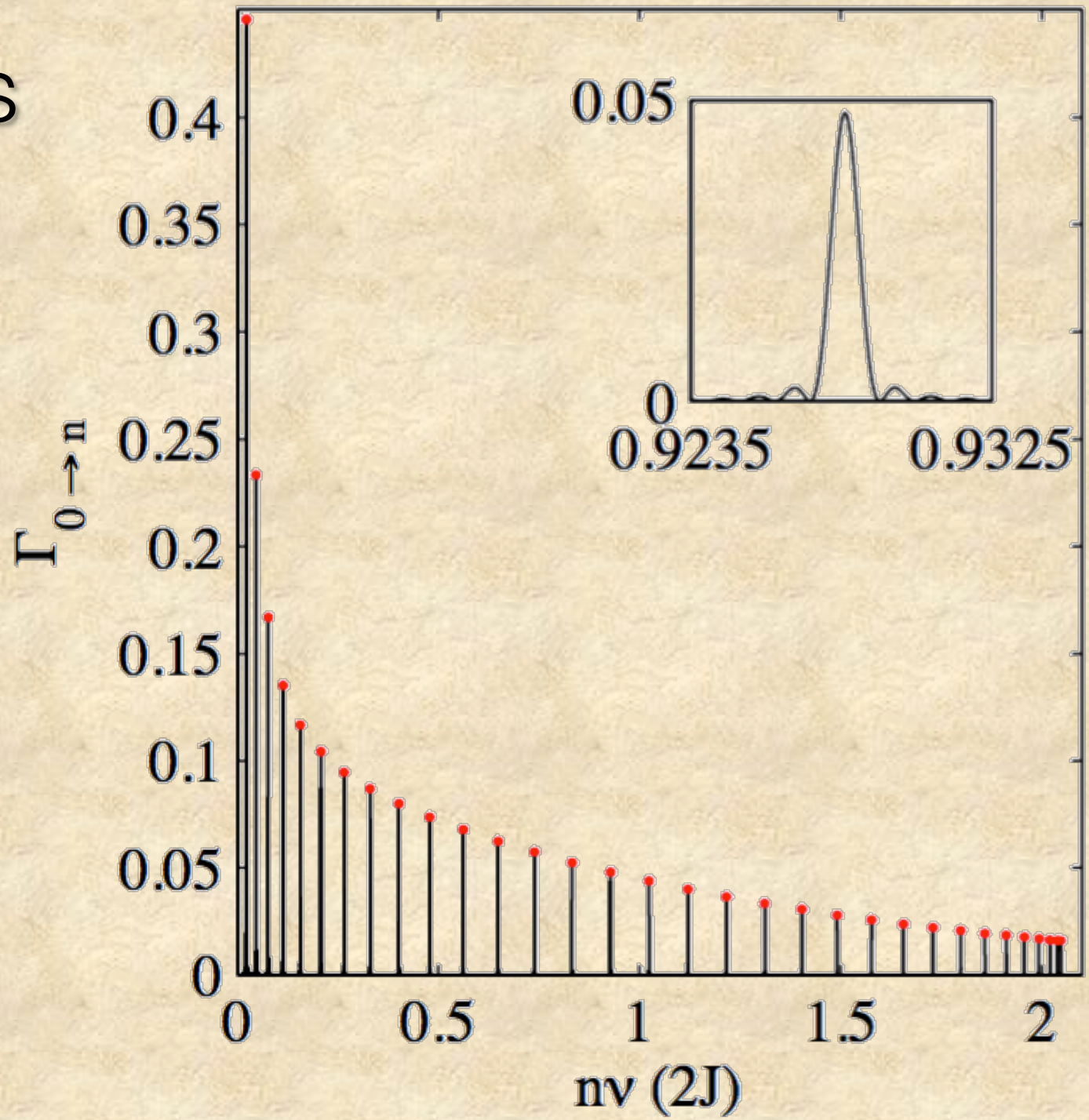
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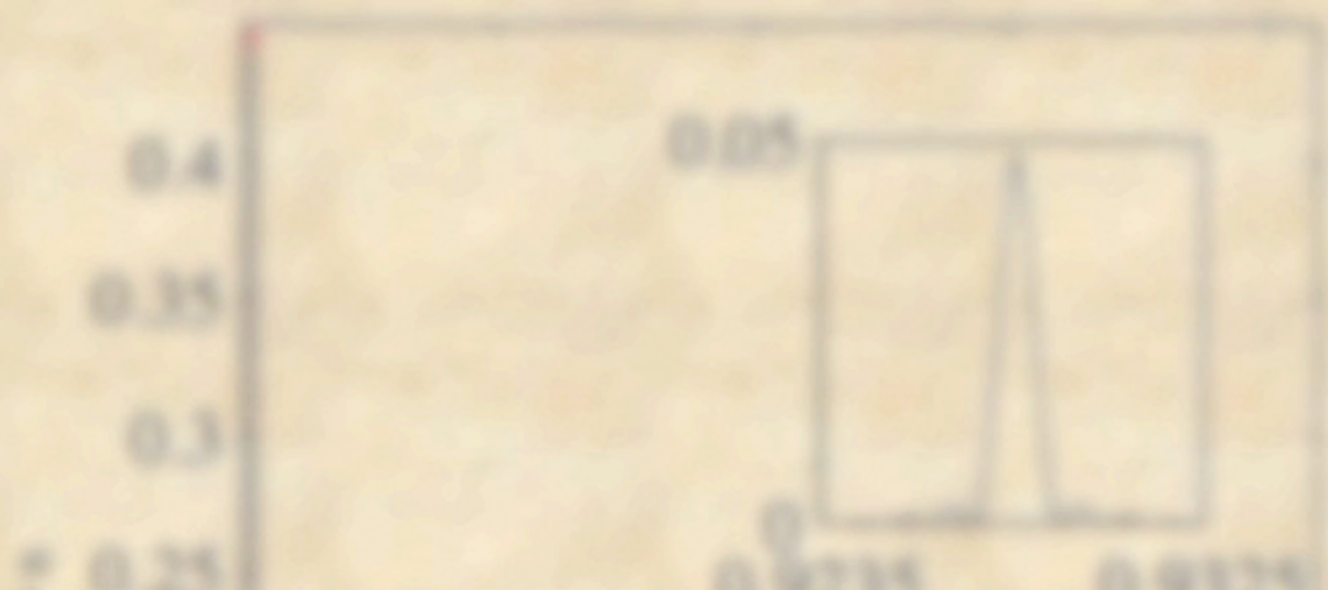
Tunability

