

Non-equilibrium thermodynamics in Coulomb Crystals



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QuProCS II
7/4/2017
IFISC (Mallorca)



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Outline



Intro to Coulomb Crystals



Non-equilibrium thermodynamics tools



Results

Coulomb Crystals

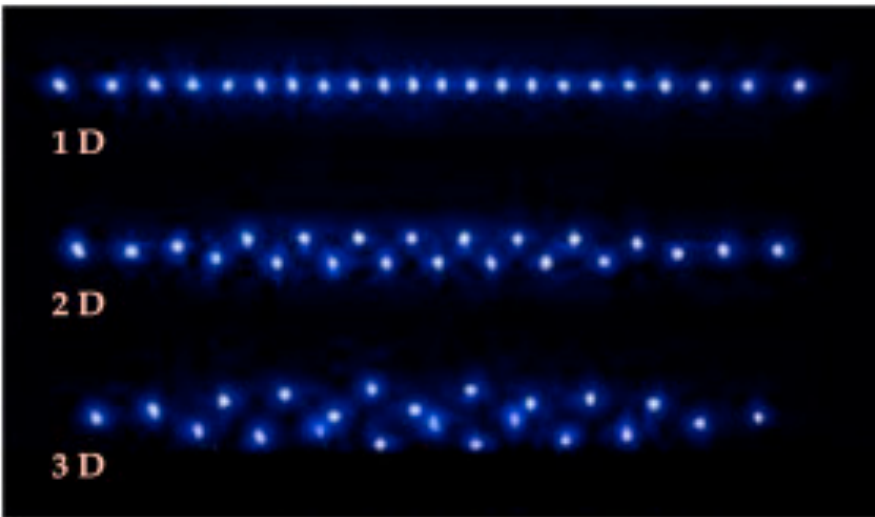
Trapped Ions

$$H_0 = \sum_{j=1}^L \left[\frac{P_{x,j}^2 + P_{y,j}^2}{2M} + \frac{M\omega_0^2}{2} Y_j^2 + V_L(X_j) \right] + \frac{Q^2}{8\pi\epsilon_0} \sum_{i \neq j} [(X_i - X_j)^2 + (Y_i - Y_j)^2]^{-1/2}$$

Interplay **confining potential**
and **Coulomb repulsion**

Coulomb Crystals

Structural phase transition



Linear Phase $\omega_0 > \omega_c$

ZigZag Phase $\omega_0 < \omega_c$

H. L. Partner et al., Physica B: Condensed Matter, Volume 460, 2015, 114–118,
<http://dx.doi.org/10.1016/j.physb.2014.11.051>

