

# Dynamics of a bosonic Josephson junction with an impurity

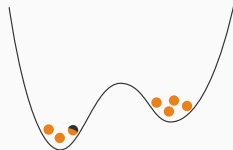
QuProCS II — IFISC (Mallorca)

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Maximilian Dirkmann, Gabriel Dufour, and Andreas Buchleitner

6-7/4/2017

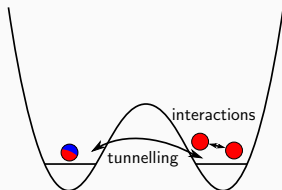
Albert-Ludwigs-Universität Freiburg



# Introduction

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# The Josephson junction with an impurity



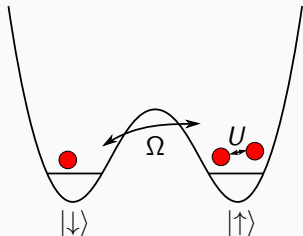
- Symmetric double well with one mode per site.
- Hopping from one mode to the other by tunnelling through the barrier.
- On-site interactions between the particles.
- All particles are indistinguishable bosons except for the impurity.

*Quantum probing*

# Single species system

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# Bose-Hubbard model<sup>1</sup> and Schwinger representation<sup>2</sup>



$$\hat{J}_x \equiv (\hat{a}_R^\dagger \hat{a}_L + \hat{a}_L^\dagger \hat{a}_R)/2,$$

$$\hat{J}_y \equiv (\hat{a}_R^\dagger \hat{a}_L - \hat{a}_L^\dagger \hat{a}_R)/2i,$$

$$\hat{J}_z \equiv (\hat{a}_R^\dagger \hat{a}_R - \hat{a}_L^\dagger \hat{a}_L)/2 = (\hat{n}_R - \hat{n}_L)/2$$

$$\hat{H} = -\Omega \hat{J}_x + U \hat{J}_z^2$$

$|j, m\rangle$  as eigenvectors of:

$$\hat{\mathbf{J}}^2 |j, m\rangle = j(j+1) |j, m\rangle, \text{ spin } j = \frac{N}{2}$$

$$\hat{J}_z |j, m\rangle = m |j, m\rangle, \text{ spin magnetic quantum number,}$$

$m \in \{-j, -j+1, \dots, j-1, j\} \hat{=} \text{population imbalance,}$

$$(\hat{n}_R - \hat{n}_L) |j, m\rangle = 2\hat{J}_z |j, m\rangle = 2m |j, m\rangle$$

<sup>1</sup>G. J. Milburn et al., *Phys. Rev. A* **55**, 4318–4324 (1997).

<sup>2</sup>B.J. Dalton and S. Ghanbari, *J. Mod. Opt.* **59**, 287–353 (2012),

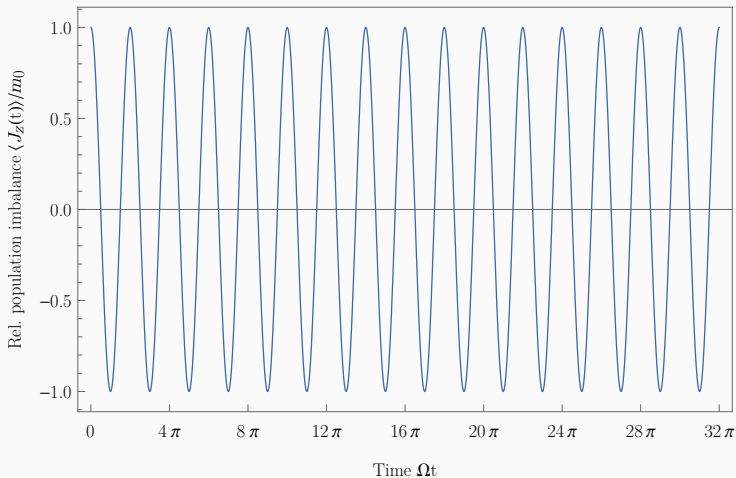
J. Schwinger, *US Atomic Energy Commission Report No. NYO-3071*, 229 (1952).