$$g_{\mathbf{I},n}^2 = \frac{g^2}{\nu} \underbrace{\left(\int dx \psi_{n_x=0}^2(x) \omega_0^2(x) \right)^2}_{\mathbf{X}} \underbrace{\left(\sqrt{m} \frac{\gamma_0}{\pi} \sqrt{\nu_y} \right)^2}_{\mathbf{Y}} \underbrace{\left((-1)^{n_z} \sqrt{m} \frac{\gamma_{n_z}^{1/2} \gamma_0^{1/2}}{\pi} \sqrt{\nu} \right)^2}_{\mathbf{Z}}$$

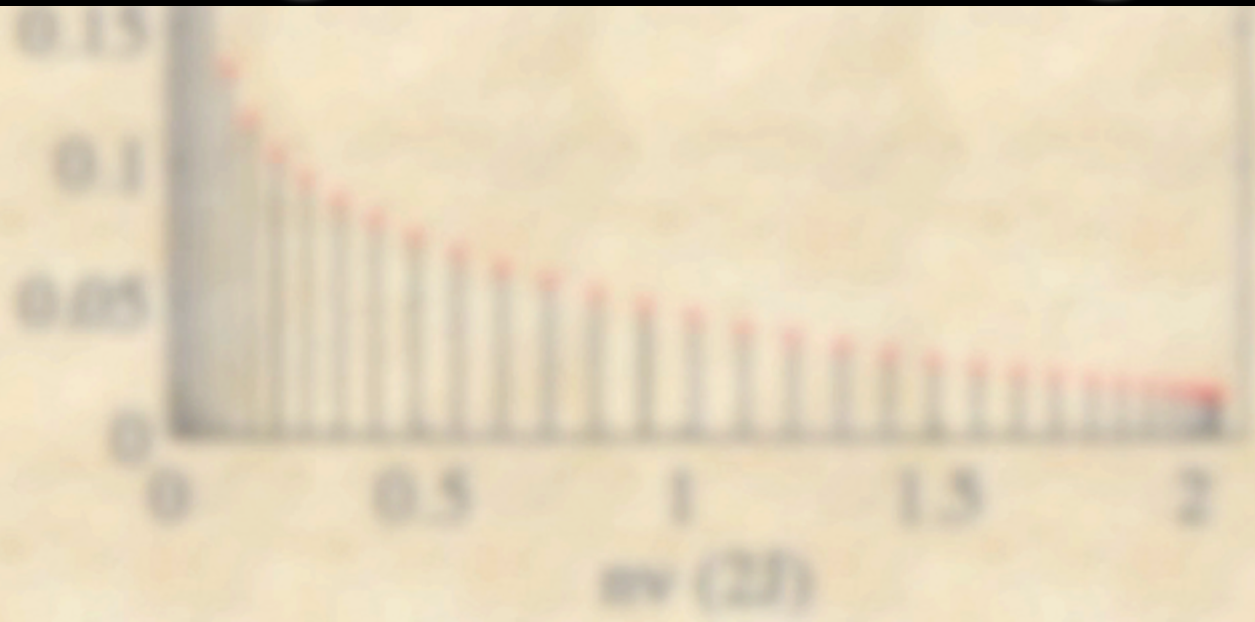
Results



Results



Something is missing



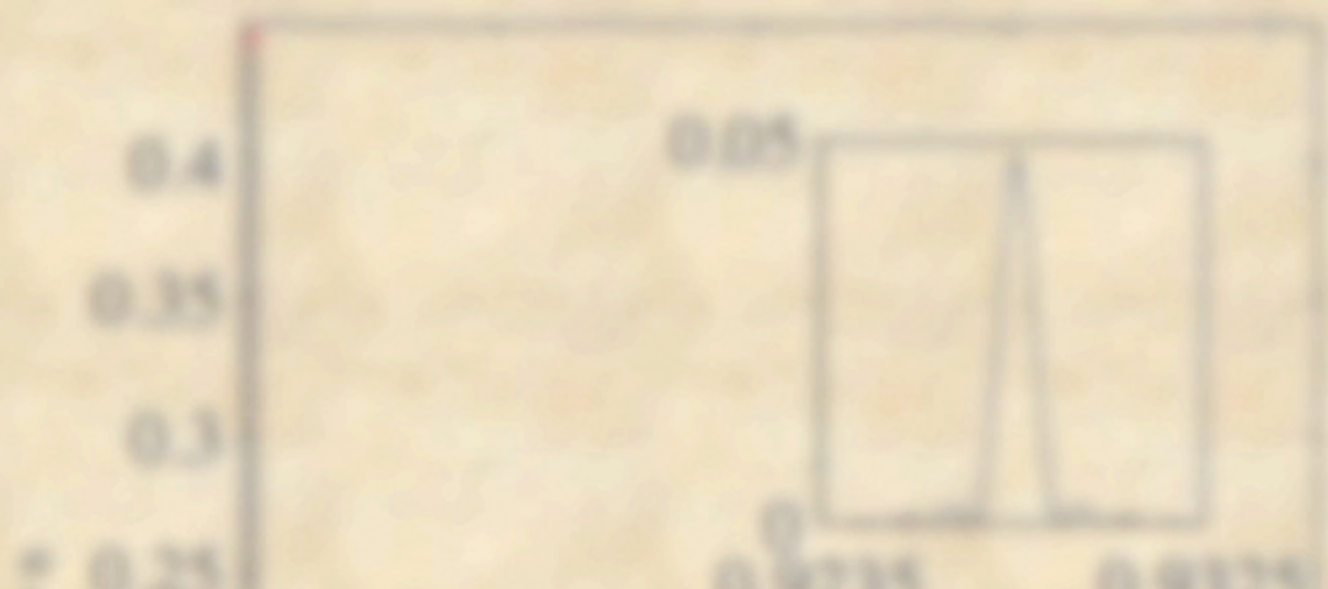
Results

Dispersion law $\omega(k)$

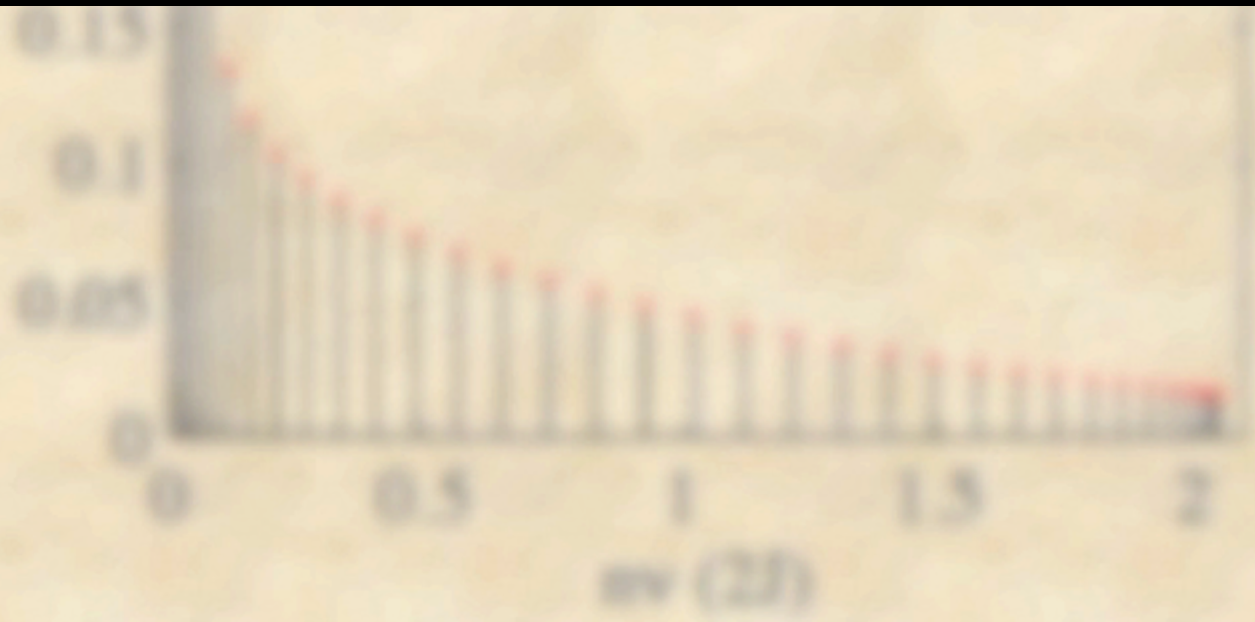


The background image contains two plots. The top plot is a line graph with a peak, and the bottom plot is a bar chart with a decaying trend.