Semi-classical approach

Equilibrium Positions

ZigZag Phase

$$y_j = (-1)^j b/2$$

Linear Phase

$$y_j = 0$$

Critical Point

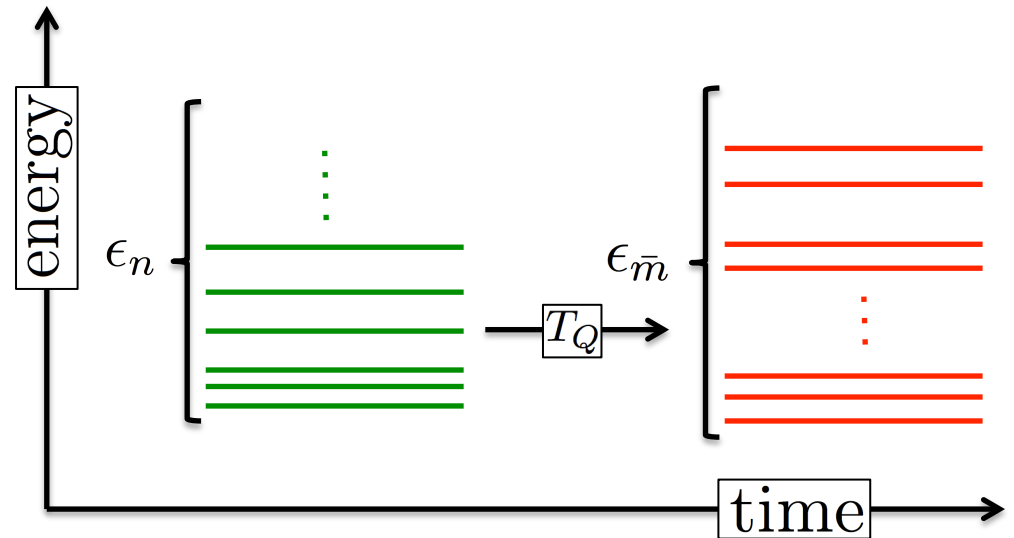
$$\omega_c = \sqrt{7\zeta(3)/2}$$

Excitations linear phase

$$\omega_k = \sqrt{\omega^2 - 4 \sum_{j=1}^{L/2} \frac{1}{j^3} \sin^2 \frac{jk}{2}}$$

Tools

Work distribution



$$P_F(W) \equiv \sum_{n, \bar{m}} p_n^0 p_{\bar{m}|n}^{t_f} \delta[W - (\epsilon_{\bar{m}} - \epsilon_n)]$$

Characteristic function

$$\chi_F(t) = \text{Tr}[e^{iH_f t} U(t_f, 0)^\dagger e^{-iH_i t} U(t_f, 0) \rho_i]$$

Crooks-Tasaki relation

$$\frac{P_F(W)}{P_B(-W)} = e^{\beta(W - \Delta F)}$$

Moments

$$M^{(n)} = (-i)^n \left. \frac{\partial^n \ln \chi_F(t)}{\partial t^n} \right|_{t=0}$$

$$M^{(1)} = \langle W \rangle$$

Measure of Irreversibility

$$\langle W_{IRR} \rangle = \langle W \rangle - \Delta F$$

Same phase quench

Transverse frequency quench

$$\hat{H}_i = V + \sum_k \hat{P}_k^2 + \frac{\omega_k^2(t)}{2} \hat{Q}_k^2$$

Time-dependency of the quench

$$\omega_k^2(t) = \omega_k^2(0) - \frac{t}{T_Q} (\omega_k^2(0) - \omega_k^2(T_Q))$$