# Single species system

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Dynamical regimes

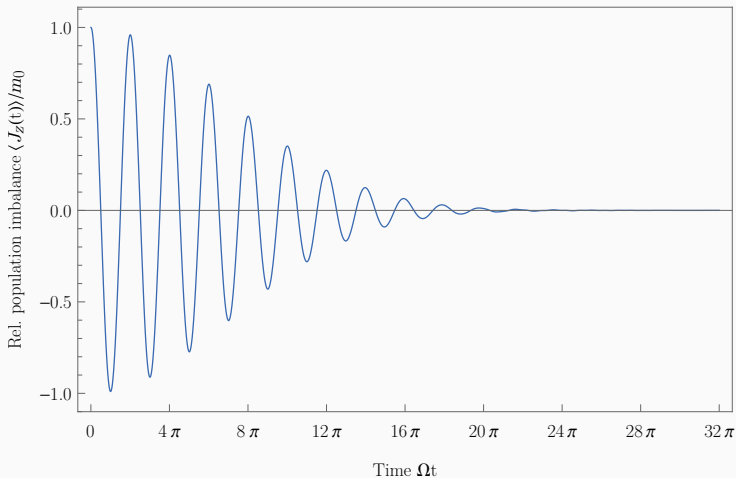
# Dynamical Regimes<sup>3</sup>



$N = 10$  particles, initial imbalance  $m_0 = 5$  and  $U = 0 \frac{\Omega}{N}$

<sup>3</sup>A. J. Leggett, *Rev. Mod. Phys.* **73**, 307–356 (2001).

# Dynamical Regimes<sup>3</sup>

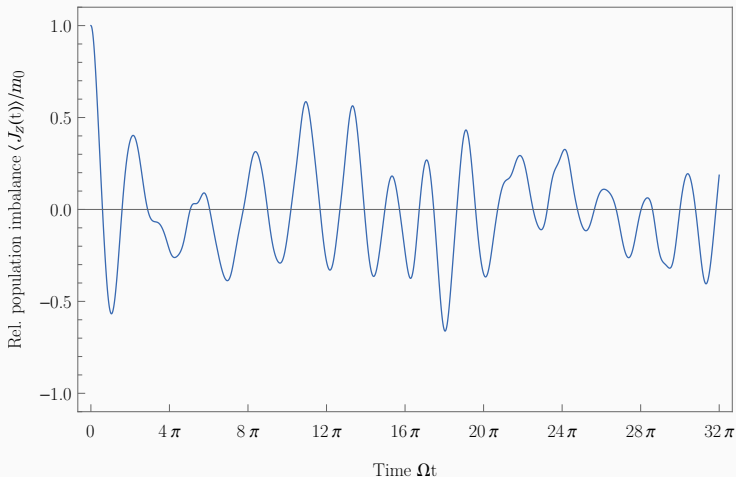


$N = 10$  particles, initial imbalance  $m_0 = 5$  and  $U = 0.3 \frac{\Omega}{N}$

<sup>3</sup>A. J. Leggett, *Rev. Mod. Phys.* **73**, 307–356 (2001).



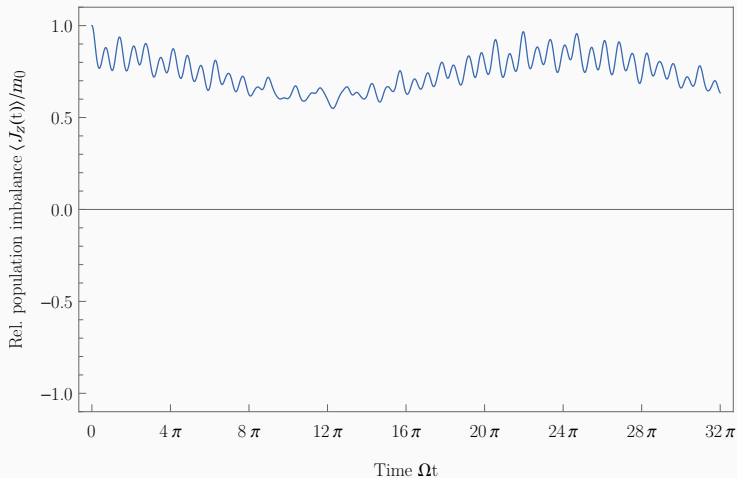
# Dynamical Regimes<sup>3</sup>



$N = 10$  particles, initial imbalance  $m_0 = 5$  and  $U = 1.5 \frac{\Omega}{N}$

<sup>3</sup>A. J. Leggett, *Rev. Mod. Phys.* **73**, 307–356 (2001).

# Dynamical Regimes<sup>3</sup>



$N = 10$  particles, initial imbalance  $m_0 = 5$  and  $U = 3.5 \frac{\Omega}{N}$

<sup>3</sup>A. J. Leggett, *Rev. Mod. Phys.* **73**, 307–356 (2001).