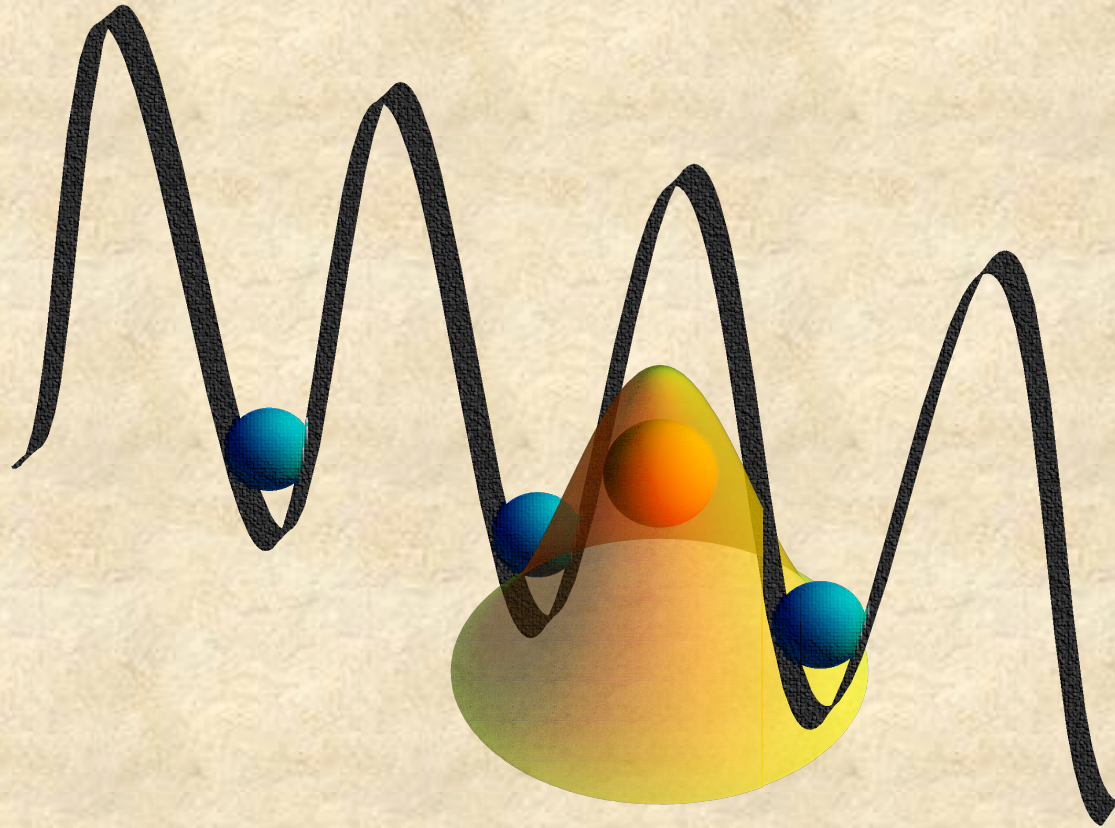
Results



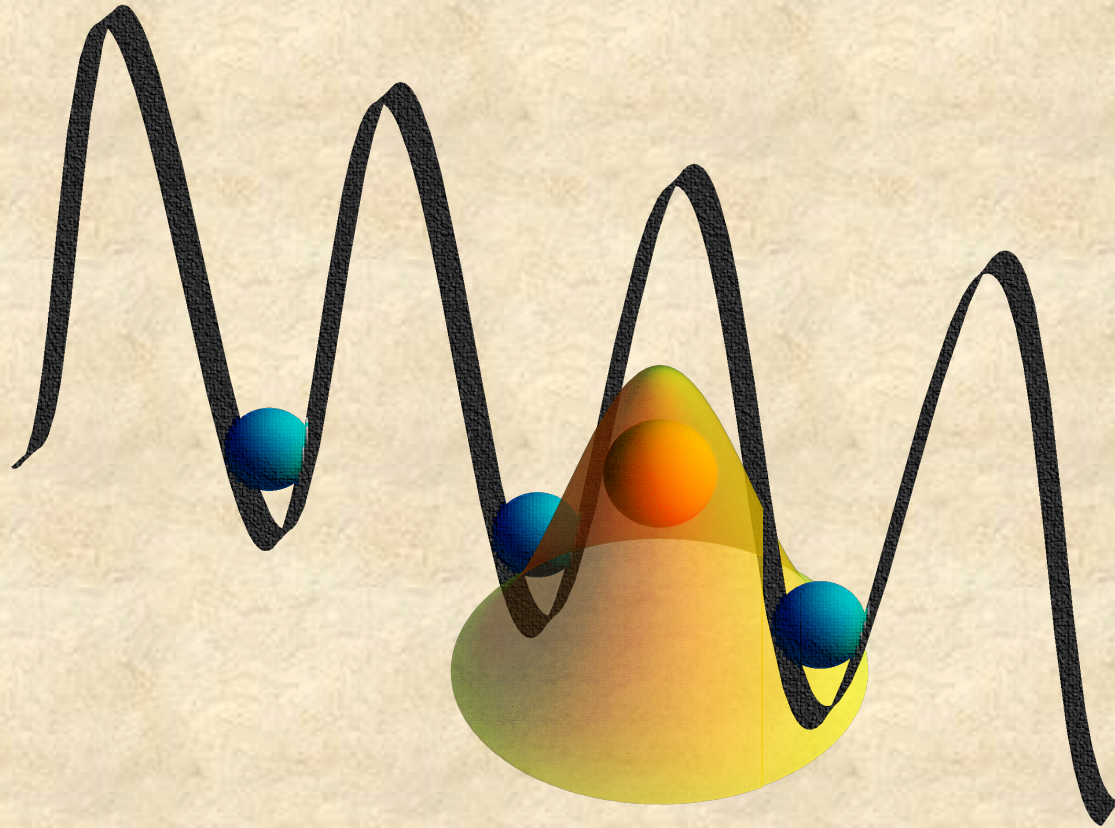
Can we solve it?



Step (IV):



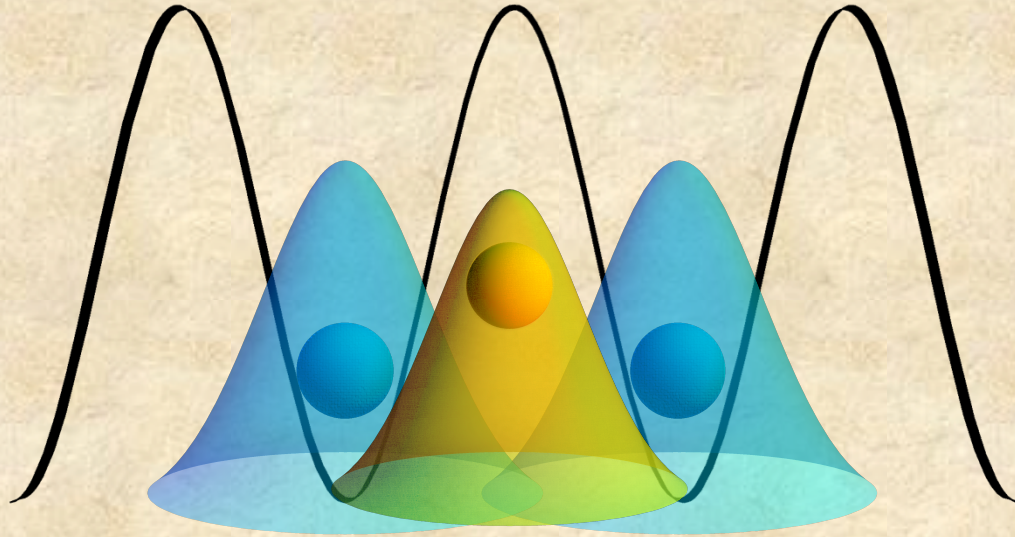
Step (IV):