Same phase quench

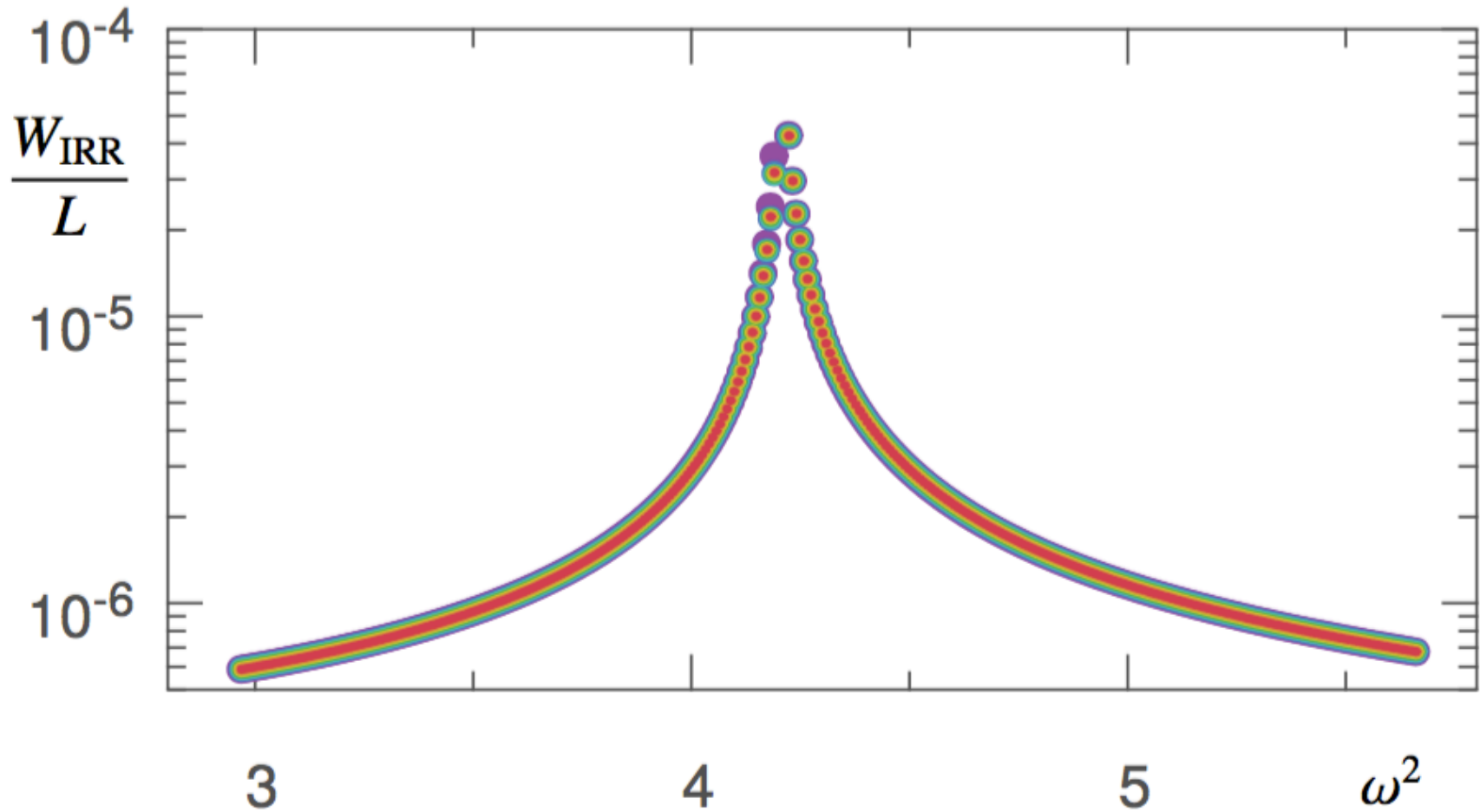
Characteristic function

$$\chi_F(t) = \prod_k \frac{\sqrt{2}(1 - e^{-\beta\omega_k(0)})e^{it\frac{(\omega_k(T_Q)\omega_k(0))}{2}}}{\sqrt{D_k(T_Q)}}$$