## Two species system

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## Two species Hamiltonian

- Tensor product structure  $\mathcal{H}_1 \otimes \mathcal{H}_2$
- New families of operators  $\hat{\mathbf{J}}_1, \hat{\mathbf{J}}_2$
- States are written as  $|j_1, m_1\rangle \otimes |j_2, m_2\rangle = |j_1, m_1, j_2, m_2\rangle$

$$\hat{H} = -\Omega_1 \hat{J}_{x1} - \Omega_2 \hat{J}_{x2} + U_{1,1} \hat{J}_{z1}^2 + U_{2,2} \hat{J}_{z2}^2 + 2U_{1,2} \hat{J}_{z1} \hat{J}_{z2}$$

### Isospecific conditions

$$\Omega = \Omega_1 = \Omega_2 \text{ and } U = U_{1,1} = U_{2,2} = U_{1,2}$$

## **Two species system**

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**Mixing behaviour of the impurity with the majority species**

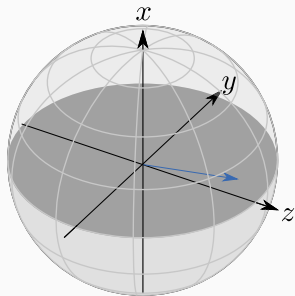
# Purity of the impurity's state

- Reduced density matrix of the impurity

$$\hat{\rho}_2(t) = \text{Tr}_1 \{ \hat{\rho}(t) \} = \frac{1}{2} (\mathbb{1} + \mathbf{s}(t) \cdot \boldsymbol{\sigma})$$

- Purity

$$P_2(t) = \text{Tr} \{ \hat{\rho}_2(t)^2 \} = \frac{1 + \|\mathbf{s}(t)\|^2}{2}$$



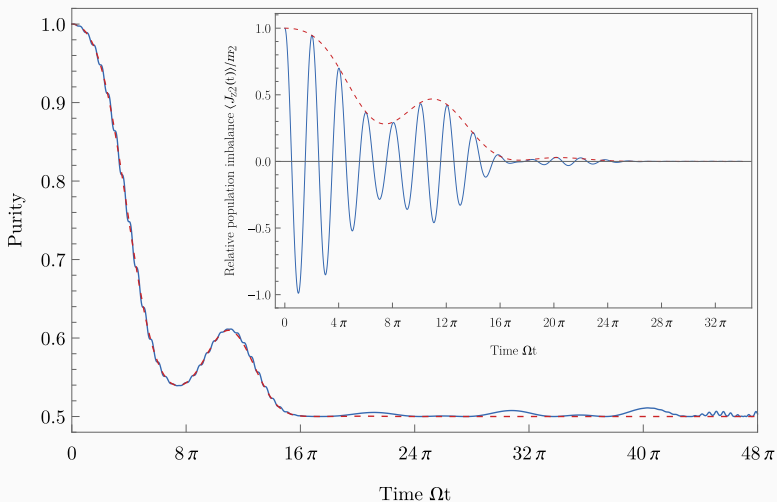
Weak interactions:

$$\mathbf{s}(t) = \frac{1}{2} \|\mathbf{s}(t)\| (\cos(\Omega t) \mathbf{e}_z + \sin(\Omega t) \mathbf{e}_y)$$

$$\langle \hat{J}_{z,2}(t) \rangle = \frac{1}{2} \|\mathbf{s}(t)\| \cos(\Omega t)$$

- Link between envelope function and purity

# Link between envelope function and purity

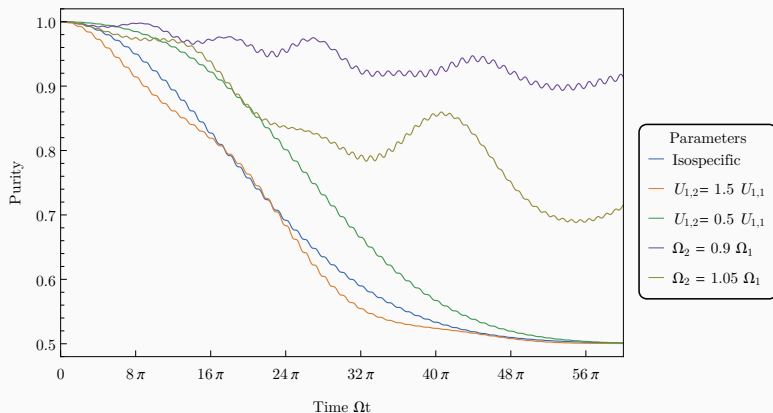


— numerically exact solution

- - - approximation through the envelope function

$$N_1 = 20 \text{ particles, } U_{1,1} = U_{1,2} = 0.4 \frac{\Omega}{N}.$$

# Outlook: Purity for non-isospecific conditions



$N_1 = 15$ ,  $U_{1,1} = 0.1 \frac{\Omega_1}{N}$ , initially all particles started in the same well.



## **Summary and conclusions**

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## Summary and conclusions

- Weak interactions introduce an envelope to the unhindered Rabi oscillations
- Direct link between purity and oscillations of the impurity particle
- "Quantum probing" possible through measuring the impurity's oscillations

## Questions and discussion