Move the apparatus

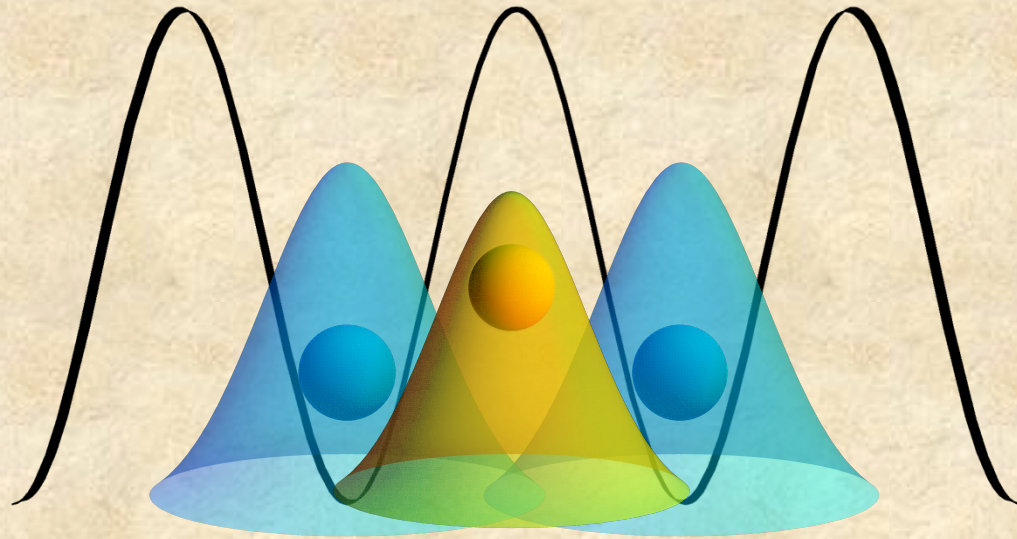
Relevant

$$\langle 0, 0, n_z | \hat{H}_{int} | \bar{0} \rangle \simeq g \sum_{i,j=0}^1 \int d\mathbf{x} \psi_{(0,0,n_z)}^*(\mathbf{x}) \psi_{(0,0,0)}(\mathbf{x}) \omega_i^*(x) \omega_j(x) \hat{c}_i^\dagger \hat{c}_j$$



Relevant

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Transition Probability

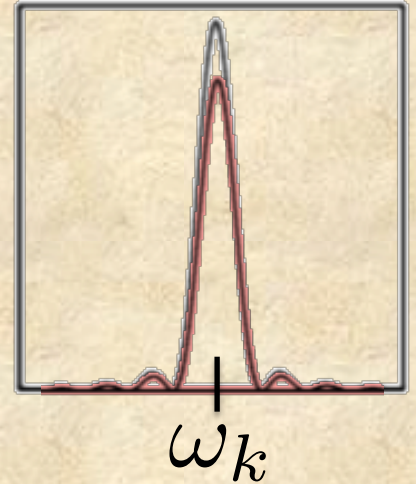
$$\Gamma_{\bar{0} \rightarrow (0,0,n_z)}^{\text{II}} = g_{\text{II},n}^2 \nu \left\{ 4n_0^2 \frac{1 - \cos(n\nu t)}{n^2 \nu^2} + \sum_k (1 + \cos(ka)) [\Gamma_k^+(n\nu, t) + \Gamma_k^-(n\nu, t)] \right\}$$

$$g_{\text{II},n}^2 = 2 \frac{g^2}{\nu} \left(\int dx \psi_{n_x=0}^2 \left(x - \frac{a}{2} \right) (\omega_0^2(x) + \omega_0(x) \omega_1(x)) \right)^2 \left(\sqrt{m} \frac{\gamma_0}{\pi} \sqrt{\nu_0} \right)^2 ((-1)^{n_z} \sqrt{m} \frac{\gamma_{n_z}^{1/2} \gamma_0^{1/2}}{\pi} \sqrt{\nu})^2$$

Resonant Probabilities

$$\Gamma_{\bar{0} \rightarrow (0,0,n_z)}^{\text{I}} \simeq 2 g_{\text{I},n}^2 \beta_k^2 n(\omega_k) T_f^2$$

$$\Gamma_{\bar{0} \rightarrow (0,0,n_z)}^{\text{II}} \simeq 2 g_{\text{II},n}^2 \beta_k^2 n(\omega_k) T_f^2 (1 + \cos(ka))$$