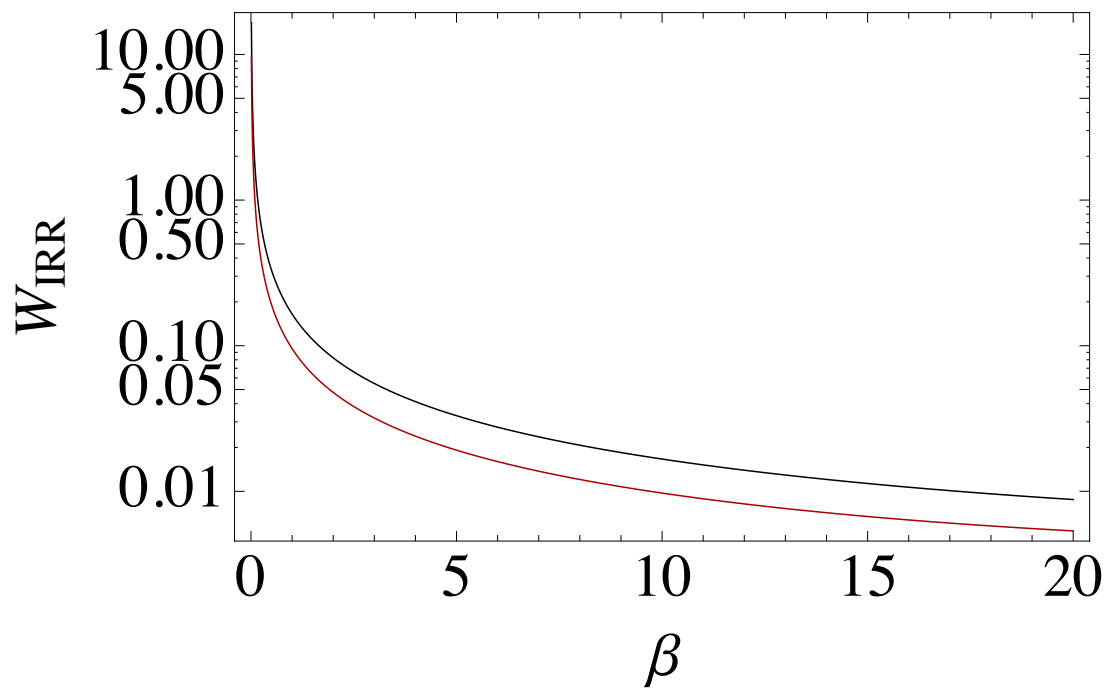
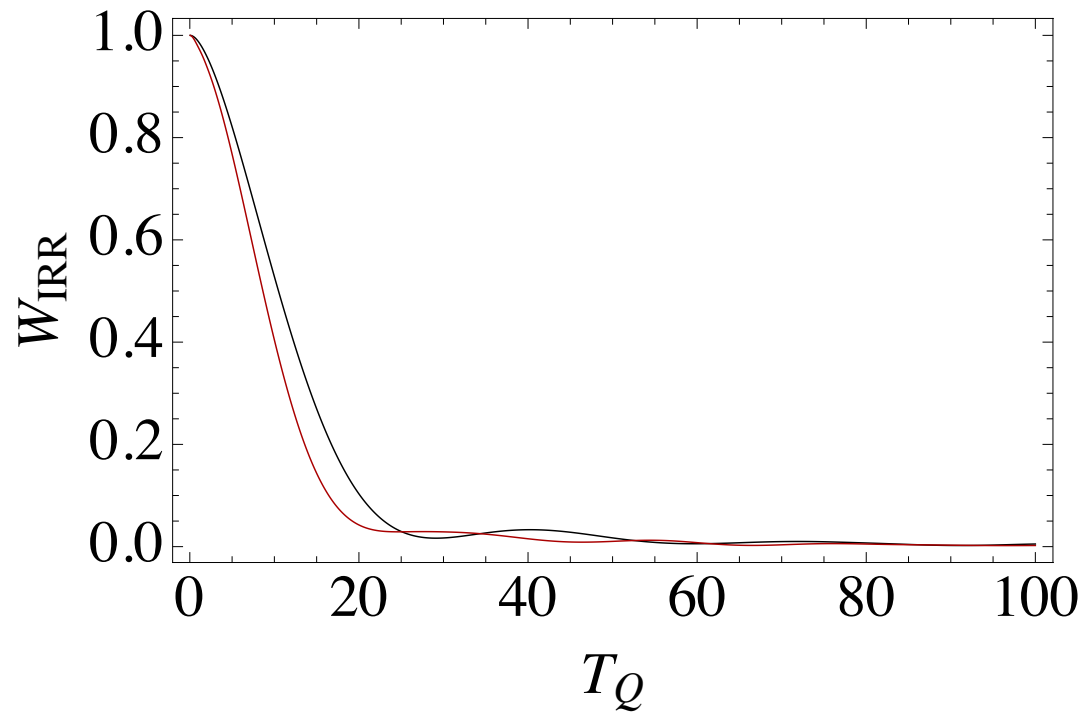
Irreversible work

$$\begin{aligned}\langle W_{IRR} \rangle &= \sum_k \frac{1}{2} (\Omega_k \omega_k(T_Q) - \omega_k(0)) \coth(\beta\omega_k(0)/2) \\ &\quad - \frac{1}{\beta} \ln \frac{\sinh(\beta\omega_k(T_Q)/2)}{\sinh(\beta\omega_k(0)/2)}\end{aligned}$$

Sudden Quench



Quenching time



Temperature

Anharmonic model: DMRG

Short range model

$$H(\omega) = \frac{1}{2} \sum_{j=1}^L \left(-g^2 \frac{\partial^2}{\partial y_j^2} + (\omega^2 - h_1) y_j^2 + h_2 (y_j + y_{j+1})^2 + h_3 y_j^4 \right)$$