Resonant Probabilities



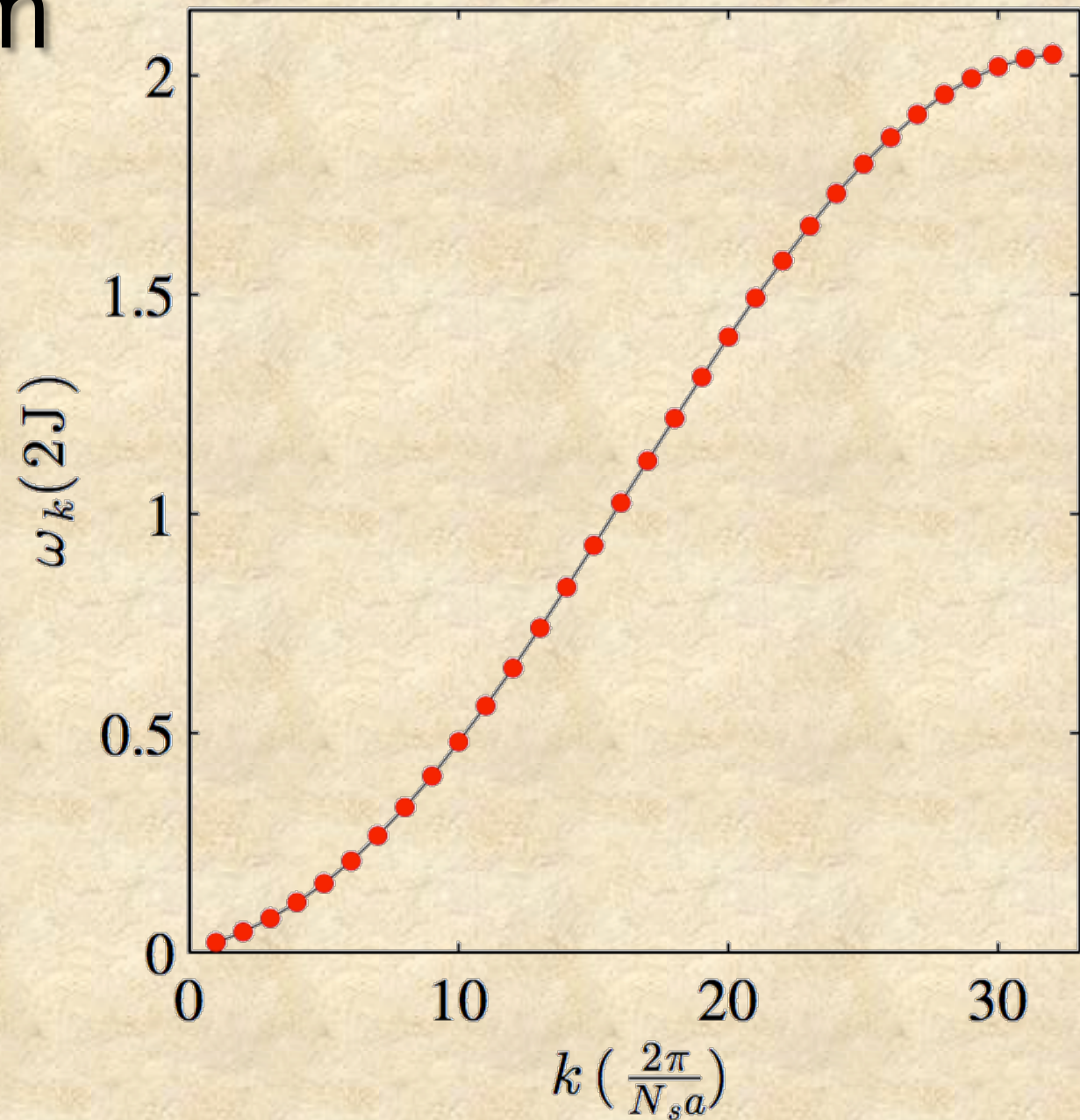
$$\Gamma_{0 \rightarrow (0,0,k_z)}^{\text{I}} \approx 2 g_{0,0}^2 \beta_k^2 n(\omega_k) T_f^2$$

$$\Gamma_{0 \rightarrow (0,0,k_z)}^{\text{II}} \approx 2 g_{0,0}^2 \beta_k^2 n(\omega_k) T_f^2 (1 + \cos(ka))$$

Ratio

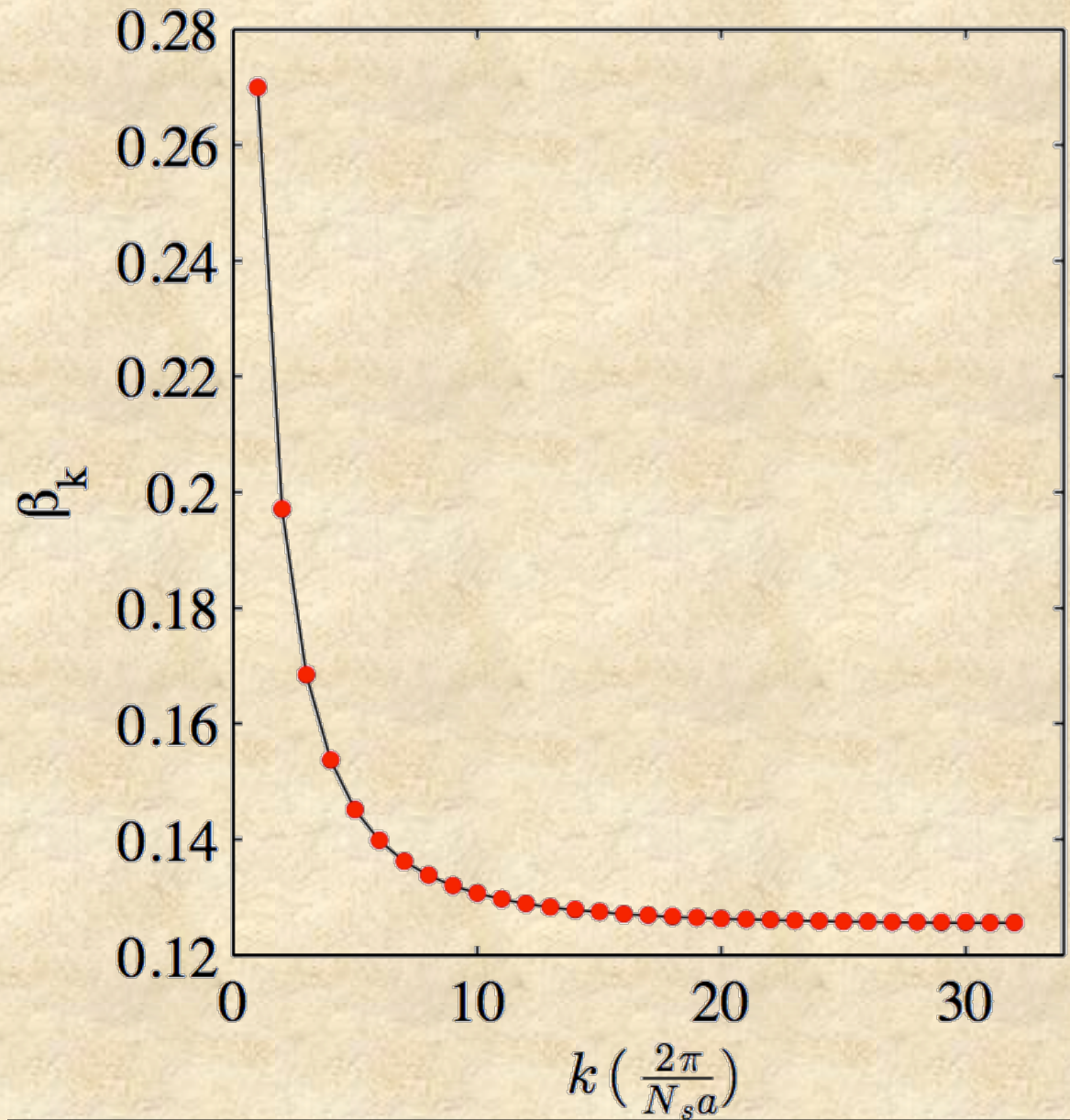
$$\frac{\Gamma^{\text{II}}}{\Gamma^{\text{I}}} \frac{g^{\text{I}}}{g^{\text{II}}} = [1 + \cos(ka)]$$