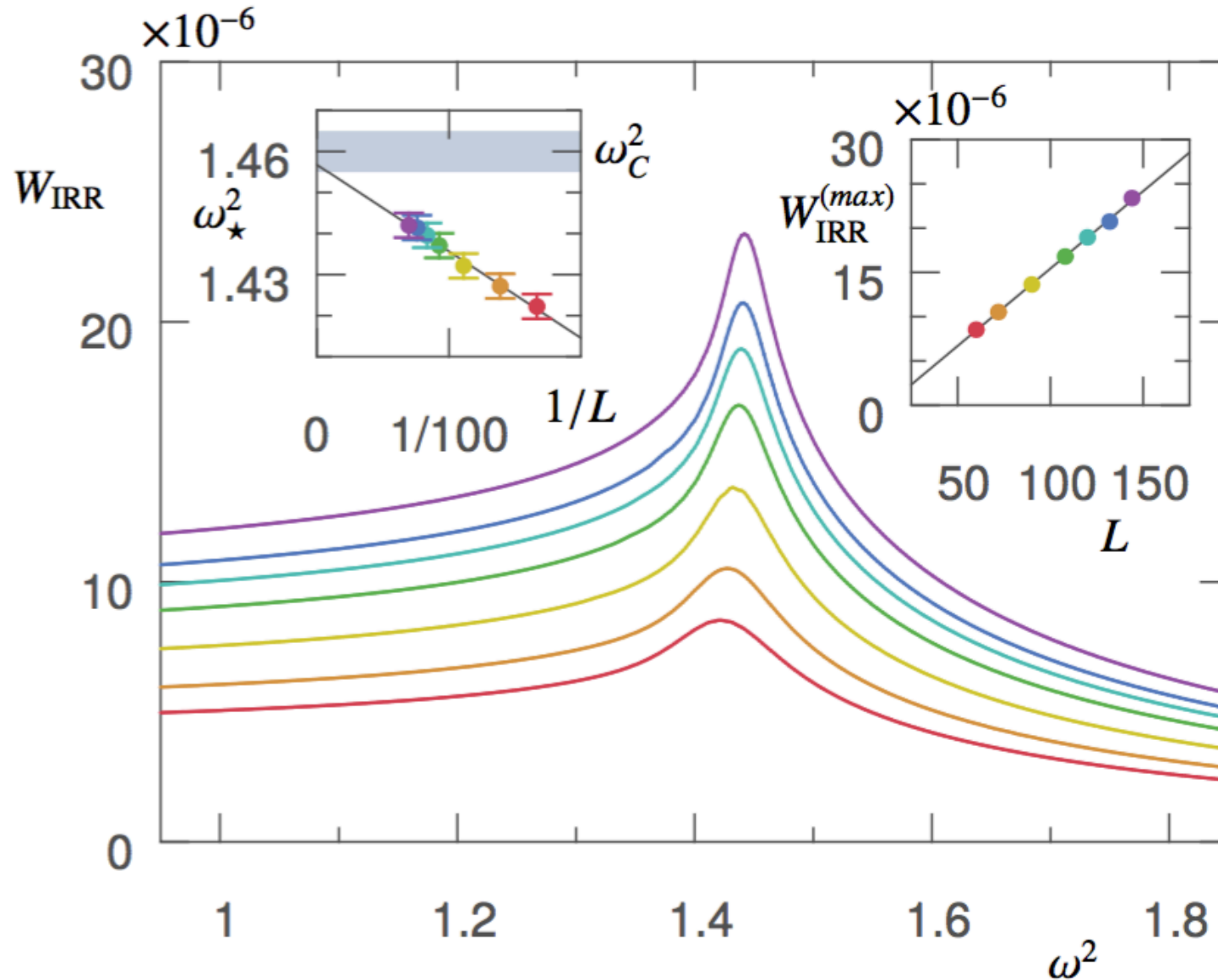
Effective Planck constant g $g = \sqrt{\frac{4\pi\epsilon_0\hbar^2}{MaQ^2}}$