Spectrum

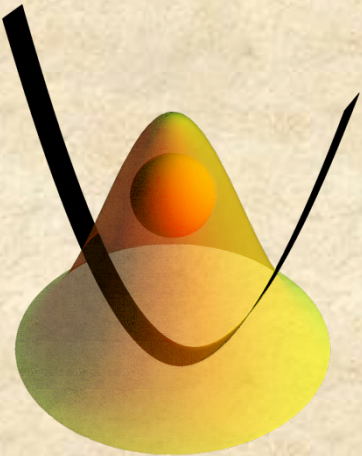


K-mode coupling

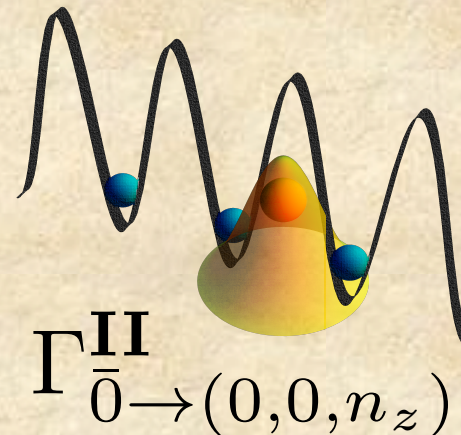
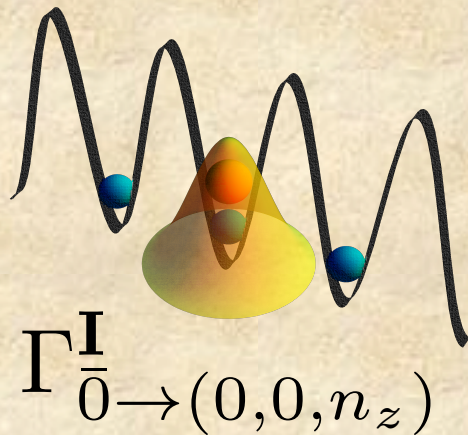
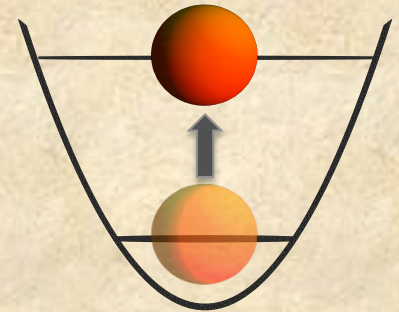
$$\Gamma_{\vec{0} \rightarrow (0,0,n_z)}^{\mathbf{I}} \simeq 2 g_{\mathbf{I},n}^2 \beta_k^2 n(\omega_k) T_f^2$$



Quantum Probe



$$\psi_0(x) \simeq e^{-\frac{x^2}{2x_0^2}} e^{-\frac{y^2}{2y_0^2}} e^{-\frac{z^2}{2z_0^2}}$$



$$\omega(k)$$

Outlook



New **MRS** protocol using atomic probes



Test the robustness of the protocol in a realistic situation



Design multiprobes scheme for the detection of quantum correlations



Investigate other geometries

**Thank you for the
attention!**