Average work

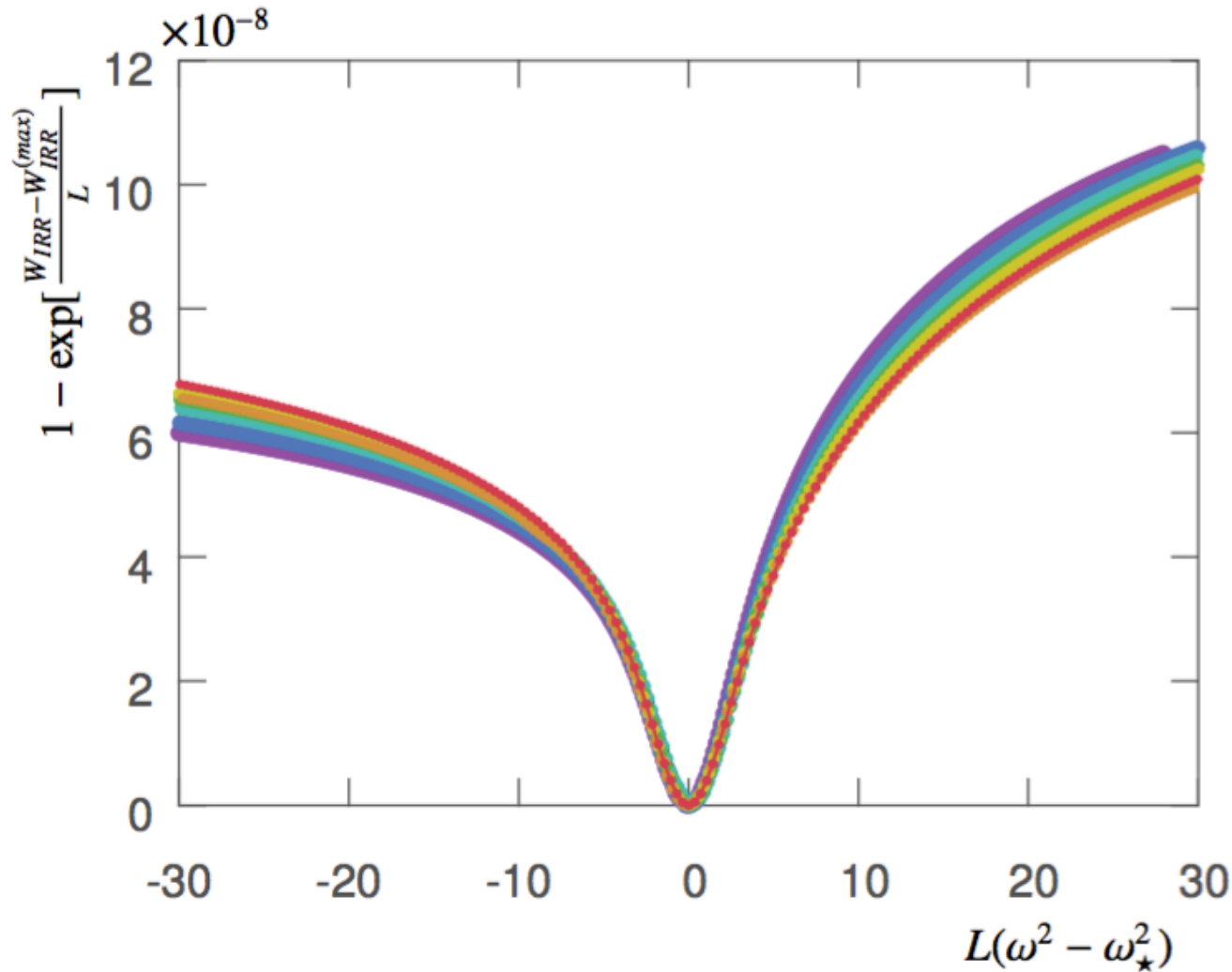
$$\langle W \rangle = E_G(\omega_f) + \frac{1}{2} (\omega_f^2 - \omega_i^2) \mathcal{Y}^2(\omega_i) - E_G(\omega_i)$$

$$\mathcal{Y}^2(\omega) = \sum_j \langle \Psi_G(\omega) | y_j^2 | \Psi_G(\omega) \rangle$$

Results: Zero temperature



Scaling: $1 - e^{\frac{W_{IRR} - W_{IRR}^{(max)}}{L}} = g(L^{\frac{1}{\nu}} (\omega^2 - \omega_{\star}^2))$



Ising model scaling

Conclusions and Future Prospectives



Characterization of the Irreversibility



Same scaling of the Ising model



Single ion as a probe

**Thank you for the
attention!**